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KAZUNAGA TANAKA

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## Non-collision solutions for a second order singular Hamiltonian system with weak force

by

**Kazunaga TANAKA**

Department of Mathematics,  
School of Science, Nagoya University,  
Chikusa-ku, Nagoya 464, Japan

ABSTRACT. — Under a weak force type condition, we consider the existence of time periodic solutions of singular Hamiltonian systems:

$$\left. \begin{aligned} \ddot{q} + V_q(q, t) &= 0, \\ q(t+T) &= q(t). \end{aligned} \right\} \quad \text{(HS)}$$

We assume  $V(q, t) < 0$  for all  $q, t$  and  $V(q, t), V_q(q, t) \rightarrow 0$  as  $|q| \rightarrow \infty$ . Moreover we assume  $V(q, t)$  is of a form:

$$V(q, t) = -\frac{1}{|q|^\alpha} + U(q, t)$$

where  $0 < \alpha < 2$  and  $U(q, t) \in C^2((\mathbf{R}^N \setminus \{0\}) \times \mathbf{R}, \mathbf{R})$  is a  $T$ -periodic function in  $t$  such that  $|q|^\alpha U(q, t), |q|^{\alpha+1} U_q(q, t), |q|^{\alpha+2} U_{qq}(q, t), |q|^\alpha U_t(q, t) \rightarrow 0$  as  $|q| \rightarrow 0$ .

For  $\alpha \in (1, 2]$ , we prove the existence of a non-collision solution of (HS). For  $\alpha \in (0, 1]$ , we prove that the generalized solution of (HS), which is introduced in [BR], enters the singularity 0 at most one time in its period. Our argument depends on a minimax argument due to [BR] and an estimate of Morse index of corresponding functional, which will be obtained via re-scaling argument.

*Key words* : Periodic solutions, Hamiltonian systems, singular potentials, minimax methods, morse index.

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*Classification A.M.S.* : 58 E 05, 58 F 05.

RÉSUMÉ. — Sous une hypothèse de type force faible, nous étudions l'existence de solutions périodiques de systèmes hamiltoniens singuliers :

$$\left. \begin{aligned} \ddot{q} + V_q(q, t) &= 0, \\ q(t+T) &= q(t). \end{aligned} \right\} \quad (\text{HS})$$

Nous supposons que  $V(q, t) < 0$  pour tout  $q, t$  et que  $V(q, t), V_q(q, t) \rightarrow 0$  si  $|q| \rightarrow \infty$ .

De plus nous supposons que  $V$  est de la forme :

$$V(q, t) = -\frac{1}{|q|^\alpha} + U(q, t)$$

où  $0 < \alpha < 2$  et  $U(q, t) \in C^2((\mathbf{R}^N \setminus \{0\}) \times \mathbf{R}, \mathbf{R})$  est une fonction  $T$ -périodique telle que  $|q|^\alpha U(q, t), |q|^{\alpha+1} U_q(q, t), |q|^{\alpha+2} U_{qq}(q, t), |q|^\alpha U_t(q, t) \rightarrow 0$  as  $|q| \rightarrow 0$ .

Pour  $\alpha \in ]1, 2]$ , nous démontrons l'existence d'une solution non collisionnelle de (HS).

Pour  $\alpha \in ]0, 1]$ , nous démontrons que la solution généralisée de (HS), introduite dans [BR], passe par la singularité 0 au plus une fois dans la période. Notre démonstration utilise un argument de minimax dû à [BR] et une estimation de l'indice de Morse de la fonctionnelle correspondante, obtenu par un argument de changement d'échelles.

## 0. INTRODUCTION

We study the existence of  $T$ -periodic solutions of the following Hamiltonian system:

$$\left. \begin{aligned} \ddot{q} + V_q(q, t) &= 0, \\ q(t+T) &= q(t), \quad t \in \mathbf{R}. \\ q(t) &\neq 0, \end{aligned} \right\} \quad (\text{HS})$$

where  $q = (q_1, q_2, \dots, q_N) \in \mathbf{R}^N$  ( $N \geq 3$ ) and  $V(q, t) : (\mathbf{R}^N \setminus \{0\}) \times \mathbf{R} \rightarrow \mathbf{R}$  is a  $T$ -periodic (in  $t$ ) function such that  $V(q, t), V_q(q, t) \rightarrow 0$  as  $|q| \rightarrow \infty$  and  $V(q, t) \rightarrow -\infty$  as  $q \rightarrow 0$ .

Classical solutions of (HS) can be characterized as critical points of functional:

$$I(q) = \int_0^T \left[ \frac{1}{2} |\dot{q}|^2 - V(q, t) \right] dt: \Lambda \rightarrow \mathbf{R}$$

where

$$\Lambda = \{ q(t) \in H_{\text{loc}}^1(\mathbf{R}, \mathbf{R}^N); q(t+T) = q(t), q(t) \neq 0 \text{ for all } t \}.$$

In case  $V(q, t)$  satisfies the *strong force condition* (SF) of Gordon [Go]:

(SF) there is a neighborhood  $\Omega$  of 0 in  $\mathbf{R}^N$  and a function  $W(q) \in C^1(\Omega \setminus \{0\}, \mathbf{R})$  such that

$$\begin{aligned} W(q) &\rightarrow \infty \text{ as } q \rightarrow 0, \\ -V(q, t) &\geq |W_q(q)|^2 \text{ for all } q \in \Omega \setminus \{0\} \text{ and } t, \end{aligned}$$

the functional  $I(q)$  satisfies the Palais-Smale compactness condition and we can apply minimax arguments to  $I(q)$ . Especially under the assumptions of (SF) and

(V1)  $V(q, t) \in C^1((\mathbf{R}^N \setminus \{0\}) \times \mathbf{R}, \mathbf{R})$  is  $T$ -periodic in  $t$ ;

(V2)  $V(q, t) < 0$  and  $V(q, t), V_q(q, t) \rightarrow 0$  as  $|q| \rightarrow \infty$ ;

(V3)  $-V(q, t) \rightarrow \infty$  as  $q \rightarrow 0$ ,

Bahri and Rabinowitz [BR] introduced a minimax method and obtained the existence of classical solutions (non-collision solutions) of (HS). See also [AC1, Gr1]. But in case (SF) does not hold, we cannot verify the Palais-Smale compactness condition for  $I(q)$  and we cannot apply minimax argument directly to  $I(q)$ . However, using a suitable approximation argument, Bahri and Rabinowitz [BR] proved the existence of *generalized  $T$ -periodic solutions*, that may enter the *singularity 0* (i.e., *collision*) under the conditions (V1)-(V3) (without (SF)).

For the study of the existence of non-collision solutions in case of *weak forces* (i.e., the case where (SF) does not hold), we refer to [AC3], [DGM], [DG], [C], [ST]. In [AC3], [DGM], [DG], they found critical points of  $I(q)$ , whose critical values are less than

$$\inf_{q(t) \in \partial \Lambda} I(q) = \inf \{ I(q); q \in H_{\text{loc}}^1(\mathbf{R}, \mathbf{R}^N), q(t+T) = q(t) \text{ for all } t \text{ and } q(t) = 0 \text{ for some } t \}.$$

In [C], [ST], they studied (HS) through minimization problems. They studied the behavior of solutions near collisions (especially [ST] studied the Morse index) and they obtained the existence of non-collision solutions.

This work is largely motivated by the works [BR], [C], [ST] and we study the existence of non-collision solutions under the weak force condition through *minimax problem*. We study the following class of weak force potentials; for  $0 < \alpha < 2$  we assume the potential  $V(q, t)$  is of a form:

$$(W1) \quad V(q, t) = -\frac{1}{|q|^\alpha} + U(q, t);$$

where

(W2)  $U(q, t) \in C^2((\mathbf{R}^N \setminus \{0\}) \times \mathbf{R}, \mathbf{R})$  is  $T$ -periodic in  $t$ ;

(W3)  $|q|^\alpha U(q, t), |q|^{\alpha+1} U_q(q, t), |q|^{\alpha+2} U_{qq}(q, t), |q|^\alpha U_t(q, t) \rightarrow 0$  as  $|q| \rightarrow 0$  uniformly in  $t$ .

We remark (V1) and (V3) follow from (W1)-(W3). We also remark (SF) holds if  $\alpha \geq 2$ .

Our main result is as follows:

**THEOREM 0.1.** — *Assume  $N \geq 3$ , (W1)-(W3), (V2) and  $1 < \alpha < 2$ . Then (HS) has at least one T-periodic (non-collision) solution.*

In case  $0 < \alpha \leq 1$ , we cannot show the existence of non-collision solution. However we can estimate the number of collisions of the generalized T-periodic solutions due to Bahri and Rabinowitz [BR]. More precisely, we get

**THEOREM 0.2.** — *Assume  $N \geq 3$ , (W1)-(W3), (V2) and  $0 < \alpha \leq 1$ . Then (HS) has a generalized T-periodic solution which has at most one collision, i. e., which enters the singularity 0 at most one time in its period T.*

The existence of a non-collision solution of (HS) will be obtained as follows; first we consider modified functional:

$$I_\varepsilon(q) = \int_0^T \left[ \frac{1}{2} |\dot{q}|^2 - V(q, t) + \frac{\varepsilon}{|q|^4} \right] dt \quad \text{for } \varepsilon \in (0, 1]$$

and obtain critical points  $q_\varepsilon \in \Lambda$  of  $I_\varepsilon(q)$ . Second, we try to pass to the limit  $\varepsilon \rightarrow 0$ . Here we remark  $I_\varepsilon(q)$  satisfies the strong force condition (SF) for each  $\varepsilon \in (0, 1]$ .

