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KALYANAPURAM RANGACHARI PARTHASARATHY

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A remark on the paper " Une martingale  
d'opérateurs bornés, non représentable en intégrale  
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by

K.R. Parthasarathy

Indian Statistical Institute, Delhi Centre  
7, S.J.S. Sansanwal Marg, New Delhi 110016

Recently, J.L. Journé has constructed a remarkable example of a bounded operator valued quantum martingale  $X$  in the usual filtration of the boson Fock space  $\mathfrak{F}_+(L_2(0, \infty))$  which does not admit the representation

$$dX = E dA^\dagger + F d\Lambda + G dA \quad (1)$$

in the sense of [1] over the domain  $\mathcal{E}$ , where  $E, F, G$  are adapted operator processes,  $A^\dagger, \Lambda, A$  are respectively the creation, conservation ( gauge ) and annihilation martingales, and  $\mathcal{E}$  is the linear manifold generated by exponential vectors. In [3], a necessary and sufficient regularity condition was found for a bounded operator valued martingale  $X$  to satisfy (1) with  $E, F, G$  being bounded operator valued processes satisfying the condition

$$\int_0^t (\|E(s)\|^2 + \|G(s)\|^2) ds < \infty \quad \text{for all } t .$$

The purpose of this note is to indicate the possibility of achieving the representation (1) for the Journé example provided  $\mathcal{E}$  is suitably restricted and the strictness of the definition of an adapted process is relaxed.

We begin with a class of examples of quantum martingales determined by second quantization of integral operators. Let  $K(\cdot, \cdot)$  be a complex valued continuous function on  $[0, \infty) \times [0, \infty)$ . For every  $t > 0$  define the bounded operator  $K_t$  on  $L_2([0, \infty))$  by

$$(K_t f)(s) = \chi_{[0, t]}(s) \int_0^t K(s, \tau) f(\tau) d\tau + \chi_{(t, \infty)}(s) f(s) \quad (2)$$

where  $\chi_C$  denotes the indicator function of the set  $C$ . Let  $K_\infty$  denote the integral operator defined by

$$(K_\infty f)(s) = \int_0^\infty K(s, \tau) f(\tau) d\tau . \quad (3)$$

If  $K_\infty$  is a bounded operator on  $L_2[0, \infty)$  then  $\|K_t\| \leq \max(1, \|K_\infty\|)$  for all  $t$ . Let  $X_K(t)$  denote the operator on Fock space defined by

the relations

$$X_K(t)\psi(f) = \psi(K_t f) \quad \text{for } f \in L_2[0, \infty) \quad (4)$$

where  $\psi(f)$  is the exponential vector corresponding to  $f$ . Then  $X_K(t)$ , the second quantization of  $K_t$ , is defined on the domain  $\mathcal{E}$  and an easy computation shows that

$$\langle \psi(fx_{[0,a]}), X_K(t)\psi(gx_{[0,a]}) \rangle = \langle \psi(fx_{[0,a]}), X_K(a)\psi(gx_{[0,a]}) \rangle$$

for all  $t \geq a$ . In other words,  $X_K = \{X_K(t), t \geq 0\}$  is a quantum martingale with domain  $\mathcal{E}$ , and  $X_K(0) = \text{identity}$ . If  $K_\infty$  is a contraction then  $K_t$  is a contraction for every  $t$  and hence  $X_K(t)$  can be extended to a contraction on Fock space. In other words,  $X_K$  becomes a contraction valued operator martingale.

