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Brownian Excursions From Extremes

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Let $B = (B_t, \mathcal{F}_t, P; t \geq 0)$ be standard Brownian motion starting at zero and define its extreme processes as

$$M_t = \max_{0 \le s \le t} B_s$$
 and $m_t = \min_{0 \le s \le t} B_s$.

The point of this note is to observe a mapping property of Brownian motion and use it to derive some results about excursions of B from its extremes which are related to the work of Groeneboom [4], Bass[1] and Pitman[9] and of Imhof[7]. It must be pointed out that these results are consequences of general excursion theory as expounded by Getoor [2],[3] and Jacobs[8], for example. However this mapping property is new and its application to excursions is direct.

Let $r_t = M_t - m_t$ be the range process and for each $\epsilon > 0$ define the increasing processes

$$a(t,\epsilon)=\int_{\epsilon}^{t}4r_{s}^{-2}ds$$

and

$$\tau(t,\epsilon) = \inf\{s: a(s,\epsilon) > t\}.$$

Let

$$X_t = \frac{2B_t - M_t - m_t}{M_t - m_t}$$

and define $X_t^{\epsilon} = X_{\tau(t,\epsilon)}$.

Proposition 1. The process $X^{\epsilon} = (X_t^{\epsilon}, \mathcal{F}_{\tau(t,\epsilon)}, P; t \geq 0)$ is a reflecting Brownian motion on [-1,1]. Its local times at ± 1 are

$$\phi_t^{\epsilon,+} = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-1} dM_s$$

and

$$\phi_t^{\epsilon,-} = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-1} d(-m_s)$$
 respectively

Proof. We may write equation (1) as $X_t = F(B_t, M_t, m_t)$ where

$$F(x,y,z) = (2x - y - z)/(y - z).$$

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Since F is smooth on $\{y \neq z\}$, we may apply Itô's formula there to obtain

(2)
$$dX_s = \frac{2dB_s}{M_s - m_s} + \frac{2(M_s - B_s)d(-m_s)}{(M_s - m_s)^2} - \frac{2(B_s - m_s)dM_s}{(M_s - m_s)^2}.$$

Because each $\tau(t,\epsilon)$ is an \mathcal{F}_{t} -stopping time we may write (2) in the integrated form

$$X_{\tau(t,\epsilon)} = X_{\epsilon} + W_{t}^{\epsilon} + \frac{1}{2}\phi_{t}^{\epsilon,-} - \frac{1}{2}\phi_{t}^{\epsilon,+}$$

where

$$\begin{cases} W_t^{\epsilon} = \int_{\epsilon}^{\tau(t,\epsilon)} 2r_s^{-1} dB_s \\ \phi_t^{\epsilon,-} = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-1} d(-m_s) \\ \phi_t^{\epsilon,+} = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-1} dM_s \end{cases}$$

To finish the proof we check that W^{ϵ} is a standard Brownian motion and that (3) is its Skorohod equation (Tanaka[11]). Clearly W^{ϵ} is an $\mathcal{F}_{\tau(t,\epsilon)}$ -martingale and

$$[W^{\epsilon}]_t = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-2} ds = a(\tau(t,\epsilon),\epsilon) = t.$$

By Lévy's criterion, W^{ϵ} is a Brownain motion, independent of X_{ϵ} . Now $\phi_t^{\epsilon,\pm}$ are continuous, incereasing $\mathcal{F}_{\tau(t,\epsilon)}$ -adapted processes which increase only when B attains a new extremum, that is only when $X^{\epsilon}=\pm 1$. \diamond

As the next propostion shows, we may write the extreme processes in terms of the local times $\phi^{\epsilon,\pm}$. Set $\phi^{\epsilon}_t = \phi^{\epsilon,+}_t + \phi^{\epsilon,-}_t$.