The proof of Theorem 0.1 will be given in the following sections; in Section 1, we study the modified functional  $I_\varepsilon(q)$ . We apply the minimax method of Bahri and Rabinowitz [BR] and get a critical point  $q_\varepsilon(t)$  of  $I_\varepsilon(q)$  for  $\varepsilon \in (0, 1]$ . Moreover we obtain the following uniform bounds

$$m \leq I_\varepsilon(q_\varepsilon) \leq M, \tag{0.1}$$

$$I'_\varepsilon(q_\varepsilon) = 0, \tag{0.2}$$

$$\text{index } I''_\varepsilon(q_\varepsilon) \leq N - 2, \tag{0.3}$$

for  $\varepsilon \in (0, 1]$ , where  $m, M > 0$  are independent of  $\varepsilon$ . Here we denote by  $\text{index } I''_\varepsilon(q_\varepsilon)$ , the Morse index of  $I''_\varepsilon(q_\varepsilon)$ .

From (0.1) and (0.2), we can deduce the uniform  $H^1$ -bound for  $(q_\varepsilon(t))_{\varepsilon \in (0, 1]}$ . Thus we may assume

$$q_{\varepsilon_n} \rightarrow q_\infty \quad \text{weakly in } H^1 \text{ and strongly in } L^\infty \tag{0.4}$$

for some sequence  $\varepsilon_n \rightarrow 0$ . However  $q_\infty(t)$  may enter the singularity 0.

In Sections 2-4, we study the behavior of critical points  $(q_{\varepsilon_n})_{n=1}^\infty$  of  $I_{\varepsilon_n}(q)$  with properties (0.1), (0.2) and (0.4). We will establish the following estimate of the Morse index

**PROPOSITION 0.3.** — *Let  $(q_n(t))_{n=1}^\infty \subset \Lambda$  be a sequence of critical points of  $I_{\varepsilon_n}(q)$  satisfying*

- (i)  $\varepsilon_n \rightarrow 0$ ;

(ii) there are constants  $0 < m < M$  independent of  $n$  such that

$$I_{\varepsilon_n}(q_n) \in [m, M] \quad \text{for all } n;$$

(iii)  $I'_{\varepsilon_n}(q_n) = 0$ ;

(iv)  $q_n \rightarrow q_\infty(t)$  weakly in  $H^1$  and strongly in  $L^\infty$ ;

and let  $v$  be the number of times  $q_\infty(t)$  enters the singularity 0; that is,

$$v = \#\{t \in (0, T]; q_\infty(t) = 0\} \in \mathbf{N} \cup \{\infty\}. \quad (0.5)$$

Then

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) \geq (N-2) i(\alpha) v, \quad (0.6)$$

where  $i(\alpha) \in \mathbf{N}$  is an integer defined by

$$i(\alpha) = \max \left\{ m \in \mathbf{N}; m < \frac{2}{2-\alpha} \right\}. \quad (0.7)$$

We remark that  $i(\alpha) = 1$  for  $\alpha \in (0, 1]$  and  $i(\alpha) \geq 2$  for  $\alpha \in (1, 2)$ . To prove the above proposition, we use re-scaling argument, which is based on the scale-invariance of the equation:

$$\ddot{q} + \frac{\alpha q}{|q|^{\alpha+2}} = 0 \quad \text{in } \mathbf{R}, \quad (0.8)$$

that is, (0.8) is invariant by the scale changes:

$$q(\cdot) \rightarrow \delta^{-1} q(\delta^{(\alpha+2)/2} \cdot).$$

In Section 5, we combine results obtained in Sections 1-4 and give proofs of our theorems 0.1 and 0.2.

## 1. MODIFIED FUNCTIONAL AND MINIMAX PROCEDURE

In this section, we study the following functional

$$I_\varepsilon(q) = \int_0^T \left[ \frac{1}{2} |\dot{q}|^2 - V(q, t) + \frac{\varepsilon}{|q|^4} \right] dt \quad \text{for } \varepsilon \in (0, 1]. \quad (1.1)$$

Here we assume only (V2), (V3) and

(V1')  $V(q, t) \in C^2((\mathbf{R}^N \setminus \{0\}) \times \mathbf{R}, \mathbf{R})$  is  $T$ -periodic in  $t$ .

We need the following notations; let  $E = H_1^1(\mathbf{R}, \mathbf{R}^N)$  denote the space of  $T$ -periodic functions on  $\mathbf{R}$  with values in  $\mathbf{R}^N$  under the norm:

$$\|q\|_E = \left( \int_0^T |\dot{q}|^2 dt + [q]^2 \right)^{1/2},$$

where  $[q] = \frac{1}{T} \int_0^T q(t) dt$ . We remark that

$$\Lambda = \{q \in E; q(t) \neq 0 \text{ for all } t\}$$

is open in  $E$  and  $I_\varepsilon(q) \in C^2(\Lambda, \mathbf{R})$ . We also use the notation:

$$\|q\|_{L^2} = \left( \int_0^T |q(t)|^2 dt \right)^{1/2}.$$

There is a one-to-one correspondence between critical points of  $I_\varepsilon(q)$  and classical  $T$ -periodic solutions of the following equation:

$$\left. \begin{aligned} \ddot{q} + V_q(q, t) + \frac{4\varepsilon q}{|q|^6} &= 0, \\ q(t+T) &= q(t), \\ q(t) &\neq 0. \end{aligned} \right\} \text{ in } \mathbf{R}, \quad (1.2)$$

We remark the potential  $V(q, t) - \frac{\varepsilon}{|q|^4}$  satisfies the strong force condition

(SF) with  $W(q) = \frac{\sqrt{\varepsilon}}{|q|}$ .

First we state some properties of  $I_\varepsilon(q)$ .

LEMMA 1.1. — Assume (V1'), (V2) and (V3).

(i) For any  $M > 0$ , there exist constants  $C_i(M) > 0$  ( $i = 1, 2$ ) independent of  $\varepsilon \in (0, 1]$  such that

$$\|\dot{q}\|_{L^2}, \quad \int_0^T -V(q, t) dt, \quad \int_0^T \frac{\varepsilon}{|q|^4} dt \leq C_1(M), \quad (1.3)$$

$$\min_{t \in [0, T]} |q(t)| \geq C_2(M) \varepsilon^{1/2} \quad (1.4)$$

for all  $q \in \Lambda$  and  $\varepsilon \in (0, 1]$  with  $I_\varepsilon(q) \leq M$ .

(ii) For any  $M > m > 0$ , there exists a constant  $C_3(m, M) > 0$  independent of  $\varepsilon \in (0, 1]$  such that

$$\|q\|_E \leq C_3(m, M) \quad (1.5)$$

for all  $q \in \Lambda$  and  $\varepsilon \in (0, 1]$  with  $I_\varepsilon(q) \in [m, M]$  and  $\|I'_\varepsilon(q)\|_{E^*} \leq m/\sqrt{2M}$ .

(iii) For any  $\varepsilon \in (0, 1]$ ,  $I_\varepsilon(q)$  satisfies the condition  $(PS^+)$  on  $\Lambda$ :

$(PS^+)$ : for any  $s > 0$ , if  $(q_n) \subset \Lambda$ ,  $I_\varepsilon(q_n) \rightarrow s$  and  $I'_\varepsilon(q_n) \rightarrow 0$ , then  $q_n$  possesses a subsequence converging to some  $q \in \Lambda$  in  $E$ .

*Proof.* — (i) By (V2) and (V3), it follows from  $I_\varepsilon(q) \leq M$  that

$$\|\dot{q}\|_{L^2} \leq \sqrt{2M}, \quad (1.6)$$

$$\int_0^T -V(q, t) dt \leq M, \quad (1.7)$$

$$\int_0^T \frac{\varepsilon}{|q|^4} dt \leq M. \quad (1.8)$$

Thus we get (1.3). Next we deal with (1.4). We get for all  $s, t \in [0, T]$  that

$$\begin{aligned} \frac{1}{|q(t)|} - \frac{1}{|q(s)|} &\leq \int_0^T \left| \frac{d}{d\tau} \frac{1}{|q(\tau)|} \right| d\tau \\ &\leq \left( \int_0^T |\dot{q}|^2 d\tau \right)^{1/2} \left( \int_0^T \frac{1}{|q(\tau)|^4} d\tau \right)^{1/2} \leq \frac{\sqrt{2}M}{\sqrt{\varepsilon}}. \end{aligned} \quad (1.9)$$

By (V3), we can find a constant  $c(M) > 0$  with the following property; for any  $q \in \Lambda$  with (1.7) there is a  $t_0 = t_0(q) \in [0, T]$  such that

$$|q(t_0)| \geq c(M).$$

We set  $s = t_0$  in (1.9), then we get for all  $t \in [0, T]$

$$\frac{1}{|q(t)|} \leq \frac{\sqrt{2}M}{\sqrt{\varepsilon}} + \frac{1}{c(M)} \leq \frac{1}{\sqrt{\varepsilon}} \left( \sqrt{2}M + \frac{1}{c(M)} \right).$$

Thus

$$|q(t)| \geq \left( \sqrt{2}M + \frac{1}{c(M)} \right)^{-1} \varepsilon^{1/2} \equiv C_2(M) \varepsilon^{1/2}.$$

Hence we get (1.4).