By straightforward differentiation we obtain the relation

$$\frac{d}{dt} \langle \psi(f), X_K(t)\psi(g) \rangle = \langle \psi(f), X_K(t)\psi(g) \rangle \{-\bar{F}(t)g(t) + \bar{F}(t) \int_0^t K(t,s)g(s)ds + g(t) \int_0^t K(s,t)\bar{F}(s)ds\} \quad (5)$$

in the generalized sense of absolute continuity. In the language of [1], (5) is equivalent to saying that  $X_K$  obeys the quantum stochastic differential equation

$$dX_K = X_K L dA^\dagger - X_K dA + M X_K dA \quad (6)$$

where  $L$  and  $M$  are adapted processes of operators defined on the domain  $\mathcal{E}$  by the relations

$$\left. \begin{aligned} L(t) &= a(x_{[0,t]} \overline{K(t, \cdot)}) \\ M(t) &= a^\dagger(x_{[0,t]} \overline{K(\cdot, t)}) \end{aligned} \right\} \quad (7)$$

and  $a, a^\dagger$  are the usual annihilation and creation fields over  $L_2[0, \infty)$ . Since  $\mathcal{E}$  is left invariant by the operators  $X_K(t)$ ,  $a(h)$  for all  $t \geq 0$ ,  $h \in L_2[0, \infty)$ , it follows that the coefficients  $X_K L$ ,  $-X_K$ ,  $M X_K$  of  $dA^\dagger$ ,  $dA$  and  $dA$  respectively in (6) are all well defined adapted processes on the domain  $\mathcal{E}$  satisfying the inequalities

$$\int_0^t \{ \|X_K(s)L(s)\psi(f)\|^2 + |f(s)|^2 \|X_K(s)\psi(f)\|^2 + \|M(s)X_K(s)\psi(f)\|^2 \} ds < \infty \quad (8)$$

for all  $t$ . We remark that the finiteness of the third integral in (8) follows from the canonical commutation rules.

We now try to relax the conditions on the kernel  $K$ . As long as  $K_T$  is a bounded operator for each  $t$  and

$$\int_0^t (|K(s,t)|^2 + |K(t,s)|^2) ds < \infty \quad \text{for each } t \quad (9)$$

it is clear that  $L(t)$  and  $M(t)$  are well defined on  $\mathcal{E}$  by (7) and

condition (8) obtains. This implies (6). In the example of Journé,  $K(s,t) = (s-t)^{-1}$  and (9) breaks down. Then the definition of  $L(t)$  and  $M(t)$  by (7) does not make any sense. To face this situation we have to interpret equation (6) in a weak sense and we proceed as follows.

Let  $\mathcal{L} \subset L_2[0,\infty)$  denote the linear manifold of functions  $f$  which satisfy the local Lipschitz's condition

$$\|f\|_{\mathcal{L},t} = \sup_{0 \leq x,y \leq t} \left| \frac{f(x)-f(y)}{x-y} \right| < \infty \quad \text{for all } t.$$

Denote by  $\mathcal{E}(\mathcal{L})$  the linear manifold generated by the set  $\{\psi(f), f \in \mathcal{L}\}$ . For the kernel  $K(s,t) = (s-t)^{-1}$  we define the martingale  $X_K$  as before and observe that for  $f, g \in \mathcal{L}$

$$\begin{aligned} \frac{d}{dt} \langle \psi(f), X_K(t)\psi(g) \rangle &= \\ \langle \psi(f), X_K(t)\psi(g) \rangle &= \left\{ -\bar{F}(t)g(t) + \int_0^t \frac{\bar{F}(t)g(s) - \bar{F}(s)g(t)}{t-s} ds \right\} \end{aligned} \quad (10)$$

The last integral in the above equation can be written as

$$\bar{F}(t) \int_0^t \frac{g(s)-g(t)}{t-s} ds + g(t) \int_0^t \frac{\bar{F}(t)-\bar{F}(s)}{t-s} ds.$$

This suggests the introduction of the Schwartz-Fock space  $\mathcal{F}_S$  of sequences of distributions in the following sense. Any element of  $\mathcal{F}_S$  is of the form  $\lambda = (c, \lambda_1, \lambda_2, \dots)$  where  $\lambda_n$  is a symmetric distribution (in the sense of Schwartz) in the space  $\mathbb{R}^n$  for  $n=1,2,\dots$ , and  $c \in \mathbb{C}$ .

