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ON THE WELLPOSED SINGULAR BOUNDARY VALUE

PROBLEMS FOR HEAT OPERATOR

by S. MIZOHATA

§ 1 - INTRODUCTION -

In [2], S. Itô treated the following initial-boundary value problem for heat operator :

$$\begin{array}{l}
 \text{(I.B.P.)} \\
 \left\{ \begin{array}{l}
 (1) \quad \frac{\partial u}{\partial t} = \Delta u, \text{ in } \Omega \subset \mathbb{R}^n, \\
 (2) \quad Bu = a(x) \frac{\partial u}{\partial n} + b(x)u = 0, \text{ } x \in \partial\Omega, \\
 \text{where } \frac{\partial u}{\partial n} \text{ is the derivative in the direction of outer-} \\
 \text{normal, and,} \\
 a(x) \geq 0, b(x) \geq 0, a(x) + b(x) = 1, \\
 (3) \quad u|_{t=0} = u_0(x)
 \end{array} \right.
 \end{array}$$

He proved that this problem is well-posed by constructing explicitly fundamental solution. His method is fairly complicated. Later several authors (cf. [1], [3], [4], [6]), treated the elliptic boundary value problem with the boundary condition (2), using the methods of functional analysis. Fairly recently, M. Terakado pointed out that in the paper of S. Itô, a crucial lemma is not clear. From that time, the author discussed often with him about this problem. Our purpose is to consider what conditions should be imposed on $a(x)$ and $b(x)$ in (2), in order that the problem (I.B.P.) to be wellposed.

To be more precise, we consider the problem under the following situation.

$$\left\{ \begin{array}{l}
 \Omega = \{(x,y) \in \mathbb{R}^2 \mid y < 0, a(x) \frac{\partial u}{\partial y} + b(x)u|_{y=0} = 0, \\
 a(x) \text{ and } b(x) \text{ are real-valued ; } a(x)^2 + b(x)^2 = 1 ; \\
 a(x), b(x) \in C^{\infty}, \text{ and } \underline{\text{bounded}} \text{ with all their derivatives.}
 \end{array} \right.$$

We are concerned with $H^\infty(\Omega)$ -wellposedness for (I.B.P.), imposing the compatibility conditions. Namely, denoting

$$(4) \quad Bu \equiv a(x) \frac{\partial u}{\partial y} + b(x)u \Big|_{y=0},$$

$$(5) \quad B(\Delta^j u_0(x,y)) = 0, \text{ for } j = 0,1,2,\dots$$

The result is the following.

Theorem. For the above problem to be H^∞ -wellposed, the following conditions are necessary and sufficient :

- (i) $a(x)$ does not change the sign. We assume therefore $a(x) \geq 0$,
- (ii) On the set $\{x; a(x) = 0\}$, $b(x) > 0$.

We are concerned here with the necessity. The sufficiency is proved fairly easily. First we observe that, if we put

$$u(.,t) = T_t u_0,$$

T_t is a semi-group. T. Komura obtained the necessary and sufficient condition to the infinitesimal generator in Fréchet spaces [5]. However it seems difficult to apply her method to the actual problem. Instead of that, suggested by this article, we use the truncated Laplace transform :

$$\hat{u}(.,\lambda) = \int_0^1 e^{-\lambda t} u(.,t) dt.$$

§ 2 - PRELIMINARIES -

We assume the problem (I.B.P.) is H^∞ -wellposed, and from this assumption some basic facts.

1)- Continuity :

By the assumption of the wellposedness, by Banach, there exist an integer q and the constant $C(T)$ such that

$$(2.1) \quad \sup_{t \in [0, T]} \|u(\cdot, t)\|_{H^2} \leq C(T) \|u(\cdot, 0)\|_{H^{2q}}.$$

In this inequality, the initial data should satisfy the compatibility conditions :

$$(2.2) \quad B(\Delta^j u(\cdot, 0)) = 0, \quad 0 \leq j \leq q-1.$$

2)- Truncated Laplace transform :

Let :

$$(2.3) \quad \hat{u}(\cdot, \lambda) = \int_0^1 e^{-\lambda t} u(\cdot, t) dt,$$

we obtain

$$(2.4) \quad (\lambda - \Delta) \hat{u}(\cdot, \lambda) = u(\cdot, 0) - e^{-\lambda} u(\cdot, 1).$$

We take $u(\cdot, 0)$ in the following form: Let $f_0(x) \in C_0^\infty$, and put

$$u_1(x, y) = \left(\sum_{j=0}^q \left(\lambda - \frac{d^2}{dx^2} \right)^j f_0(x) \cdot \frac{y^{2j}}{(2j)!} \right) \alpha(y),$$

where $\alpha(y) \in C_0^\infty$, which takes the value 1 in a neighborhood of the origin.

