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# T. K. CARNE <br> Brownian motion and stereographic projection 

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# Brownian motion and stereographic projection 

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Abstract. - Stereographic projection from $R^{N}$ to $S^{N}$ maps Brownian paths in $\mathbf{R}^{\mathrm{N}}$ to the paths of Brownian motion on $S^{\mathrm{N}}$ conditioned to be at the centre of the projection at a negative exponential time.

Key-words: Stereographic projection; Conditioned Brownian motion; Conformal transformations.

Résumé. - La projection stéréographique de $\mathrm{R}^{\mathrm{N}}$ à $\mathrm{S}^{\mathrm{N}}$ applique les trajectoires Browniennes de $\mathrm{R}^{\mathrm{N}}$ sur les trajectoires Browniennes de $\mathrm{S}^{\mathbf{N}}$ conditionnées par le fait d'être au centre de projection à un instant de loi exponentielle.

In this brief note we shall discuss how Brownian motion in $\mathbf{R}^{\mathbf{N}}$, for $\mathbf{N} \geqslant 3$, can be interpreted as a Brownian bridge conditioned to go to the «ideal point at infinity ». This question was posed by Prof. L. Schwartz [2]. Prof. M. Yor [3] presents an alternative, more probabilistic, approach.

## 1. STEREOGRAPHIC PROJECTION

Consider the unit sphere $\mathrm{S}^{\mathrm{N}}$ in $\mathrm{R}^{\mathrm{N}+1}$ and the hyperplane

$$
\mathbf{R}^{\mathbf{N}}=\left\{y=\left(y_{1}, \ldots, y_{\mathbf{N}+1}\right): y_{\mathbf{N}+1}=0\right\}
$$

Stereographic projection from the point $\mathrm{P}=(0, \ldots, 0,1)$ of $S^{\mathrm{N}}$ maps $y \in \mathbb{R}^{\mathrm{N}}$
to the point $x \in \mathrm{~S}^{\mathbf{N}} \backslash\{\mathrm{P}\}$ which lies on the straight line from P through $y$; see the diagram. This is a diffeomorphism between $\mathbf{S}^{\mathbf{N}} \backslash\{\mathbf{P}\}$ and $\mathbf{R}^{\mathbf{N}}$, so we regard $P$ as being the point of $S^{\mathbf{N}}$ which corresponds to the «ideal point at infinity of $\mathrm{R}^{\mathrm{N}}$ ».

Proposition 1. - Brownian motion on $\mathrm{R}^{\mathrm{N}}$ is mapped by stereographic projection onto a time changed version of the Brownian motion on $\mathrm{S}^{\mathrm{N}}$ together with a drift towards P at speed $\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta$ on the sphere.

Proof. - Brownian motion on a Riemannian manifold with metric $g_{a b} d x_{a} d x_{b}$ has as its infinitesimal generator one half of the Laplacian, viz.

$$
\frac{1}{2} \Delta=\frac{1}{2 \sqrt{g}} \sum \frac{\partial}{\partial x_{a}}\left(\sqrt{g} g^{a b} \frac{\partial}{\partial x_{b}}\right)
$$

where $g=\operatorname{det}\left(g_{a b}\right)$ and $\left(g^{a b}\right)=\left(g_{a b}\right)^{-1}$. On $S^{\mathrm{N}}$ take co-ordinates $(\theta, z)$ for $x \in \mathrm{~S}^{\mathrm{N}}$ where $0 \leqslant \theta \leqslant \pi$ is the angle shown in the diagram and $z=y /\|y\| \in \mathbf{S}^{\mathbf{N}-1}=\mathbf{S}^{\mathbf{N}} \cap \mathbf{R}^{\mathbf{N}}$.