Let  $\mathcal{M}$  be an arbitrary dense linear manifold in  $L_2[0,\infty)$  and let  $\mathcal{E}(\mathcal{M})$  be the linear manifold in Fock space generated by  $\{\psi(f), f \in \mathcal{M}\}$ . Let  $E = \{E(t), t \geq 0\}$  be a family of linear maps  $E(t) : \mathcal{E}(\mathcal{M}) \rightarrow \mathcal{F}_S$  satisfying the following conditions:

1) For any  $C_c^\infty$  function  $f$  on  $[0,\infty)$ , the series

$$\langle \psi(f), E(t)\psi(g) \rangle = 1 + \sum_{n=1}^{\infty} n!^{-1/2} (E(t)\psi(g))_n (\bar{F}^{\otimes n}) \quad (11)$$

converges absolutely, where  $(E(t)\psi(g))_n$  denotes the  $n$ -th term of the sequence  $E(t)\psi(g)$  (a symmetric distribution on  $\mathbb{R}^n$ ).

2) The scalar quantity

$$\langle \psi(f), E(t)\psi(g) \rangle \exp\left(-\int_t^\infty \bar{F}(s)g(s)ds\right)$$

depends only on the values of  $f$  and  $g$  on the interval  $[0,t]$  and is a Borel function of  $t$ .

Then we say that  $E$  is a generalized adapted process with domain  $\mathcal{E}(\mathcal{M})$ . If  $X, E, F, G, H$  are five generalized adapted processes over  $\mathcal{E}(\mathcal{M})$  such that

$$\begin{aligned} \langle \psi(f), X(t_2)\psi(g) - X(t_1)\psi(g) \rangle &= \int_{t_1}^{t_2} \langle \psi(f), \bar{F}(s)E(s)\psi(g) + \\ &+ \bar{F}(s)g(s)F(s)\psi(g) + H(s)\psi(g) \rangle ds \quad \text{for all } f \in C_c^\infty, g \in \mathcal{M}, t_1 < t_2 \end{aligned}$$

we say that

$$dX = EdA^\dagger + Fd\Lambda + GdA + Hdt$$

on the domain  $\mathcal{E}(\mathcal{M})$ . With these conventions it is straightforward to verify that the Journé martingale  $X_K$  obeys the generalized quantum stochastic differential equation

$$dX_K = EdA^\dagger + Fd\Lambda + GdA$$

where E,F,G are generalized adapted processes over the domain  $\mathcal{E}(t)$ , defined by

$$E(t)\psi(g) = \left( \int_0^t \frac{g(s)-g(t)}{t-s} ds \right) \psi(K_t g),$$

$$F(t)\psi(g) = -\psi(K_t g),$$

$$G(t)\psi(g) = (1, \lambda_1, \lambda_2, \dots) e^{\mathcal{F}_S},$$

where for any test function  $\phi$  on  $\mathbb{R}^n$

$$\lambda_n(\phi) = n!^{-1/2} \sum_{j=1}^n \int_{\mathbb{R}^n} \chi_{[0,t]}(x_j) h(x_1) \dots \hat{h}(x_j) \dots h(x_n) \\ \times \frac{\phi(x_1, \dots, x_{j-1}, t, x_{j+1}, \dots, x_n) - \phi(x_1, \dots, x_n)}{t-x_j} dx_1 \dots dx_n,$$

here  $h=K_t g$ , and the symbol  $\hat{\phantom{h}}$  over a term implies its omission.

Just as the derivative of a bounded function on the line could be a distribution, it seems that the 'partial derivatives' of a bounded operator valued quantum martingale  $X$  in Fock space, with respect to the fundamental creation, conservation and annihilation martingales could be generalized adapted processes determining  $X$ .

I wish to thank P.A. Meyer for presenting me a detailed account of the example in [2], W. von Waldenfels and M. Schurmann for their patient hearing of several preliminary accounts of the contents of the present exposition.

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Note. The norm conventions on Fock space are those of [1], and are slightly different from those used elsewhere in this volume ( these would require a factor  $n!^{-1}$  instead of  $n!^{-1/2}$  in formula (11) ).