Then :

$$g(x, y) = (\lambda - \Delta) u_1 = \left(\lambda - \frac{d^2}{dx^2} \right)^{q+1} f_0(x) \cdot \frac{y^{2q}}{(2q)!},$$

in a neighborhood of the origin. g satisfies

$$B(\Delta^j g) = 0 \quad (0 \leq j \leq q-1), \quad \text{and}$$

$$(2.5) \quad B u_1 = -b(x) f_0(x).$$

Observe that g is determined by $f_0(x)$. We take

$$f_0(x) \in H^{4q+2}.$$

We put $u(\cdot, 0) = g(x, y)$. Then by hypothesis, there exists a unique

solution $u(.,t)$ with initial data g . Thus $\hat{u}(.,\lambda)$ is defined by (2.3). Since $(\lambda-\Delta)u_1 = u(.,0)$, from (2.4) we get

$$(\lambda-\Delta) (\hat{u}(.,\lambda) - u_1) = - e^{-\lambda} u(.,1).$$

Moreover, since $B\hat{u}(.,\lambda) = 0$, we obtain

$$B(\hat{u}-u_1) = - Bu_1 = - b(x) f_0(x).$$

We proceed further. Let $u_D(x,y)$ be the solution of

$$\begin{cases} (\lambda-\Delta)u_D = u(.,1) \\ u_D|_{y=0} = 0. \end{cases}$$

Then denoting

$$(2.6) \quad v(x,y) = \hat{u}(.,\lambda) - u_1(x,y) - e^{-\lambda}u_D$$

$$(2.7) \quad \begin{cases} (\lambda-\Delta) v(x,y) = 0 \\ B v = b(x) f_0(x) + e^{-\lambda} \tilde{f}(x) , \end{cases}$$

$$\text{where } \tilde{f}(x) = a(x) \frac{\partial}{\partial y} u_D|_{y=0} .$$

3) - Functional equation on the boundary :

If we use the Poisson representation of $v(x,y)$,

$$v(x,y) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ix\xi} e^{\sqrt{\lambda+\xi^2} y} \hat{v}(\xi,0) d\xi ,$$

putting $v(x,0) = v(x)$, the above equation can be written

$$(2.8) \quad A(x,D) v(x) = (a(x) \sqrt{\lambda+D^2} + b(x)) v(x) = f(x) + e^{-\lambda} \tilde{f}(x)$$

$$\text{where } f(x) = b(x) f_0(x).$$

We assume hereafter $b(x) \neq 0$, in a neighborhood of the origin, and $f_0(x)$ has its support in a small neighborhood of the origin. We obtain :

Proposition 1 :

$$(2.9) \quad \tilde{\Delta}(x,D) = a(x) + b(x) \sqrt{\lambda + D^2}^{-1} \equiv a(x) + b(x) \Lambda^{-1}.$$

Then for any $f(x) \in H^{4q+2}$, there exists a solution $w(x) \in H^{\frac{1}{2}}$

$$(2.10) \quad \tilde{\Delta}(x,D) w(x) = f(x) + e^{-\lambda} \tilde{f}(x),$$

satisfying the following type inequalities.

$$\|w(x)\|_{\frac{1}{2}}, \|\tilde{f}(x)\|_{\frac{3}{2}} \leq \text{const.} \|f(x)\|_{4q+2}.$$

Remark :

Hereafter we denote $\|\cdot\|_s = \|\cdot\|_{H^s}$, and

$$\|f(x)\|_s^2 = \int (\lambda + \xi^2)^s |\hat{f}(\xi)|^2 d\xi.$$

§ 3 - ALMOST NULL SOLUTIONS -

For simplicity, we consider the problem under the following assumptions :

$$(3.1) \quad \left\{ \begin{array}{l} \text{(i)} \quad a(x) = x^2 \text{ in a neighborhood } V \text{ of the origin,} \\ \text{(ii)} \quad b(x) = -1, \quad x \in \mathbb{R}. \end{array} \right.$$