(ii) By (1.6), it suffices to prove  $\|q\|_{L^\infty} \leq C_3(m, M)$ . We have for  $q \in \Lambda$  with  $\|I'_\varepsilon(q)\|_{E^*} \leq m/\sqrt{2M}$  and  $I_\varepsilon(q) \leq M$  that

$$\begin{aligned} I_\varepsilon(q) &= \frac{1}{2} I'_\varepsilon(q)(q - [q]) + \frac{1}{2} \int_0^T V_q(q, t)(q - [q]) dt \\ &\quad + \int_0^T -V(q, t) dt + 2 \int_0^T \frac{\varepsilon}{|q|^6} (q, q - [q]) dt + \int_0^T \frac{\varepsilon}{|q|^4} dt \\ &\leq \frac{1}{2} \frac{m}{\sqrt{2M}} \|\dot{q}\|_{L^2} + \frac{1}{2} \int_0^T |V_q(q, t)| |q - [q]| dt \\ &\quad + \int_0^T -V(q, t) dt + \int_0^T \frac{2|q - [q]|}{|q|^5} dt + \int_0^T \frac{1}{|q|^4} dt. \end{aligned}$$

Note that we have from (1.6)

$$\|q(t) - [q]\|_{L^\infty} \leq \sqrt{T} \|\dot{q}\|_{L^2} \leq \sqrt{2TM}.$$

Thus we get

$$\begin{aligned}
 I_\varepsilon(q) &\leq \frac{1}{2} \frac{m}{\sqrt{2M}} \sqrt{2M} + \frac{\sqrt{2TM}}{2} \int_0^T |V_q(q, t)| dt \\
 &\quad + \int_0^T -V(q, t) dt + \int_0^T \frac{2}{|q|^5} \sqrt{2TM} dt + \int_0^T \frac{1}{|q|^4} dt \\
 &\leq \frac{1}{2} m + T \Phi([q] - \sqrt{2TM}), \quad (1.10)
 \end{aligned}$$

where

$$\Phi(R) = \max_{|y| \geq R, t \in [0, T]} \left[ \frac{\sqrt{2TM}}{2} |V_q(y, t)| - V(y, t) \right] + \frac{2\sqrt{2TM}}{R^5} + \frac{1}{R^4}.$$

We remark that

$$\Phi(R) \rightarrow 0 \quad \text{as } R \rightarrow \infty. \quad (1.11)$$

Now we assume  $I_\varepsilon(q) \in [m, M]$ , then we have from (1.10) that

$$\frac{1}{2} m \leq T \Phi([q] - \sqrt{2TM}).$$

By (1.11), we can see there is a constant  $C_3(m, M) > 0$  independent of  $\varepsilon \in (0, 1]$  such that

$$|[q]| \leq C_3(m, M),$$

*i. e.*,

$$\|q\|_{L^\infty} \leq C_3(m, M).$$

Thus we get (1.5).

(iii) Assume  $(q_n) \subset \Lambda$  satisfies  $I_\varepsilon(q_n) \rightarrow s > 0$  and  $I'_\varepsilon(q_n) \rightarrow 0$  in  $E^*$ . From (1.4)-(1.6), we can extract a subsequence – we denote it still by  $q_n$  – such that

$$q_n \rightarrow q \in \Lambda \text{ weakly in } E \text{ and strongly in } L^\infty.$$

Thus the form of  $I'_\varepsilon(q)$  shows  $q_n \rightarrow q$  strongly in  $E$ . ■

Next we apply minimax method, which is essentially due to Bahri and Rabinowitz [BR], to  $I_\varepsilon(q)$  for each  $\varepsilon \in (0, 1]$ . Consider the family of mappings  $C(S^{N-2}, \Lambda)$ . Identifying  $[0, T]/\{0, T\} \simeq S^1$ , we can associate each  $\gamma \in C(S^{N-2}, \Lambda)$  with a mapping  $\tilde{\gamma}: S^{N-2} \times S^1 \rightarrow S^{N-1}$  by

$$\tilde{\gamma}(x, t) = \frac{\gamma(x)(t)}{|\gamma(x)(t)|} \quad \text{for } x \in S^{N-2} \quad \text{and} \quad t \in S^1 \simeq [0, T]/\{0, T\}.$$

We denote the Brouwer degree of  $\tilde{\gamma}$  by  $\text{deg } \tilde{\gamma}$ . We define

$$\Gamma^* = \{ \gamma \in C(S^{N-2}, \Lambda); \text{deg } \tilde{\gamma} \neq 0 \}. \quad (1.12)$$

We can see  $\Gamma^* \neq \emptyset$  as in [BR], Lemma 1.2.

We define minimax values of  $I_\varepsilon(q)$  as follows:

$$b_\varepsilon = \inf_{\gamma \in \Gamma^*} \max_{x \in S^{N-2}} I_\varepsilon(\gamma(x)) \quad \text{for } \varepsilon \in (0, 1], \quad (1.13)$$

$$b_0 = \inf_{\gamma \in \Gamma^*} \max_{x \in S^{N-2}} I(\gamma(x)). \quad (1.14)$$

Since  $I(q) \leq I_\varepsilon(q) \leq I_1(q)$  for all  $q \in \Lambda$  and  $\varepsilon \in (0, 1]$ , we have

$$b_0 \leq b_\varepsilon \leq b_1 \quad \text{for } \varepsilon \in (0, 1]. \quad (1.15)$$

We argue as in [BR], Proposition 1.4, we get

PROPOSITION 1.2. —  $b_0 > 0$ . ■

Thus we have

PROPOSITION 1.3. — For  $\varepsilon \in (0, 1]$ , there is a critical point  $q_\varepsilon(t) \in \Lambda$  of  $I_\varepsilon(q)$  such that

$$(i) \quad I_\varepsilon(q_\varepsilon) = b_\varepsilon, \quad (1.16)$$

$$(ii) \quad I'_\varepsilon(q_\varepsilon) = 0, \quad (1.17)$$

$$(iii) \quad \text{index } I''_\varepsilon(q_\varepsilon) \leq N - 2, \quad (1.18)$$

where  $\text{index } I''_\varepsilon(q_\varepsilon)$  is the Morse index of  $I''_\varepsilon(q_\varepsilon)$ .

Moreover there are constants  $M > m > 0$  such that

$$m \leq b_\varepsilon = I_\varepsilon(q_\varepsilon) \leq M \quad \text{for } \varepsilon \in (0, 1]. \quad (1.19)$$

*Proof.* — (1.19) follows from (1.15) and Proposition 1.2. Since  $I_\varepsilon(q)$  satisfies the strong force condition (SF) for  $\varepsilon \in (0, 1]$ , we have the following “Deformation Theorem”:

PROPOSITION 1.4. ([BR], Proposition 1.17). — Suppose  $\varepsilon \in (0, 1]$  and assume  $s > 0$  is not a critical value of  $I_\varepsilon(q)$ . Then for each  $\bar{a} > 0$  there is an  $a \in (0, \bar{a})$  and  $\eta \in C([0, 1] \times \Lambda, \Lambda)$  such that

1.  $\eta(1, q) = q$  if  $I_\varepsilon(q) \notin (s - \bar{a}, s + \bar{a})$ ,
2.  $I_\varepsilon(\eta(\tau, q)) \leq I_\varepsilon(q)$  for  $\tau \in [0, 1]$ ,
3.  $\eta(1, [I_\varepsilon \leq s + a]) \subset [I_\varepsilon \leq s - a]$ , where  $[I_\varepsilon \leq \sigma] = \{q \in \Lambda; I_\varepsilon(q) \leq \sigma\}$ . ■

By Proposition 1.2 and (1.15), we can see

$$b_\varepsilon \geq b_0 > 0 \quad \text{for all } \varepsilon \in (0, 1].$$

Using the property 3. of Proposition 1.4 in a standard way (cf. [R]), we can see  $b_\varepsilon > 0$  is a critical value of  $I_\varepsilon(q)$ .

As to the property (1.18), we can obtain it in a similar way to the proof of Theorem A of Tanaka [T]. In [T], we studied properties of Morse indices of critical values related to the symmetric mountain pass theorem and we got (1.16)-(1.18) for the symmetric mountain pass theorem. See also [BL], [BenFo], [Sc], [V], [LS]. ■

The above proposition ensures the existence of approximate solutions  $q_\varepsilon(t) \in \Lambda$  together with uniform estimates (1.17) and (1.18). We will get a solution of the original problem (HS) as a limit of  $q_\varepsilon(t)$  as  $\varepsilon \rightarrow 0$ .

To do so, we study the behavior of critical points of  $I_\varepsilon(q)$  whose critical values and Morse indices are uniformly bounded, that is, we study the behavior of critical points  $q_n(t) \in \Lambda$  such that

$$\begin{aligned} \varepsilon_n &\rightarrow 0, \\ I_{\varepsilon_n}(q_n) &\in [m, M], \\ I'_{\varepsilon_n}(q_n) &= 0, \\ \text{index } I''_{\varepsilon_n}(q_n) &\leq N - 2. \end{aligned}$$

The following proposition, which is due to Bahri and Rabinowitz [BR], ensures the existence of convergent subsequence of  $(q_n(t))$  and it shows the limit of the subsequence is a *generalized solution* of (HS).