Then

$$
\|d x\|^{2}=|d \theta|^{2}+\sin ^{2} \theta \cdot\|d z\|^{2}
$$

so the Laplacian on $S^{N}$ is

$$
\Delta_{\mathrm{S}^{\mathrm{N}}}=\frac{1}{\sin ^{\mathrm{N}-1} \theta} \frac{\partial}{\partial \theta}\left(\sin ^{\mathrm{N}-1} \theta \frac{\partial}{\partial \theta}\right)+\frac{1}{\sin ^{2} \theta} \Delta_{\mathrm{S}^{\mathrm{N}-1}}
$$

Similarly, if we take co-ordinates $(r, z)$ for $y \in \mathbb{R}^{\mathbf{N}}$, where $r=\|y\|$, then

$$
\|d y\|^{2}=|d r|^{2}+r^{2}\|d z\|^{2}
$$

so the usual Laplacian on $\mathbb{R}^{N}$ is

$$
\Delta_{\mathrm{R}^{\mathrm{N}}}=\frac{1}{r^{\mathbf{N}-1}} \frac{\partial}{\partial r}\left(r^{\mathrm{N}-1} \frac{\partial}{\partial r}\right)+\frac{1}{r^{2}} \Delta_{\mathrm{S}^{\mathrm{N}-1}}
$$

The infinitesimal generator for the deterministic motion given by a drift towards $P$ at speed $\frac{1}{2}(N-2) \tan \frac{1}{2} \theta$ is clearly

$$
\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta \cdot \frac{\partial}{\partial \theta}
$$

Hence, to prove the proposition we need to show that, under stereo-
graphic projection $\frac{1}{2} \Delta_{\mathbf{R}^{\mathrm{N}}}$ corresponds to some strictly positive function timi,

$$
\mathscr{G}_{\mathrm{P}}=\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}+\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta \cdot \frac{\partial}{\partial \theta} .
$$

Under stereographic projection we have $r=\tan \frac{1}{2} \theta$ so

$$
\begin{aligned}
\Delta_{\mathrm{S}^{\mathrm{N}}} & =\left(\frac{2 r}{1+r^{2}}\right)^{1-\mathrm{N}}\left(\frac{1+r^{2}}{2}\right) \frac{\partial}{\partial r}\left[\left(\frac{2 r}{1+r^{2}}\right)^{\mathrm{N}-1}\left(\frac{1+r^{2}}{2}\right) \frac{\partial}{\partial r}\right]+\left(\frac{1+r^{2}}{2 r}\right)^{2} \Delta_{\mathrm{S}^{\mathrm{N}-1}} \\
& =\left(\frac{1+r^{2}}{2}\right)^{2}\left\{\left(\frac{2}{1+r^{2}}\right)^{2-\mathrm{N}} \frac{1}{r^{\mathrm{N}-1}} \frac{\partial}{\partial r}\left[\left(\frac{2}{1+r^{2}}\right)^{\mathrm{N}-2} r^{\mathrm{N}-1} \frac{\partial}{\partial r}\right]+\frac{1}{r^{2}} \Delta_{\mathrm{S}^{\mathrm{N}-1}}\right\} \\
& =\left(\frac{1+r^{2}}{2}\right)^{2}\left\{\frac{1}{r^{\mathrm{N}-1}} \frac{\partial}{\partial r}\left[r^{\mathrm{N}-1} \frac{\partial}{\partial r}\right]-(\mathrm{N}-2)\left(\frac{2 r}{1+r^{2}}\right) \frac{\partial}{\partial r}+\frac{1}{r^{2}} \Delta_{\mathrm{S}^{\mathrm{N}-1}}\right\} \\
& =\left(\frac{1+r^{2}}{2}\right)^{2}\left\{\Delta_{\mathrm{R}^{\mathrm{N}}}-(\mathrm{N}-2)\left(\frac{2 r}{1+r^{2}}\right) \frac{\partial}{\partial r}\right\}
\end{aligned}
$$



Equivalently,

$$
\begin{aligned}
\frac{1}{2} \Delta_{\mathrm{R}^{\mathrm{N}}} & =\left(\frac{2}{1+r^{2}}\right)^{2}\left\{\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}+\frac{1}{2}(\mathrm{~N}-2) r\left(\frac{1+r^{2}}{2}\right) \frac{\partial}{\partial r}\right\} \\
& =(1+\cos \theta)^{2}\left\{\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}+\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta \frac{\partial}{\partial \theta}\right\}
\end{aligned}
$$