Proposition 2. For $t \geq \epsilon$,

(i)
$$M_t = M_{\epsilon} + r(\epsilon) \int_0^{a(t,\epsilon)} \exp \left\{ \phi_s^{\epsilon}/4 \right\} d\phi_s^{\epsilon,+}$$

(ii)
$$m_t = m_{\epsilon} + r(\epsilon) \int_0^{a(t,\epsilon)} \exp \left\{ \phi_s^{\epsilon}/4 \right\} d\phi_s^{\epsilon,-}$$

where $a(t, \epsilon) = \inf\{s : \tau(s, \epsilon) > t\}$ and

$$au(t,\epsilon) = \epsilon + rac{1}{4}r(\epsilon)^2 \int_0^t \exp\left\{\phi_s^\epsilon/2\right\} ds.$$

Proof. By (4) we have

$$\phi_t^{\epsilon} = \int_{\epsilon}^{\tau(t,\epsilon)} 4r_s^{-1} dr_s = 4 \log \frac{r(\tau(t,\epsilon))}{r(\epsilon)},$$

hence

(5)
$$r(\tau(t,\epsilon)) = r(\epsilon) \exp \left\{ \phi_t^{\epsilon} / 4 \right\}$$

Since $a(\tau(t,\epsilon),\epsilon)=t$ it follows that $d\tau(t,\epsilon)=r(\tau(t,\epsilon))^2dt/4$. Thus by (5),

$$egin{aligned} au(t,\epsilon) &= & au(0,\epsilon) + rac{1}{4} \int_0^t r(au(t,\epsilon))^2 \, ds \ &= &\epsilon + rac{1}{4} r(\epsilon)^2 \int_0^t \exp\left\{\phi_s^\epsilon/2\right\} ds. \end{aligned}$$

It follows that τ and hence a are defined solely in terms of M_{ϵ} , m_{ϵ} and $\phi^{\epsilon,\pm}$. Next, by (4)

(6)
$$\phi_t^{\epsilon,+} = \int_0^{\tau(t,\epsilon)} 4r_s^{-1} dM_s = \int_0^t 4r(\tau(s,\epsilon))^{-1} dM_{\tau(s,\epsilon)},$$

and so

(7a)
$$M_{\tau(t,\epsilon)} = M_{\epsilon} + \frac{1}{4} r_{\epsilon} \int_{0}^{t} \exp\left\{\phi_{s}^{\epsilon}/4\right\} d\phi_{s}^{\epsilon,+},$$

Similarly, we have

(7b)
$$m_{\tau(t,\epsilon)} = m_{\epsilon} + \frac{1}{4} r_{\epsilon} \int_{0}^{t} \exp \left\{ \phi_{s}^{\epsilon}/4 \right\} d\phi_{s}^{\epsilon,-},$$

and the proposition follows from a time change in (7a) and (7b).

These propositions allow us to compare excursions of B from its extremes with excursions of reflecting Brownian motion in [-1,1]. To be precise, let

(8)
$$f(t,\epsilon) = \inf\{s : \phi_s^{\epsilon} > t\}$$

be the inverse of boundary local time and let q^{ϵ} be the point process of excursions of X^{ϵ} . That is, let

$$D_{q^{\epsilon}} = \{s: f(s, \epsilon) > f(s-, \epsilon)\}$$

and for each $s \in D_{q'}$ let

$$q_s^{\epsilon}(u) = X^{\epsilon}(f(s-,\epsilon) + u \wedge l_s^{\epsilon}), \qquad u \geq 0$$

where $l_s^{\epsilon} = f(s, \epsilon) - f(s-, \epsilon)$ is the duration of the excursion. Similarly, consider the point process p of excursions of B from its extremes. Let

(10)
$$\mu(t) = \inf\{s : r(s) > t\},\,$$

let the domain of p be $D=\{t: \mu(t)>\mu(t-)\}$ and for each $t\in D$ let

$$(11) p_t(u) = B(\mu(t-) + u \wedge \lambda(t)), u \geq 0$$

where $\lambda(t) = \mu(t) - \mu(t-)$. Proposition 4 provides a formula for p in terms of q^{ϵ} . To ensure the formula is well defined we need the

Lemma 3. Let $D^{\epsilon} = \{t : \mu(t) > \mu(t-) \text{ and } \mu(t) > \epsilon\}$. Then

(i) $f(t,\epsilon) = a(\mu(r(\epsilon)e^{t/4}),\epsilon)$

(ii)
$$D_{q'} = \{s(u,\epsilon) : u \in D^{\epsilon}\}$$
 where $s(t,\epsilon) = 4 \log(tr(\epsilon)^{-1})$.