PROPOSITION 1.5 (cf. [BR], Theorem 3.24). — *Let  $(\varepsilon_n)_{n=1}^\infty \subset (0, 1]$  be a sequence such that  $\varepsilon_n \rightarrow 0$ . Suppose  $(q_n(t))_{n=1}^\infty \subset \Lambda$  is a sequence of critical points of  $I_{\varepsilon_n}(q)$  such that*

$$I'_{\varepsilon_n}(q_n) = 0, \tag{1.20}$$

$$I_{\varepsilon_n}(q_n) \in [m, M] \text{ for all } n, \tag{1.21}$$

where  $0 < m < M$  are constants independent of  $n$ .

Then there is a subsequence — still denoted by  $n$  — and  $q_\infty(t) \in E$  such that

- (i)  $q_n(t)$  converges to  $q_\infty(t)$  weakly in  $E$  and strongly in  $L^\infty$ ;
- (ii)  $\int_0^T -V(q_\infty, t) dt < \infty$ ;
- (iii)  $q_\infty(t)$  vanishes on a set  $D$ , of measure 0;
- (iv)  $q_\infty(t) \in C^2(\mathbf{R} \setminus D, \mathbf{R})$ ;
- (v)  $q_\infty(t)$  satisfies (HS) on  $\mathbf{R} \setminus D$ .

Remark 1.6. — (i)  $(q_\varepsilon(t))_{\varepsilon \in (0, 1]}$  given in Proposition 1.3 satisfies the assumptions of the above proposition.

(ii)  $q_\infty(t)$  is a *generalized T-periodic solution* of (HS) in the sense of [BR].

Proof of Proposition 1.5. — By Lemma 1.1, we get

$$\|q_n\|_E, \quad \int_0^T -V(q_n, t) dt \leq C_6 \tag{1.22}$$

where  $C_6 > 0$  is independent of  $n$ .

Thus we get (i). By (1.22) and Fatou's lemma, we get (ii). We have (iii) easily from (ii). Since  $q_n(t)$  satisfies (1.2) with  $\varepsilon = \varepsilon_n$  and  $q_n(t) \rightarrow q_\infty(t)$  in  $L^\infty$ , we can deduce (iv) and (v). ■

If  $D = \emptyset$  in the above proposition, the limit function  $q_\infty(t)$  is a classical solution (non-collision solution) of the original problem (HS). In the following sections, we will show that for the sequence  $(q_\varepsilon(t))_{\varepsilon \in (0, 1]}$  given in Proposition 1.3.

(i) If  $V(q, t)$  satisfies (W1)-(W3) with  $\alpha \in (1, 2)$  in addition to (VI)-(V3), then  $D = \emptyset$ ;

(ii) if  $V(q, t)$  satisfies (W1)-(W3) with  $\alpha \in (0, 1]$  in addition to (V1)-(V3), then  $D \cap (0, T]$  consists of at most one point, that is,  $q_\infty(t)$  enters the singularity 0 at most one time in period  $T$ .

To get the above properties (i)-(ii), the uniform estimate of Morse indices (1.18) plays an important role. We remark that in Proposition 1.5, we used only the uniform bound of critical values.

Lastly in this section, we assume (W1)-(W3) in addition to (V1)-(V3) and get some a priori estimate, which will be used in the following sections.

PROPOSITION 1.7. — Assume (W1)-(W3) and (V2). For any  $0 < m < M$ , there are constants  $C_7(m, M)$ ,  $C_8(m, M) > 0$  independent of  $\varepsilon \in (0, 1]$  such that for all  $q \in \Lambda$  and  $\varepsilon \in (0, 1]$  with  $I_\varepsilon(q) \in [m, M]$  and  $I'_\varepsilon(q) = 0$

$$\begin{aligned} \text{(i)} \quad & \|q\|_E, \quad \int_0^T \frac{1}{|q|^\alpha} dt, \quad \int_0^T \frac{\varepsilon}{|q|^4} dt \leq C_7(m, M), \\ \text{(ii)} \quad & \left| \frac{1}{2} |\dot{q}(t)|^2 - \frac{1}{|q|^\alpha} + U(q, t) - \frac{\varepsilon}{|q|^4} \right| \leq C_8(m, M) \end{aligned} \quad (1.23)$$

for all  $t \in \mathbf{R}$ .

*Proof.* — We can get the assertion (i) from (W1)-(W3) and (i), (ii) of Lemma 1.1. To obtain (ii), we set

$$E(t) \equiv \frac{1}{2} |\dot{q}(t)|^2 - \frac{1}{|q|^\alpha} + U(q, t) - \frac{\varepsilon}{|q|^4}.$$

By (i), we get

$$\begin{aligned} \int_0^T |E(t)| dt &\leq \frac{1}{2} \int_0^T |\dot{q}|^2 dt \\ &+ \int_0^T -V(q, t) dt + \int_0^T \frac{\varepsilon}{|q|^4} dt \leq C_7(m, M). \end{aligned} \quad (1.24)$$

Since  $q(t) \in \Lambda$  is a solution of (1.2),

$$\frac{d}{dt} E(t) = U_t(q, t).$$

Thus by (W3)

$$\int_0^T \left| \frac{d}{dt} E(t) \right| dt \leq \int_0^T |U_t(q, t)| dt \leq C \int_0^T \frac{1}{|q(t)|^\alpha} dt \leq C_7''(m, M). \quad (1.25)$$

Combining (1.24) and (1.25), we get

$$\|E(t)\|_{L^\infty} \leq C_8(m, M).$$

Therefore we obtain (ii). ■

## 2. ASYMPTOTIC BEHAVIOR OF $q_n(t)$ NEAR COLLISION

In what follows, we assume (V2) and (W1)-(W3). Suppose  $(q_n(t)) \subset \Lambda$  be a sequence of critical points of  $I_{\varepsilon_n}(q)$  satisfying

$$\varepsilon_n \rightarrow 0, \quad (2.1)$$

$$I_{\varepsilon_n}(q_n) \in [m, M], \quad (2.2)$$

$$I'_{\varepsilon_n}(q_n) = 0, \quad (2.3)$$

$$q_n(t) \rightarrow q_\infty(t) \text{ weakly in } E \text{ and strongly in } L^\infty, \quad (2.4)$$

where  $0 < m < M$  are constants independent of  $n$ . By Proposition 1.5, a suitable subsequence of critical points  $(q_\varepsilon(t))_{\varepsilon \in (0, 1]} \subset \Lambda$ , which is obtained in Proposition 1.3, satisfies the conditions (2.1)-(2.4).

The main purpose of the following 3 sections is to prove Proposition 0.3, that is, to estimate the Morse index of  $I'_{\varepsilon_n}(q_n)$  from below by the number of collisions  $v$ :

$$v \equiv \#D = \# \{ t \in (0, T]; q_\infty(t) = 0 \}.$$

We can obtain Theorems 0.1 and 0.2 from Proposition 0.3 and (1.18). First we study the asymptotic behavior of  $q_n(t)$  near collisions. Suppose  $t_\infty \in (0, T]$  satisfies

$$q_\infty(t_\infty) = 0.$$

We may assume  $t_\infty \in (0, T)$  without loss of generality. Extracting a subsequence—still denoted by  $n$ —, we can choose  $t_n \in (0, T]$  such that

$$1. \quad |q_n(t_n)| \text{ takes its local minimum at } t = t_n, \quad (2.5)$$

$$2. \quad t_n \rightarrow t_\infty \text{ as } n \rightarrow \infty, \quad (2.6)$$

$$3. \quad |q_n(t_n)| \rightarrow 0 \text{ as } n \rightarrow \infty. \quad (2.7)$$

In fact, by (iii) of Proposition 1.5, we can find a sequence  $a_n, b_n \in (0, T)$  such that

$$\left. \begin{aligned} t_\infty - \frac{1}{n} < a_n < t_\infty < b_n < t_\infty + \frac{1}{n}, \\ q_\infty(a_n) > 0, \quad q_\infty(b_n) > 0. \end{aligned} \right\} \quad (2.8)$$

Thus we can find a sequence of integers  $m(1) < m(2) < \dots$  such that

$$|q_{m(n)}(t_\infty)| \leq \frac{1}{2} \min \{ |q_{m(n)}(a_n)|, |q_{m(n)}(b_n)| \}. \quad (2.9)$$

Suppose  $|q_{m(n)}(t_{m(n)})| = \min_{[a_n, b_n]} |q_{m(n)}(t)|$  for  $t_{m(n)} \in [a_n, b_n]$ . By (2.9),  $t_{m(n)} \in (a_n, b_n)$ . Thus  $|q_{m(n)}(t)|$  takes its local minimum at  $t = t_{m(n)}$ . Moreover we have  $t_{m(n)} \rightarrow t_\infty$  by (2.8) and  $|q_{m(n)}(t_{m(n)})| \leq |q_{m(n)}(t_\infty)| \rightarrow 0$ . Therefore we get (2.5)-(2.7) for the subsequence  $m(n)$ .