This completes the proof.
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We now wish to obtain the random process with infinitesimal generator $\mathscr{G}_{\mathrm{P}}$ by conditioning the standard Brownian motion $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ on the sphere to be at P at an appropriate time. To do this we will follow the analysis of conditioning given by J. L. Doob [1, Chapter 10]. Note that we are seeking a time-homogeneous process, so that conditioning $\operatorname{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ to be at P at a fixed time will not do. Furthermore, we cannot simply condition $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ to hit P at some time since, to do so, we would require a positive harmonic function on $S^{N} \backslash\{P\}$ with a singularity at $P$. No such function exists. However, we do obtain time homogeneous processes by conditioning $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ to be at P at a random time T which is independent of $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ and has a negative exponential distribution.

Proposition 2. - Let T be a random time which is independent of $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ and has a negative exponential distribution with parameter $\lambda=\mathrm{N}(\mathrm{N}-2) / 8$. Then $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ conditioned to be at P at time T has infinitesimal generator

$$
\mathscr{G}_{\mathrm{P}}=\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}+\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta \frac{\partial}{\partial \theta}
$$

on $\mathbf{S}^{\mathbf{N}} \backslash\{\mathrm{P}\}$. Hence, $\mathrm{BM}\left(\mathbf{R}^{\mathbf{N}}\right)$ is mapped by stereographic projection to a time-changed version of $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ conditioned to be at P at the time T .

Proof. - To condition $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ to be at P at time T we need to find a positive function $h$ on $S^{\mathrm{N}} \backslash\{\mathrm{P}\}$ with a singularity at P and

$$
\left(\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}-\lambda \mathrm{I}\right) h=0
$$

Then the conditioned process will have the $h$-transform:

$$
u \rightarrow h^{-1}\left(\frac{1}{2} \Delta_{\mathrm{S}^{N}}-\lambda \mathrm{I}\right)(h . u)
$$

as its infinitesimal generator. Such a function $h$ must be a multiple of the Green's function for $\frac{1}{2} \Delta_{\mathbf{S}^{N}}-\lambda I$ with a pole at $P$ and hence it must be a function of $\theta$ only. Thus we wish to solve

$$
\frac{1}{2 \sin ^{\mathbf{N}-1} \theta} \frac{\partial}{\partial \theta}\left[\sin ^{\mathrm{N}-1} \theta \frac{\partial h}{\partial \theta}\right]-\lambda h=0 .
$$

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When $\lambda=\mathrm{N}(\mathrm{N}-2) / 8$ the required function $h$ is given by $h=\left(\cos \frac{1}{2} \theta\right)^{-w+2}$ Consequently, the conditioned process has infinitesimal generator

$$
\begin{aligned}
u & \rightarrow h^{-1}\left(\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}-\lambda \mathrm{I}\right)(h \cdot u) \\
& =h^{-1}\left(\frac{1}{2} h \Delta_{\mathrm{S}^{\mathrm{N}}} u+\nabla h \cdot \nabla u+\frac{1}{2} u \Delta_{\mathrm{S}^{\mathrm{N}}} h-\lambda u \cdot h\right) \\
& =\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}} u+h^{-1} \nabla h \cdot \nabla u \\
& =\frac{1}{2} \Delta_{\mathrm{S}^{\mathfrak{N}}} u+\frac{1}{2}(\mathrm{~N}-2) \tan \frac{1}{2} \theta \frac{\partial}{\partial \theta}
\end{aligned}
$$

where $\nabla$ is the gradient for the Euclidean metric on $S^{N}$. This proves the first assertion and the second follows from Proposition 1.
(Note that the conditioning described above does correspond to the naïve idea of conditioning a process by its position at time T. For suppose that U is a subset of $\mathrm{S}^{\mathrm{N}}$ with a smooth boundary. If $\left(x_{t}\right)$ is the Brownian motion on $S^{\mathrm{N}}$, then we may form a new process