We consider a solution $\psi_0(x)$ of the following equation

$$(3.2) \quad (x^2 - \Lambda^{-1}) \psi_0(x) = 0.$$

By taking the Fourier transformation,

$$(3.3) \quad \left(\frac{d^2}{d\xi^2} + (\lambda + \xi^2)^{-\frac{1}{2}} \right) \hat{\psi}_0(\xi) = 0.$$

We see that, there exists the solution satisfying

$$(3.4) \quad \begin{cases} |\psi_0(\xi)| \sim \lambda^{-\frac{1}{8}} (\lambda + \xi^2)^{-\frac{1}{8}}, (\lambda \rightarrow \infty) \\ |(\frac{d}{d\xi})^p \psi_0(\xi)| \leq C_p (\lambda + \xi^2)^{-\frac{1}{8} - \frac{1}{8}p} \lambda^{-\frac{1}{8}} \end{cases}$$

Observe that :

$$(3.5) \quad \begin{cases} \psi_0(x) \in H^{-\delta} \quad (\delta > \frac{3}{4}) \\ \|\Lambda^{-2} \psi_0\|_1 \sim C_1 \lambda^{-\frac{1}{4}} \quad (C_1 : \text{positive constant}) \end{cases}$$

Take a $\beta(x) \in C_0^\infty(V)$, which takes the value 1 in a neighborhood of the origin. From (3.2), it follows

$$(a(x) - \Lambda^{-1})(\beta\psi_0) - [\beta, \Lambda^{-1}]\psi_0 = 0.$$

Let us introduce the symbol " ≈ 0 ". $f \approx 0$ means that, for any k and any p ,

$$\|\Lambda^k f\|_0 \leq C_{kp} / \lambda^p \text{ when } \lambda \text{ is large.}$$

From the property (ii) of (3.4), it follows

$$[\beta, \Lambda^{-1}]\psi_0 \approx 0, [a(x) - \Lambda^{-1}](1 - \beta(x))\psi_0 \approx 0.$$

Therefore, it holds

$$(3.6) \quad (a(x) - \Lambda^{-1})\psi_0(x) \approx 0.$$

Following the notation (2.9), we denote $\tilde{A} = a(x) - \Lambda^{-1}$. Using the results of proposition 1, we get.

Proposition 2 :

For $f (= \Lambda^{-2}\psi_0) \in H^1$, there exists an element $w \in H^1$, satisfying

$$(i) \quad \|\tilde{A}w - f\|_1 \leq \frac{1}{2} \|f\|_1,$$

(ii) There exist positive constant k_0 and C such that

$$\|w\|_1 \leq C \lambda^{k_0} \quad (\text{if } \lambda \text{ is large}).$$

Now it is easy to see that these properties are not compatible with (3.6), which shows, in the case (3.1). The problem (I.B.P.) is not well-posed.

In fact, in view of (3.5) :

$$(3.7) \quad \operatorname{Re}(\tilde{A}w, \Lambda^{-2}\psi_0)_1 \geq \frac{1}{2} \|\Lambda^{-2}\psi_0\|_1^2 \sim \frac{1}{2} C_1^2 \lambda^{-\frac{1}{2}}$$

On the other hand, the left-hand side is equal to

$$\operatorname{Re} \langle \tilde{A}w, \overline{\psi_0} \rangle = \operatorname{Re} \langle w, \overline{\tilde{A}\psi_0} \rangle.$$

This is estimated from above by $\|w\|_1 \|\tilde{A}\psi_0\|_{-1}$. By the property of w , and (3.6), we see that, this last quantity is estimated of the form $C'_p \lambda^{-p}$ for any p . This is not compatible with (3.7).