By Proposition 1.7,  $q_n(t)$  satisfies

$$\ddot{q}_n + \frac{\alpha q_n}{|q_n|^{\alpha+2}} - U_q(q_n, t) + \frac{4\varepsilon_n q_n}{|q_n|^6} = 0, \quad (2.10)$$

$$q_n(t+T) = q_n(t), \quad \text{in } \mathbf{R}, \quad (2.11)$$

$$\left| \frac{1}{2} |\dot{q}_n(t)|^2 - \frac{1}{|q_n|^\alpha} + U(q_n, t) - \frac{\varepsilon}{|q_n|^4} \right| \leq C_7(m, M). \quad (2.12)$$

We set

$$\delta_n = |q_n(t_n)| > 0 \quad (2.13)$$

and define  $x_n: \mathbf{R} \rightarrow \mathbf{R}^N \setminus \{0\}$  by

$$x_n(s) = \delta_n^{-1} q_n(\delta_n^{(\alpha+2)/2} s + t_n) \quad \text{for } s \in \mathbf{R}. \quad (2.14)$$

We consider the asymptotic behavior of  $x_n(s)$  as  $n \rightarrow \infty$ . From the definition of  $x_n(s)$  and (2.5)-(2.7), (2.10)-(2.13), we can easily see

LEMMA 2.1. —  $x_n(s)$  and  $\delta_n > 0$  satisfies

$$(i) \quad \delta_n \rightarrow 0, \quad (2.15)$$

$$(ii) \quad x_n(s) \text{ takes its local minimum at } s=0,$$

$$(iii) \quad |x_n(0)| = 1, \quad x_n(0) \perp \dot{x}_n(0), \quad (2.16)$$

$$(iv) \quad \ddot{x}_n(s) + \frac{\alpha x_n}{|x_n|^{\alpha+2}} - \delta_n^{\alpha+1} U_q(\delta_n x_n, \delta_n^{(\alpha+2)/2} s + t_n) + \frac{4\varepsilon_n}{\delta_n^{4-\alpha}} \frac{x_n}{|x_n|^6} = 0 \quad \text{in } \mathbf{R} \quad (2.17)$$

$$(v) \quad \left| \frac{1}{2} |\dot{x}_n(s)|^2 - \frac{1}{|x_n|^\alpha} + \delta_n^\alpha U(\delta_n x_n, \delta_n^{(\alpha+2)/2} s + t_n) - \frac{\varepsilon_n}{\delta_n^{4-\alpha}} \frac{1}{|x_n|^4} \right| \leq C_8(m, M) \delta_n^\alpha \quad \text{for all } s \in \mathbf{R} \text{ and } n \in \mathbf{N}. \quad \blacksquare \quad (2.18)$$

The following lemma gives us an estimate of the coefficient of the equation (2.17).

LEMMA 2.2.

$$\limsup_{n \rightarrow \infty} \frac{\varepsilon_n}{\delta_n^{4-\alpha}} \leq \frac{2-\alpha}{2}.$$

*Proof.* — Since  $|x_n(s)|^2$  takes its local minimum at  $s=0$ , we have

$$0 \leq \frac{1}{2} \frac{d^2}{ds^2} \Big|_{s=0} |x_n(s)|^2 = (\ddot{x}_n(0), x_n(0)) + |\dot{x}_n(0)|^2.$$

Using (2.16)-(2.18), we get

$$0 \leq (2-\alpha) - \frac{2\varepsilon_n}{\delta_n^{4-\alpha}} + \delta_n^{\alpha+1} (x_n(0), U_q(\delta_n x_n(0), t_n)) - 2\delta_n^\alpha U(\delta_n x_n(0), t_n) + 2C_8(m, M)\delta_n^\alpha.$$

By the assumption (W3), we can see

$$\limsup_{n \rightarrow \infty} \frac{\varepsilon_n}{\delta_n^{4-\alpha}} \leq \frac{2-\alpha}{2}. \quad \blacksquare$$

By Lemma 2.2, we can extract a subsequence—we still denote it by  $n$ —such that

$$\frac{\varepsilon_n}{\delta_n^{4-\alpha}} \rightarrow d \in \left[0, \frac{2-\alpha}{2}\right] \text{ as } n \rightarrow \infty.$$

Then we can deduce the following from (2.18).

$$|\dot{x}_n(0)| \rightarrow \sqrt{2(1+d)} \text{ as } n \rightarrow \infty. \tag{2.19}$$

We extract a subsequence again—still denoted by  $n$ —and by (2.16) we may assume

$$x_n(0) \rightarrow e_1, \tag{2.20}$$

$$\dot{x}_n(0) \rightarrow \sqrt{2(1+d)}e_2, \tag{2.21}$$

where  $e_1, e_2, \dots, e_N$  are an orthonormal basis of  $\mathbf{R}^N$ .

By the continuous dependence of solutions on initial data and equation, we have

**PROPOSITION 2.3.** — *For any  $l > 0$ ,  $x_n(s)$  converges to a function  $y_{\alpha,d}(s)$  in  $C^2([-l, l], \mathbf{R}^N)$ , where  $y_{\alpha,d}(s)$  is a solution of*

$$\ddot{y} + \frac{\alpha y}{|y|^{\alpha+2}} + 4d \frac{y}{|y|^6} = 0 \text{ in } \mathbf{R}, \tag{2.22}$$

$$y(0) = e_1, \tag{2.23}$$

$$\dot{y}(0) = \sqrt{2(1+d)}e_2. \tag{2.24}$$

*Proof.* — By (W3), we have for any  $R > 1$

$$\delta_n^{\alpha+1} U_q(\delta_n x, \delta_n^{(\alpha+2)/2} s + t_n) \rightarrow 0$$

in  $C^1(\{x \in \mathbf{R}^N; 1/R \leq |x| \leq R\} \times \mathbf{R}, \mathbf{R}^N)$  as  $n \rightarrow \infty$ . On the other hand, (2.22)-(2.24) has a global solution  $y_{\alpha,d}(s)$  satisfying

$$|y_{\alpha,d}(s)| \geq 1 \text{ for all } s \in \mathbf{R} \tag{2.25}$$

for  $0 < \alpha < 2$  and  $d \in \left[0, \frac{2-\alpha}{2}\right]$ . [The proof of (2.25) will be given in Lemma 4.2] Therefore we can see

$$x_n(s) \rightarrow y_{\alpha, d}(s) \quad \text{in } C^2([-l, l], \mathbf{R}^N)$$

for any  $l > 0$ . ■

Using Proposition 2.3, we will estimate the Morse index of  $I''_{\varepsilon_n}(q_n)$  for large  $n$  in the following sections.

### 3. MORSE INDEX OF $I''_{\varepsilon}(q)$ AND THE LIMIT PROBLEM

For arbitrary given  $l > 0$ , we define linear operator  $T_n: H_0^1(-l, l; \mathbf{R}) \rightarrow H_0^1(0, T; \mathbf{R})$  by

$$(T_n \varphi)(t) = \delta_n \varphi(\delta_n^{-(\alpha+2)/2}(t-t_n)) \quad (3.1)$$

for  $n \in \mathbf{N}$  and  $\varphi \in H_0^1(-l, l; \mathbf{R})$ . Remark that  $T_n$  is well-defined for large  $n$ .

Extending  $(T_n \varphi)(t)$  periodically, we regard it as a  $T$ -periodic function on  $\mathbf{R}$ .

We have for  $j = 3, \dots, N$

$$\begin{aligned} & \delta_n^{-(2-\alpha)/2} I''_{\varepsilon_n}(q_n)((T_n \varphi) e_j, (T_n \varphi) e_j) \\ &= \delta_n^{-(2-\alpha)/2} \int_0^T \left[ \left| \frac{d}{dt} (T_n \varphi) \right|^2 - \frac{\alpha |T_n \varphi|^2}{|q_n|^{\alpha+2}} + \frac{\alpha(\alpha+2)(q_n, e_j)^2 |T_n \varphi|^2}{|q_n|^{\alpha+4}} \right. \\ & \quad \left. - U_{qq}(q_n, t)((T_n \varphi) e_j, (T_n \varphi) e_j) - \frac{4\varepsilon_n |T_n \varphi|^2}{|q_n|^6} + \frac{24\varepsilon_n (q_n, e_j)^2 |T_n \varphi|^2}{|q_n|^8} \right] dt \\ &= \int_{-l}^l \left[ |\dot{\varphi}(s)|^2 - \frac{\alpha |\varphi|^2}{|x_n|^{\alpha+2}} + \frac{\alpha(\alpha+2)(x_n, e_j)^2 |\varphi|^2}{|x_n|^{\alpha+4}} \right. \\ & \quad \left. - \delta_n^{\alpha+2} U_{qq}(\delta_n x_n, \delta_n^{(\alpha+2)/2} s + t_n)(\varphi e_j, \varphi e_j) \right. \\ & \quad \left. - \frac{4\varepsilon_n |\varphi|^2}{\delta_n^{4-\alpha} |x_n|^6} + \frac{24\varepsilon_n (x_n, e_j)^2 |\varphi|^2}{\delta_n^{4-\alpha} |x_n|^8} \right] ds. \end{aligned}$$

By (W3) and Proposition 2.3, we have

$$\begin{aligned} & \delta_n^{-(2-\alpha)/2} I''_{\varepsilon_n}(q_n)((T_n \varphi) e_j, (T_n \varphi) e_j) \\ & \rightarrow \int_{-l}^l \left[ |\dot{\varphi}(s)|^2 - \frac{\alpha |\varphi|^2}{|y_{\alpha, d}|^{\alpha+2}} + \frac{\alpha(\alpha+2)(y_{\alpha, d}, e_j)^2 |\varphi|^2}{|y_{\alpha, d}|^{\alpha+4}} \right. \\ & \quad \left. - \frac{4d |\varphi|^2}{|y_{\alpha, d}|^6} + \frac{24d (y_{\alpha, d}, e_j)^2 |\varphi|^2}{|y_{\alpha, d}|^8} \right] ds \\ & = \int_{-l}^l \left[ |\dot{\varphi}(s)|^2 - \frac{\alpha |\varphi|^2}{|y_{\alpha, d}|^{\alpha+2}} - \frac{4d |\varphi|^2}{|y_{\alpha, d}|^6} \right] ds \quad (3.2) \end{aligned}$$

as  $n \rightarrow \infty$ .