Proof. The lemma follows easily from the equality $\tau(f(t,\epsilon),\epsilon) = \mu(r(\epsilon)e^{t/4})$, which we now show. Let $\alpha(t)$ and $\beta(t)$ denote the left and right side of this equality, respectively. On the one hand, by (7a) and (7b), we have

$$r(\alpha(t)) = r(\epsilon) \exp\{\phi^{\epsilon}(f(t,\epsilon))\} = r(\epsilon)e^{t/4} \stackrel{\text{def.}}{=} g(t).$$

On the other hand, by definition $r(\beta(t)) = g(t)$. Thus $r(\alpha(t)) = r(\beta(t))$. Since g(t) is strictly increasing, for any $\delta > 0$

$$\alpha(t+\delta) \ge \mu(r(\alpha(t+\delta))-) = \mu(g(t+\delta)-)$$

$$\ge \mu(g(t)) = \mu(r(\beta(t))) = \beta(t).$$

Letting $\delta \downarrow 0$ we get $\alpha(t) \geq \beta(t)$. Since the reverse inequality is similar, the lemma is proved.

Proposition 4. Let $\{p_t; t \in D\}$ be the point process of excursions of B from its extremes. For each $t \in D^{\epsilon}$

$$p_t(u) = \frac{t}{2}q_{s(t,\epsilon)}^{\epsilon}\left(\frac{4u}{t^2}\right) + \frac{1}{2}\left(M_{\mu(t)} + m_{\mu(t)}\right)$$

Proof. First note that the statement makes sense, by Lemma 3. Let $s \in D_{q'}$ where $s = s(t, \epsilon)$ and $t \in D^{\epsilon}$. The durations $l^{\epsilon}(s)$ of q_s^{ϵ} and $\lambda(t)$ of p_t are related, according to Lemma 3, by

(12)
$$l^{\epsilon}(s) = f(s(t,\epsilon),\epsilon) - f(s(t,\epsilon),\epsilon)$$
$$= a(\mu(t),\epsilon) - a(\mu(t),\epsilon)$$
$$= 4 \int_{\mu(t-)}^{\mu(t)} r_{u}^{-2} du$$
$$= \frac{\mu(t) - \mu(t-)}{r(\mu(t))^{2}} = 4 \frac{\lambda(t)}{t^{2}}.$$

Thus by the formulas

$$X_{u} = \frac{2B_{u} - M_{u} - m_{u}}{M_{u} - m_{u}}, \qquad X_{v}^{\epsilon} = X_{\tau(v,\epsilon)}$$

and the definition of q^{ϵ} and p we get

$$q_{s(t,\epsilon)}^{\epsilon}(u) = rac{1}{t}\left(2p_t\left(rac{t^2u}{4}
ight) - M_{\mu(t)} - m_{\mu(t)}
ight),$$

from which the proposition follows. >

An immediate corollary is the identification of the conditional law of excursions of B from its extremes. Indeed, let $-\infty < c < d < \infty$ and introduce the transition density of Brownian motion in [c,d] with absorption at the endpoints (Port-Stone[10]):

$$(14) p_0^{c,d}(t,x,y) = \frac{2}{d-c} \sum_{n=0}^{\infty} \sin\left(n\pi \frac{x-c}{d-c}\right) \sin\left(n\pi \frac{y-c}{d-c}\right) \exp\left\{-\frac{n^2\pi^2}{(d-c)^2} \frac{t}{2}\right\}$$

as well as the functions

(15)
$$\begin{cases} g^{c,d}(t,y;a) = \frac{1}{2} \frac{\partial}{\partial n_a} p_0^{c,d}(t,a,y), & a = c,d \\ \theta^{c,d}(t,a,b) = \frac{1}{4} \frac{\partial^2}{\partial n_a \partial n_b} p_0^{c,d}(t,a,b), & a,b = c,d. \end{cases}$$