§ 4 - PROOF OF PROPOSITION 2 -

We define \tilde{f}, w_0, w in the following order

$$f \rightarrow \tilde{f} \rightarrow w_0 \rightarrow w$$

$$(4.1) \quad \left\{ \begin{array}{l} (i) \quad \tilde{f} = \rho_\delta * f \\ (ii) \quad \tilde{A} w_0 = \tilde{f} + e^{-\lambda} f_1 \\ (iii) \quad w = \rho_\varepsilon * w_0 \end{array} \right.$$

where $\rho_\delta, \rho_\varepsilon$ are mollifiers.

$$(i) \quad \|f - \tilde{f}\|_1 \leq C \delta^s \|f\|_{1+s} \quad (0 < s < \frac{1}{4}). \quad \text{We suppose here } \delta \leq \lambda^{-\frac{1}{2}};$$

C is independent of δ . Since $\|f\|_{1+s} \sim C \lambda^{\frac{s}{2}} \|f\|_1$, we take $\delta = \varepsilon_0 \lambda^{-\frac{1}{2}}$ (ε_0 being a small constant). This, we have

$$(4.2) \quad \|\tilde{f} - \tilde{f}^\vee\|_1 \leq \frac{1}{6} \|f\|_1.$$

(ii) Observe $\tilde{f} \in H^\infty$, and

$$\|\tilde{f}^\vee\|_{4q+2} \leq \text{const.} \delta^{-(4q+2)} \|f\|_1 \leq \text{const.} \lambda^{2q+1} \|f\|_1,$$

(since $\delta = \varepsilon_0 \lambda^{-\frac{1}{2}}$).

By proposition 1, there exists $w_0 \in H^{\frac{1}{2}}$ such that

$$\tilde{A} w_0 = \tilde{f} + e^{-\lambda} f_1,$$

where

$$\|w_0\|_{\frac{1}{2}} \leq \text{const.} \|\tilde{f}^\vee\|_{4q+2} \leq \text{const.} \lambda^{2q+1} \|f\|_1,$$

$$\|f_1\|_{\frac{s}{2}} \leq \text{const.} \|\tilde{f}^\vee\|_{4q+2} \leq \text{const.} \lambda^{2q+1} \|f\|_1.$$

Observe that $\|\tilde{A} w_0 - \tilde{f}\|_1 = e^{-\lambda} \|f_1\|_1 \leq \text{const.} e^{-\lambda} \lambda^{2q+1} \|f\|_1$, which is negligible.

$$(iii) \quad f - \tilde{A} w = (f - \tilde{f}) + (\tilde{f} - \tilde{A} w_0) + (\tilde{A} w_0 - \tilde{A} w).$$

We are concerned with the last term. Observe $\tilde{A} w_0 \in H^{\frac{3}{2}}$, and

$$(4.3) \quad \|\tilde{A} w_0\|_{\frac{3}{2}} \leq \|\tilde{f}^\vee\|_{\frac{3}{2}} + e^{-\lambda} \|f_1\|_{\frac{3}{2}},$$

where the last term is negligible when $\lambda \rightarrow \infty$. Now,

$$\|\tilde{f}^\vee\|_{\frac{3}{2}} \leq \text{const.} \frac{1}{\sqrt{\delta}} \|f\|_1 \leq \text{const.} \lambda^{\frac{1}{4}} \|f\|_1.$$

$$\tilde{A} w_0 - \tilde{A}(\rho_\varepsilon * w_0) = (\tilde{A} w_0 - \rho_\varepsilon * \tilde{A} w_0) + (\rho_\varepsilon * \tilde{A} w_0 - \tilde{A}(\rho_\varepsilon * w_0)).$$

The first term is estimated by

$$\|\tilde{A} w_0 - \rho_\varepsilon * (\tilde{A} w_0)\|_1 \leq \text{const.} \|\tilde{A} w_0\|_{\frac{3}{2}} \sqrt{\varepsilon}.$$

By virtue of (4.5), this is estimated, when λ is large, by

$$\text{const.} (\lambda^{\frac{1}{4}} \sqrt{\varepsilon}) \|f\|_1.$$

Now put, $\varepsilon = \varepsilon_1' \lambda^{-(4q+2)}$, thus

$$\|[\rho_\varepsilon * , \tilde{A}] w_0\|_1 \leq \frac{1}{6} \|f\|_1,$$

which implies

$$\|\tilde{A} w_0 - \tilde{A} w\|_1 \leq (\frac{1}{6} + \text{const.} \lambda^{-2q}) \|f\|_1,$$

which completes the proof of proposition 2.

§ 5 - FINAL REMARKS -

In the case when $a(x)$ has a finite order of zero at the origin, we can argue essentially in the same way. However the technical complication arises when $a(x)$, and $b(x)$ are general. In the case when $a(x)$ has a infinite order of zero at the origin, we can argue in a fairly different way.

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