Here we used the fact:

$$y_{\alpha, d}(s) \in \text{span} \{ e_1, e_2 \} \quad \text{for } s \in \mathbf{R}.$$

We set

$$J_{\alpha, d, l}(\varphi) = \int_{-l}^l \left[ |\dot{\varphi}(s)|^2 - \frac{\alpha |\varphi|^2}{|y_{\alpha, d}|^{\alpha+2}} \right] ds : H_0^1(-l, l; \mathbf{R}) \rightarrow \mathbf{R}$$

for  $\alpha \in (0, 2)$ ,  $d \in [0, (2-\alpha)/2]$  and  $l > 0$ . Then we can see

$$\lim_{n \rightarrow \infty} \delta_n^{-(2-\alpha)/2} I''_{\varepsilon_n}(q_n) ((T_n \varphi) e_j, (T_n \varphi) e_j) \leq J_{\alpha, d, l}(\varphi) \tag{3.3}$$

for all  $\varphi \in H_0^1(-l, l; \mathbf{R})$ .

We define

$$N(\alpha, d, l) = \max \{ \dim H; H \subset H_0^1(-l, l; \mathbf{R}) \text{ is a subspace such that } J_{\alpha, d, l}(\varphi) < 0 \text{ for } \varphi \in H \setminus \{0\} \}. \tag{3.4}$$

Clearly

$N(\alpha, d, l)$  = the number of negative eigenvalues of the following eigenvalue problem:

$$\left. \begin{aligned} -\ddot{u} - \frac{\alpha}{|y_{\alpha, d}(s)|^{\alpha+2}} u &= \lambda u \quad \text{in } (-l, l). \\ u(-l) &= u(l) = 0. \end{aligned} \right\} \tag{3.5}$$

We remark that  $N(\alpha, d, l)$  is a non-decreasing function of  $l$  for each  $\alpha$  and  $d$ . Let  $\varphi_i(s) \in H_0^1(-l, l; \mathbf{R})$  ( $i = 1, 2, \dots, N(\alpha, d, l)$ ) be eigenfunctions of the problem (3.5) with negative eigenvalues, in particular, we have

$$J_{\alpha, d, l}(\varphi) < 0 \quad \text{for } \varphi \in \text{span} \{ \varphi_i(s); i = 1, \dots, N(\alpha, d, l) \} \setminus \{0\}. \tag{3.6}$$

We consider the set of functions:

$$H(t_\infty, n) = \text{span} \{ (T_n \varphi_i) e_j; 1 \leq i \leq N(\alpha, d, l), 3 \leq j \leq N \} \subset E. \tag{3.7}$$

By (3.3) and (3.6), we can see for sufficiently large  $n$  that

$$I''_{\varepsilon_n}(q_n)(h, h) < 0 \quad \text{for all } h \in H(t_\infty, n) \setminus \{0\}. \tag{3.8}$$

We remark

$$\dim H(t_\infty, n) = (N-2) N(\alpha, d, l).$$

Finally we set

$$i(\alpha) = \sup_{l > 0} \min_{d \in [0, (2-\alpha)/2]} N(\alpha, d, l). \tag{3.9}$$

Choosing  $l > 0$  sufficiently large, we may assume

$$\dim H(t_\infty, n) \geq (N-2) i(\alpha). \tag{3.10}$$

In Section 4, we will give a representation (0.7) of  $i(\alpha)$ .

PROPOSITION 3.1. — Assume (V2) and (W1)-(W3) and suppose  $(q_n(t))_{n=1}^\infty \subset \Lambda$  satisfies (2.1)-(2.4). Let  $\nu$  be the number of times  $q_\infty(t)$  enters the singularity 0:

$$\nu = \# \{ t \in (0, T]; q_\infty(t) = 0 \}.$$

Then we have

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) \geq (N-2) i(\alpha) \nu. \quad (3.11)$$

*Proof.* — Suppose  $\nu < \infty$  and

$$\{ t_{\infty, 1}, t_{\infty, 2}, \dots, t_{\infty, \nu} \} = \{ t \in (0, T]; q_\infty(t) = 0 \}.$$

For any given subsequence  $n_m \rightarrow \infty$ , we can extract a subsequence — we still denote it by  $n_m$  — such that Proposition 2.3 holds for each  $t_{\infty, k}$  for suitable orthonormal basis  $e_1^{(k)}, e_2^{(k)}, \dots, e_N^{(k)}$  and  $d^{(k)} \in \left[ 0, \frac{2-\alpha}{2} \right]$ . Thus we can construct subspaces  $H(t_{\infty, k}, n_m) \subset E$  for each  $t_{\infty, k}$  ( $k=1, 2, \dots, \nu$ ) as in (3.7). From the construction, we have

$$\dim H(t_{\infty, k}, n_m) \geq (N-2) i(\alpha) \quad \text{for all } k.$$

For any  $\delta > 0$ , we find a constant  $m_0(\delta) \in \mathbb{N}$  such that

$$\text{supp } h(t) \subset [t_{\infty, k} - \delta, t_{\infty, k} + \delta]$$

for all  $h(t) \in H(t_{\infty, k}, n_m)$  and  $m \geq m_0(\delta)$ . Thus we get

$$H(t_{\infty, i}, n_m) \cap H(t_{\infty, j}, n_m) = \{0\} \quad (i \neq j)$$

for sufficiently large  $m$ . Set

$$H_{n_m} = H(t_{\infty, 1}, n_m) \oplus H(t_{\infty, 2}, n_m) \oplus \dots \oplus H(t_{\infty, \nu}, n_m).$$

Choosing sufficiently large  $l > 0$ , we obtain from (3.8) and (3.10) that

$$\begin{aligned} \dim H_{n_m} &\geq (N-2) i(\alpha) \nu, \\ I''_{\varepsilon_n}(q_{n_m})(h, h) &< 0 \quad \text{for } h \in H_{n_m} \setminus \{0\} \end{aligned}$$

for sufficiently large  $m$ .

Therefore we get (3.11). In case  $\nu = \infty$ , for any  $k \in \mathbb{N}$  we can see in a similar way that

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) \geq (N-2) i(\alpha) k.$$

Thus we conclude

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) = \infty. \quad \blacksquare$$

#### 4. REPRESENTATION OF THE NUMBER $i(\alpha)$ AND PROOF OF PROPOSITION 0.3

The aim of this section is to give a representation (0.7) of the number  $i(\alpha)$ , that is, to prove

PROPOSITION 4.1. — *Let  $i(\alpha) \in \mathbf{N}$  be the number defined in (3.4)-(3.9). Then for any  $\alpha \in (0, 2)$  the number  $i(\alpha)$  can be represented as*

$$i(\alpha) = \max \left\{ m \in \mathbf{N}; m < \frac{2}{2-\alpha} \right\}.$$

We remark

$$i(\alpha) = 1 \quad \text{for } 0 < \alpha \leq 1, \quad (4.1)$$

$$i(\alpha) \geq 2 \quad \text{for } 1 < \alpha < 2, \quad (4.2)$$

$$i(\alpha) \rightarrow \infty \quad \text{as } \alpha \rightarrow 2. \quad (4.3)$$

First, we consider the solution  $y_{\alpha, d}(s)$  of (2.22)-(2.24).

LEMMA 4.2. — *For any  $0 < \alpha < 2$ ,  $d \in \left[ 0, \frac{2-\alpha}{2} \right]$ , the equation (2.22)-(2.24) has a global solution  $y_{\alpha, d}(s)$ . Moreover,  $y_{\alpha, d}(s)$  satisfies*

$$1 \leq |y_{\alpha, d}(s)| \leq |y_{\alpha, 0}(s)| \quad (4.4)$$

for all  $d \in \left[ 0, \frac{2-\alpha}{2} \right]$  and  $s \in \mathbf{R}$ .