$$
\begin{aligned}
x_{t}^{*} & =x_{t} \quad \text { for } \quad t<\mathrm{T} \\
& =\partial \quad \text { for } \quad t \geqslant \mathrm{~T}
\end{aligned}
$$

which jumps to a coffin state $\partial$ at the random time T . If we condition ( $x_{t}^{*}$ ) so that $x_{\mathrm{T}-}^{*} \in \mathrm{U}$ then we obtain the transition semigroup $\mathrm{P}_{t}$ given by

$$
\begin{aligned}
\mathrm{P}_{t} f(x) & =\mathrm{E}^{x}\left(f\left(x_{t}^{*}\right) \mid x_{\mathrm{T}}^{*}-\in \mathrm{U}\right) \\
& =\mathrm{E}^{x}\left(f\left(x_{t}\right) 1_{(t<\mathrm{T})} \mid x_{\mathrm{T}} \in \mathrm{U}\right) \\
& =\frac{\mathrm{E}^{x}\left(f\left(x_{t}\right) 1_{(t<\mathrm{T})} 1_{\mathrm{U}}\left(x_{\mathrm{T}}\right)\right)}{\mathrm{E}^{x}\left(1_{\mathrm{U}}\left(x_{\mathrm{T}}\right)\right)}
\end{aligned}
$$

Setting

$$
h(x)=\mathrm{E}^{x}\left(1_{\mathrm{U}}\left(x_{\mathrm{T}}\right)\right)
$$

we find that

$$
\begin{aligned}
\mathrm{P}_{t} f(x) & =h(x)^{-1} \mathrm{E}^{x}\left(f\left(x_{t}\right) 1_{(t<\mathrm{T})} h\left(x_{\mathrm{T}}\right)\right) \\
& =h(x)^{-1} \int_{t}^{\infty} \mathrm{E}^{x}\left(f\left(x_{t}\right) h\left(x_{s}\right)\right) \lambda e^{-\lambda s} d s \\
& =h(x)^{-1} e^{-\lambda t} \mathrm{E}^{x}\left(f\left(x_{t}\right) h\left(x_{t}\right)\right)
\end{aligned}
$$

by using the Markov property of the Brownian motion. Thus the condiVol. 21, $\mathrm{n}^{\circ}$ 2-1985.
tioned process is the $h$-transform of the Brownian motion for $h$ the distributional solution of

$$
\left(\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}-\lambda \mathrm{I}\right) h=1_{\mathrm{U}}
$$

We can now decompose this process into an average of the processes conditioned to be at a point $X \in U$ at the time T. See J. L. Doob [1] for further details.)

For each $Y \in \mathbf{S}^{\mathrm{N}}$ let $h(\mathrm{Y},$.$) be the Green's function of \frac{1}{2} \Delta_{\mathbf{S}^{\mathrm{N}}}-\frac{\mathrm{N}(\mathrm{N}-2)}{8} \mathrm{I}$ with a pole at Y . Then the Brownian motion conditioned to be at Y at the negative exponential time T has infinitesimal generator

$$
u \rightarrow h(\mathrm{Y}, x)^{-1}\left(\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}-\frac{\mathrm{N}(\mathrm{~N}-2)}{8} \mathrm{I}\right)(h(\mathrm{Y}, x) u(x))
$$

on $S^{\mathbf{N}} \backslash\{Y\}$. As in Proposition 2 we find that this is

$$
u \rightarrow \frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}} u(x)-(\mathrm{N}-2)\|x-\mathrm{Y}\|^{-1} \nabla\|x-\mathrm{Y}\| . \nabla u(x)
$$

Call this generator $\mathscr{G}_{\mathrm{Y}}$.
Corollary. - Let $\left(x_{t}: 0 \leqslant t \leqslant \mathrm{~S}\right)$ be the process with generator $\mathscr{G}_{\mathrm{P}}$ which starts from Y at time $t=0$ and stops at the time S when it first hits P . Then the time reversed process ( $\tilde{x}_{\tau}: 0 \leqslant \tau \leqslant S$ ) given by

$$
\tilde{x}_{\tau}=x_{\mathbf{S}-\tau}
$$

has infinitesimal generator $\mathscr{G}_{\mathbf{Y}}$, starts from P at $\tau=0$ and stops at the time S when it first hits Y.