There exist unique probability laws $P_{c,d}^{a,b;l}$ on $C([0,\infty),[c,d])$ with absolute distribution:

(16)
$$P_{c,d}^{a,b;l}(e(u) \in dy) = \frac{g^{c,d}(u,y;a)g^{c,d}(l-u,y;b)}{\theta^{c,d}(l,a,b)}dy, \qquad 0 \le u \le l$$

and transition density

$$(17) \quad P_{c,d}^{a,b;l}\left(e(v) \in dy \middle| e(u) = x\right) = p_0^{c,d}(v-u,x,dy) \frac{g^{c,d}(l-v,y;b)}{g^{c,d}(l-u,x;b)} \qquad 0 \leq u < v \leq l.$$

Indeed, if $X^{c,d}$ is reflecting Brownian motion in [c,d] then $P_{c,d}^{a,b;l}$ is just the law of the excursion process of $X^{c,d}$ conditioned to begin at a, end at b and have duration l. This is a simple extension of the well-known case of one reflecting barrier (e.g. Ikeda-Watanabe[6]) and also can be proved by imitating the calculations of Hsu[5]. Finally let us note a scaling property of the laws $P_{c,d}^{a,b;l}$ which follows from the invariance of the family $\{p_0^{c,d}, -\infty < c < d < \infty\}$ under affine changes of variable:

(18) If $Z = \{Z(t); 0 \le t \le l\}$ has the law $P_{c,d}^{a,b;l}$ then $\{\alpha Z(\alpha^{-2}t) + \beta; 0 \le t \le \alpha^2 l\}$ has the law $P_{\alpha c + \beta, \alpha b + \beta; \alpha^2 l}^{\alpha a + \beta, \alpha b + \beta; \alpha^2 l}$.

Theorem 5. Let $t \in D$. Let $-\infty < c < d < \infty$ and let l > 0. Then conditional on the event $\xi = [m_{\mu(t)} = c, M_{\mu(t)} = d, p_t(0) = a, p_t(\lambda(t)) = b, \lambda(t) = l]$, the law of the excursion process $p_t(\cdot)$ is $P_{c,d}^{a,b;l}$.

Proof. Fix some ϵ with $t \in D^{\epsilon}$ and let $s = s(t, \epsilon)$. By (9) and (12), we have

$$\xi = [q_s^{\epsilon}(0) = \operatorname{sgn}(a), q_s^{\epsilon}(l^{\epsilon}(s)) = \operatorname{sgn}(b), l^{\epsilon}(s) = |d - c|^2 l/4]$$

But then conditional on ξ , the process $q_s^{\epsilon}(\cdot)$ has law $P_{-1,1}^{e,f;m}$ with $e = \operatorname{sgn}(a), f = \operatorname{sgn}(b)$ and $m = |d - c|^2 l/4$. So by Proposition 4 and the invariance property (18) we find that conditional on ξ , $p_t(\cdot)$ has the law $P_{c,d}^{a,b;l}$. \diamond

It is known that if X is reflecting Brownian motion in an interval then conditional on the σ -field generated by the boundary local time of X, the various excursions of X from the boundary are mutually independent. This is evident from the construction of the excursions law characterizing the excursion point process in the one reflecting barrier case (Ikeda-Watanabe[6]). Or again, one can either imitate the argument of Hsu[5] or simply quote the results in Jacobs[8]. Let us show that this conditional independence property is shared by excursions of Brownian motion B from its extremes, conditional on $\sigma\{M_s, m_s; s \geq 0\}$.

Lemma 6. Let $\mathcal{B}_{\epsilon} = \sigma\{\phi_s^{\epsilon,+}, \phi_s^{\epsilon,-}; s \geq 0\}$ and $\mathcal{B} = \sigma\{M_s, m_s; s \geq 0\}$. Then $\mathcal{B}_{\epsilon} \subset \mathcal{B}$ and $\lim_{\epsilon \to 0} \mathcal{B}_{\epsilon} = \mathcal{B}$.