*Proof.* — First we remark that  $y_{\alpha, d}(s)$  satisfies

$$\frac{1}{2} |\dot{y}_{\alpha, d}(s)|^2 - \frac{1}{|y_{\alpha, d}|^\alpha} - \frac{d}{|y_{\alpha, d}|^4} = 0 \quad \text{for } s \in \mathbf{R}. \quad (4.5)$$

We fix here  $\alpha \in (0, 2)$  and set  $R_d(s) = |y_{\alpha, d}(s)|^2$ . Using (2.22) and (4.5), we get

$$\ddot{R}_d = 2(\ddot{y}_{\alpha, d}, y_{\alpha, d}) + 2|\dot{y}_{\alpha, d}|^2 \quad (4.6)$$

$$= 2(2-\alpha) \frac{1}{R_d^{\alpha/2}} - 4d \frac{1}{R_d^2},$$

$$R_d(0) = 1, \quad (4.7)$$

$$\dot{R}_d(0) = 0. \quad (4.8)$$

We can easily see from (4.6)-(4.8) that  $R_{(2-\alpha)/2}(s) \equiv 1$  and for  $d \in \left[ 0, \frac{2-\alpha}{2} \right)$

$$\ddot{R}_d(0) = 2(2-\alpha) - 4d > 0$$

and

$$\ddot{R}_d(s) \geq 2(2-\alpha) \left( \frac{1}{R_d^{\alpha/2}} - \frac{1}{R_d^2} \right) > 0 \quad \text{if } R_d(s) > 1.$$

Thus we get for  $d \in \left[ 0, \frac{2-\alpha}{2} \right)$

$$\left. \begin{array}{l} R_d(s) > 1 \quad \text{for all } s \neq 0, \\ s \dot{R}_d(s) > 0 \quad \text{for all } s \neq 0. \end{array} \right\} \quad (4.9)$$

Next we fix  $d \in \left( 0, \frac{2-\alpha}{2} \right)$  and prove  $R_d(s) < R_0(s)$  for all  $s$ . Since

$\ddot{R}_d(0) = 2(2-\alpha) - 4d < 2(2-\alpha) = \ddot{R}_0(0)$  for  $d \in \left( 0, \frac{2-\alpha}{2} \right)$ , we have

$R_d(s) < R_0(s)$  for sufficiently small  $s > 0$ .

Suppose there is an  $s_1 > 0$  such that  $R_d(s_1) = R_0(s_1)$ . Then there is an  $s_0 > 0$  such that

$$\left. \begin{array}{l} R_d(s) < R_0(s) \quad \text{for } s \in (0, s_0), \\ R_d(s_0) = R_0(s_0). \end{array} \right\} \quad (4.10)$$

Since  $R_d(s)$  satisfies (4.6)-(4.8), we have

$$\frac{1}{2} \dot{R}_d(s)^2 - 4 R_d^{(2-\alpha)/2} - \frac{4d}{R_d} = -4(1+d).$$

Thus we get

$$\frac{1}{2} \dot{R}_d(s_0)^2 - \frac{1}{2} \dot{R}_0(s_0)^2 = 4d \left( \frac{1}{R_d(s_0)} - 1 \right) < 0.$$

By (4.9), we get  $\dot{R}_d(s_0) < \dot{R}_0(s_0)$ . But this contradicts with (4.10). Therefore we have

$$R_d(s) < R_0(s) \quad \text{for } s > 0.$$

Similarly we get  $R_d(s) < R_0(s)$  for  $s < 0$ . ■

**COROLLARY 4.3.** — For any  $0 < \alpha < 2$ ,  $d \in \left[ 0, \frac{2-\alpha}{2} \right]$  and  $l > 0$ ,

$$N(\alpha, 0, l) \leq N(\alpha, d, l),$$

i. e.,

$$i(\alpha) = \sup_{l > 0} N(\alpha, 0, l). \quad (4.11)$$

*Proof.* — By (4.4), we have

$$J_{\alpha, d, l}(\varphi) \leq J_{\alpha, 0, l}(\varphi) \quad \text{for all } \varphi \in H_0^1(-l, l; \mathbf{R}).$$

Thus we get the desired result from the definition of  $N(\alpha, d, l)$  and  $i(\alpha)$ . ■

By (4.11), from now on, we deal with only the case  $d=0$ . The following lemma is a consequence of Sturm Comparison Theorem.

LEMMA 4.4. — *The number  $i(\alpha) + 1$  is equal to the maximal number of zeros of nontrivial solutions  $u(s)$  of*

$$-\ddot{u} - \frac{\alpha}{|y_{\alpha,0}(s)|^{\alpha+2}} u = 0 \quad \text{in } \mathbf{R}. \tag{4.12}$$

That is,

$$i(\alpha) + 1 = \max \{ \# \{s \in \mathbf{R}; u(s) = 0\}; u(s) \text{ is a nontrivial solution of (4.12)} \}.$$

*Proof.* — Suppose  $i(\alpha) = k$  and let  $l > 0$  be sufficiently large so that  $N(\alpha, 0, l) = k$ . Then  $k$ -th eigenvalue  $\lambda_k$  of (3.5) is negative, that is, there is an eigenfunction  $u_k(s)$  of

$$\begin{aligned} -\ddot{u}_k - \frac{\alpha}{|y_{\alpha,0}(s)|^{\alpha+2}} u_k &= \lambda_k u_k \quad \text{in } (-l, l), \\ u_k(\pm l) &= 0, \end{aligned}$$

which has exactly  $(k + 1)$  zeros in  $[-l, l]$ . Consider initial value problem (4.12) with initial data  $u(-l) = 0$  and  $\dot{u}(-l) = 1$ , then by Sturm Comparison Theorem,  $u(s)$  has at least  $(k + 1)$  zeros in  $[-l, l]$ .

Conversely, suppose (4.12) has a nontrivial solution with  $(k + 1)$  zeros  $t = t_1 < t_2 < \dots < t_{k+1}$  and consider the eigenvalue problem:

$$\begin{aligned} -\ddot{u} - \frac{\alpha}{|y_{\alpha,0}(s)|^{\alpha+2}} u &= \lambda u \quad \text{in } (t_1, t_{k+1}), \\ u(t_1) = u(t_{k+1}) &= 0. \end{aligned}$$

Then we can see that the  $k$ -th eigenvalue  $\lambda_k$  equals to 0. Choosing  $l > 0$  such that  $[t_1, t_{k+1}] \subset (-l, l)$ , we have  $N(\alpha, 0, l) \geq k$ . ■

Therefore we will consider the number of zeros of nontrivial solutions  $u(s)$  of (4.12). We write  $y_i^{(\alpha)}(s) = (y_{\alpha,0}(s), e_i) : \mathbf{R} \rightarrow \mathbf{R}$  ( $i = 1, 2$ ). Then  $\{y_1^{(\alpha)}, y_2^{(\alpha)}\}$  are linearly independent solutions of (4.12). Thus any solution  $u(s)$  of (4.12) can be represented by their linear combinations. That is, we can write

$$u(s) = \sin \beta y_1^{(\alpha)}(s) + \cos \beta y_2^{(\alpha)}(s) \quad (\beta \in \mathbf{R})$$

up to multiplicative constants. Using polar coordinate  $(r_\alpha, \theta_\alpha)$ , we write

$$(y_1^{(\alpha)}(s), y_2^{(\alpha)}(s)) = (r_\alpha(s) \cos \theta_\alpha(s), r_\alpha(s) \sin \theta_\alpha(s)) \tag{4.13}$$

where  $r_\alpha(s) > 0$  and  $\theta_\alpha(s) \in \mathbf{R}$  with  $\theta_\alpha(0) = 0$ . Then any solution  $u(s)$  of (4.12) can be written (up to multiplicative constants) as

$$u(s) = r_\alpha(s) \sin(\theta_\alpha(s) + \beta) \quad (\beta \in \mathbf{R}). \tag{4.14}$$

From (4.14), we can easily see

LEMMA 4.5. — *The maximal number of zeros of nontrivial solutions of (4.12) is equal to the number*

$$\max \left\{ m \in \mathbf{Z}; m < \frac{\theta_{\alpha}^{+} - \theta_{\alpha}^{-}}{\pi} \right\} + 1. \quad (4.15)$$

Here  $\theta_{\alpha}^{\pm}$  is defined by

$$\theta_{\alpha}^{\pm} = \lim_{s \rightarrow \pm \infty} \theta_{\alpha}(s). \quad \blacksquare$$

REMARK 4.6. — The number (4.15) describes twice of the number of times the point  $(y_1^{(\alpha)}(s), y_2^{(\alpha)}(s))$  turns around the singularity 0 while  $-\infty < s < \infty$ .

*Proof.* — For  $u_{\beta}(s) = r_{\alpha}(s) \sin(\theta_{\alpha}(s) + \beta)$ , we can easily see

$$u_{\beta}(s) = 0 \quad \text{if and only if} \quad \theta_{\alpha}(s) + \beta = m\pi \quad \text{for some } m \in \mathbf{Z}.$$

Thus we can see the maximal number of zeros of nontrivial solutions of (4.12) is equal to the number (4.15), that is,

$$\begin{aligned} \max \{ \# \{ s \in \mathbf{R}; u_{\beta}(s) = 0 \}; \beta \in (0, 2\pi] \} \\ = \max \left\{ m \in \mathbf{Z}; m < \frac{\theta_{\alpha}^{+} - \theta_{\alpha}^{-}}{\pi} \right\} + 1. \quad \blacksquare \end{aligned}$$

*Proof of Proposition 4.1.* — Since

$$y(s) = y_{\alpha, 0}(s) = (r_{\alpha}(s) \cos \theta_{\alpha}(s), r_{\alpha}(s) \sin \theta_{\alpha}(s))$$

satisfies (2.22)-(2.24) with  $d=0$ , we have

$$r_{\alpha}(s)^2 \dot{\theta}_{\alpha}(s) = \sqrt{2} \quad \text{for all } s \in \mathbf{R} \quad (4.16)$$

(conservation of the angular momentum). Thus we can make a change of independent variables  $s \rightarrow \theta = \theta_{\alpha}$ . We set  $\rho_{\alpha} = \rho_{\alpha}(\theta) = \frac{1}{r_{\alpha}(\theta)}$ . Then  $\rho_{\alpha}(\theta)$  satisfies

$$(\rho_{\alpha})_{\theta\theta} + \rho_{\alpha} - \frac{\alpha}{2} (\rho_{\alpha})^{\alpha-1} = 0, \quad (4.17)$$

$$\rho_{\alpha}(0) = 1, \quad (4.18)$$

$$(\rho_{\alpha})_{\theta}(0) = 0, \quad (4.19)$$

and  $\theta_{\alpha}^{\pm}$  can be characterized as

$$\theta_{\alpha}^{\pm} = \pm \sup \{ \theta > 0; \rho_{\alpha}(\tau) \text{ exists and is positive for all } \tau \in [0, \theta) \} \quad (4.20)$$

By (4.16)-(4.19), we have

$$(\rho_{\alpha})_{\theta}(\theta)^2 + \rho_{\alpha}(\theta)^2 - \rho_{\alpha}(\theta)^{\alpha} = 0 \quad \text{for all } \theta \in (\theta_{\alpha}^{-}, \theta_{\alpha}^{+}).$$

Since  $(\rho_{\alpha})_{\theta}(\theta) < 0$  for all  $\theta > 0$  (it follows from (4.9)), we have

$$\frac{-(\rho_{\alpha})_{\theta}(\theta)}{\sqrt{\rho_{\alpha}(\theta)^{\alpha} - \rho_{\alpha}(\theta)^2}} = 1.$$

Integrating over  $[0, \theta]$ , we get

$$\int_{\rho_\alpha(\theta)}^1 \frac{d\rho}{\sqrt{\rho^\alpha - \rho^2}} = \theta \quad \text{for all } \theta \in [0, \theta_\alpha^+).$$

By (4.20), we can see

$$\theta_\alpha^\pm = \pm \int_0^1 \frac{d\rho}{\sqrt{\rho^\alpha - \rho^2}} = \pm \frac{\pi}{2 - \alpha}. \quad (4.21)$$

Thus by Lemmas 4.4, 4.5 and (4.21), we obtain Proposition 4.1. ■

*Proof of Proposition 0.3.* — We can easily deduce Proposition 0.3 from Proposition 3.1 and 4.1. ■

## 5. PROOFS OF THEOREMS 0.1 AND 0.2

Now we can deduce Theorems 0.1 and 0.2 from Propositions 1.3, 0.3 and (4.1)-(4.2).

*Proof of Theorem 0.1.* — Let  $(q_\varepsilon(t))_{\varepsilon \in (0, 1]}$  be a sequence of critical points given in Proposition 1.3. By Proposition 1.5, we can extract a subsequence  $\varepsilon_n \rightarrow 0$  such that  $q_n(t) = q_{\varepsilon_n}(t)$  satisfies the assumptions of Proposition 0.3. Since  $i(\alpha) \geq 2$  for  $\alpha \in (1, 2)$ , we have from Proposition 0.3

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) \geq 2(N-2)v.$$

Comparing with (1.18), we can see

$$v = 0.$$

That is,  $q_\infty(t)$  does not enter the singularity 0 and  $q_\infty(t)$  is a non-collision T-periodic solution of (HS). ■

*Proof of Theorem 0.2.* — Proof of Theorem 0.2 can be done in a similar way to the proof of Theorem 0.1. However, by (4.1),  $i(\alpha) = 1$  for  $\alpha \in (0, 1]$ . Thus

$$\liminf_{n \rightarrow \infty} \text{index } I''_{\varepsilon_n}(q_n) \geq (N-2)v.$$

Comparing with (1.18), we get

$$v \leq 1.$$

This is the desired result. ■

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## REFERENCES

- [AC1] A. AMBROSETTI and V. COTI ZELATI, Critical points with lack of compactness and applications to singular dynamical system, *Annali Mat. Pura Appl.*, Vol. **149**, 1987, pp. 237-259.
- [AC2] A. AMBROSETTI and V. COTI ZELATI, Periodic solutions of singular dynamical systems, in *Periodic solutions of Hamiltonian systems and related topics*, P. H. RABINOWITZ *et al.* Eds., V209, NATO ASI Series, Reidel, 1987, pp. 1-10.
- [AC3] A. AMBROSETTI and V. COTI ZELATI, Noncollision orbits for a class of Keplerian-like potentials, *Ann. Inst. Henri Poincaré, Analyse non linéaire*, Vol. **5**, 1988, pp. 287-295.
- [AC4] A. AMBROSETTI and V. COTI ZELATI, Perturbation of Hamiltonian systems with Keplerian potentials, *Math. Z.*, Vol. **201**, 1989, pp. 227-242.
- [BL] A. BAHRI and P. L. LIONS, Morse index of some min-max critical points, I. Application to multiplicity results, *Comm. Pure Appl. Math.*, Vol. **41**, 1988, pp. 1027-1037.
- [BR] A. BAHRI and P. H. RABINOWITZ, A minimax method for a class of Hamiltonian systems with singular potentials, *J. Funct. Anal.*, Vol. **82**, 1989, pp. 412-428.
- [C] V. COTI ZELATI, Periodic solutions for a class of planar, singular dynamical systems, *J. Math. Pure Appl.*, Vol. **68**, 1989, pp. 109-119.
- [DG] M. DEGIOVANNI and F. GIANNONI, Dynamical systems with Newtonian type potentials, *Annali Scuola Norm. Sup. Pisa*, Vol. **15**, 1988, pp. 467-494.
- [DGM] M. DEGIOVANNI, F. GIANNONI and A. MARINO, Periodic solutions of dynamical systems with Newtonian type potentials, in *Periodic solutions of Hamiltonian systems and related topics*, P. H. RABINOWITZ *et al.* Eds., V209, NATO ASI Series, Reidel, 1987, pp. 111-115.
- [Go] W. B. GORDON, Conservative dynamical systems involving strong forces, *Trans. Amer. Math. Soc.*, Vol. **204**, 1975, pp. 113-135.
- [Gr1] C. GRECO, Periodic solutions of a class of singular Hamiltonian systems, *Nonlinear Analysis*, T.M.A., Vol. **12**, 1988, pp. 259-269.
- [GR2] C. GRECO, Remarks on periodic solutions for some dynamical systems with singularities, in *Periodic solutions of Hamiltonian systems and related topics*, P. H. RABINOWITZ *et al.* Eds., V209, NATO ASI Series, Reidel, 1987, pp. 169-173.
- [LS] A. C. LAZER and S. SOLIMINI, Nontrivial solutions of operator equations and Morse indices of critical points of min-max type, *Nonlinear Analysis*, T.M.A., Vol. **12**, 1988, pp. 761-775.
- [LF] L. A. LYUSTERNIK and A. I. FET, Variational problems on closed manifolds, *Dokl. Akad. Nauk. USSR*, N.S., Vol. **81**, 1951, pp. 17-18.
- [R] P. H. RABINOWITZ, Minimax methods in critical point theory with applications to differential equations, CBMS Regional Conf. Ser. in Math., Vol. **65**, Amer. Math. Soc., Providence, 1986.
- [Sc] J. T. SCHWARTZ, *Nonlinear functional Analysis*, Gordon and Breach, New York, 1969.
- [ST] E. SERRA and S. TERRACINI, Noncollision solutions to some singular minimization problems with Keplerian-like potentials, preprint.

- [T] K. TANAKA, Morse indices at critical points related to the symmetric mountain pass theorem and applications, *Comm. Partial Diff. Eq.*, Vol. **14**, 1989, pp. 99-128.
- [V] C. VITERBO, Indice de Morse des points critiques obtenus par minimax, *Ann. Inst. Henri Poincaré, Analyse non linéaire*, Vol. **5**, 1988, pp. 221-225.

After the Author submitted this paper, he learned works [BenFo] from Professor V. Benci, [Bea] from Professor A. Bahri, and [CS] from professor P. H. Rabinowitz. In particular, in [Bea], the Morse indices of collision solutions is studied and in [CS], the result closely related our Theorem 0-1 is obtained by different argument.

- [BenFo] V. BENCI and D. FORTUNATO, Subharmonic solutions of prescribed minimal period for autonomous differential equations, in *Recent advances in Hamiltonian systems*. Dell'Antonio e D'Onofrio Eds., World Sci., Singapore, 1986.
- [Bea] A. BEALIEU, Étude de solutions généralisées pour un système hamiltonien avec potentiel singulier, *Duke Math. J.*, Vol. **67**, 1992, pp. 21-37.
- [CS] V. COTI ZELATI and E. SERRA, *Collisions and non-collisions solutions for a class of Keplerian-like dynamical system*, preprint.

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