Proof. Since Proposition 2 exhibits M and m as explicit functions of $\phi^{\epsilon,\pm}$, we have the inclusions

$$\sigma\{M_s-M_\epsilon,m_s-m_\epsilon;s\geq\epsilon\}\subset\sigma\{M_\epsilon,m_\epsilon,\phi_s^{\epsilon,\pm};s\geq0\}\subset\sigma\{M_s,m_s;s\geq\epsilon\}$$

and the lemma follows from this.

Theorem 7. Conditional on $\beta = \sigma\{M_s, m_s; s \geq 0\}$, the excursions $[p_t(\cdot); t \in D]$ are mutually independent.

Proof. For $n \geq 1$ consider functionals $F: C([0,\infty),R)^n \to R$ of the form

$$F(\omega_1, \omega_2, \ldots, \omega_n) = \prod_{j=1}^n f_j(\omega_j(s_{j,1}), \ldots, \omega_j(s_{j,m(j)}))$$

for bounded continuous functions f_j . Let $t_1, \ldots, t_n \in D$. Using Proposition 1, for all sufficiently small ϵ ,

$$E\left[F(p_{t_1},\ldots,p_{t_n})\middle|\mathcal{B}_{\epsilon}\right] = \prod_{j=1}^n E\left[f_j(p_{t_j}(s_{j,1}),\ldots,p_{t_j}(s_{j,m(j)}))\middle|\mathcal{B}_{\epsilon}\right]$$

by the conditional independence property of q^{ϵ} . Thus by the martingale convergence theorems and Lemma 6; taking the limit as $\epsilon \downarrow 0$ yields

$$E\left[F(p_{t_1},\ldots,p_{t_n})\middle|\mathcal{B}\right] = \prod_{j=1}^n E\left[f_j(p_{t_j}(s_{j,1}),\ldots,p_{t_j}(s_{j,m(j)})\middle|\mathcal{B}\right]. \quad \diamond$$

We close by remarking that Theorem 5 and 7 show that Brownian motion consists of conditionally independent Brownian excursions properly interpolated between endpoints of flat stretches of the extreme process M and m.

References

- [1] Bass, R.F., Markov Processes and Convex Minorants, Sem. Prob. XVIII, Lecture Notes in Math. No. 1059, 24-41.
- [2] Getoor, R.K., Splitting Times and Shift Functionals, Z. fur Wahr. verw. Gebiete, 47(1979), 69-81.
- [3] Getoor, R.K., Excursions of Markov Processes, Ann. Prob. 7(1979), 244-266.
- [4] Groeneboom, P., The Concave Majorant of Brownian Motion, Ann. Prob. 6(1983), 1016-1027.
- [5] Hsu, P., On Excursions of Reflecting Brownian Motion, Trans. Amer. Math. Soc., vol.296, No.1(1986), p.239-263.
- [6] Ikeda, N. and Watanabe, S., Stochastic Differential Equations and Diffusion Processes, North Holland and Kodansha, Amsterdam, Oxford, New York, 1981.
- [7] Imhof, J-P., On the Range of Brownian Motion and its Inverse Process, Ann. Prob. 13(1985), 1011-1017.
- [8] Jacobs, P.A., Excursions of a Markov Process Induced by Continuous Additive Functionals, Z. fur Wahr. verw. Gebiete, 44 (1978), 325-336.
- [9] Pitman, J.W., Remarks on the Convex Minorant of Brownian Motion, Seminar on Stochastic Processes, Birkhäuser, 1982.
- [10] Port, S.C. and Stone, C.J., Brownian Motion and Classical Potential Theory, Academic Press, New York, 1978.
- [11] Tanaka, H., Stochastic Differential Equations with Reflecting Boundary Conditions in Convex Domains, Hiroshima Math. J. 9(1979), 163-179.