Proof. - Since stereographic projection maps $\left(x_{t}\right)$ onto Brownian motion in $\mathrm{R}^{\mathrm{N}}$ it is clear that ( $x_{t}: t>0$ ) almost surely never hits Y . Thus the reversed process certainly starts from $\mathbf{P}$ at $\tau=0$ and stops at the time $\mathbf{S}$ when it first hits Y. It remains to find its infinitesimal generator.

Let $g(\mathrm{Y},$.$) be the Green's function for \mathscr{G}_{\mathrm{P}}$ with pole at Y , then, for any smooth function $f$ which is compactly supported within $\mathrm{S}^{\mathbf{N}} \backslash\{\mathrm{P}, \mathrm{Y}\}$, we have

$$
\mathrm{E} \int_{0}^{\mathrm{s}} f\left(x_{t}\right) d t=\int g(x, \mathrm{Y}) f(x) d \mathrm{~V}(x)=\mathrm{E} \int_{0}^{\mathrm{s}} f\left(\tilde{x}_{\tau}\right) d \tau
$$

where $d \mathrm{~V}$ is the N -dimensional Lebesgue measure on $\mathrm{S}^{\mathrm{N}}$.
Consequently, if we denote by $\mathscr{G}_{\mathrm{P}},\left(\mathrm{P}_{t}\right)$ the generator and transition
semigroup for $\left(x_{t}\right)$ and by $\tilde{\mathscr{G}}_{\mathrm{P}},\left(\tilde{\mathrm{P}}_{\tau}\right)$ the corresponding operators for $\left(\tilde{x}_{\tau}\right)$, then we obtain

$$
\begin{aligned}
\int g(x, \mathrm{X}) f(x) \mathrm{P}_{r} k(x) d \mathrm{~V}(x) & =\mathrm{E} \int_{0}^{\mathrm{s}} f\left(x_{t}\right) \mathrm{P}_{r} k\left(x_{t}\right) d t \\
& =\mathrm{E} \int_{0}^{\mathrm{s}} f\left(x_{t}\right) k\left(x_{t+r}\right) d t \\
& =\mathrm{E} \int_{0}^{\mathrm{s}} f\left(\tilde{x}_{\tau+r}\right) k\left(\tilde{x}_{\tau}\right) d \tau \\
& =\int g(x, \mathrm{Y}) k(x) \tilde{\mathrm{P}}_{r} f(x) d \mathrm{~V}(x) .
\end{aligned}
$$

So

$$
\tilde{\mathrm{P}}_{r} k(x)=g(x, \mathrm{Y})^{-1} \mathrm{P}_{r}^{*}(g(x, \mathrm{Y}) k(x))
$$

and

$$
\tilde{\mathscr{G}}_{\mathrm{P}} k(x)=g(x, \mathrm{Y})^{-1} \mathscr{G}_{\mathrm{P}}^{*}(g(x, \mathrm{Y}) k(x)) .
$$

Now recall that $\mathscr{G}_{\mathrm{P}}=h(\mathrm{P}, .)^{-1}\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right) h(\mathrm{P},$.$) so$

$$
g(x, \mathrm{Y})=\frac{h(\mathrm{Y}, x) h(\mathrm{P}, x)}{h(\mathrm{P}, \mathrm{Y})}
$$

and consequently

$$
\tilde{\mathscr{G}}_{\mathrm{P}} k(x)=h(\mathrm{Y}, x)^{-1}\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right)^{*}(h(\mathrm{Y}, x) k(x)) .
$$

Since the Laplacian is self-adjoint, this gives the desired result.

## 2. CONFORMAL TRANSFORMATIONS

In this section we wish to set the results of $\S 1$ in a more general context.
For any $\lambda>0$ we can condition $\mathrm{BM}\left(\mathrm{S}^{\mathrm{N}}\right)$ to be at P at the independent random time T which has negative exponential distribution with parameter $\lambda$. Indeed, to do so we need only find a positive function $h$ of $\theta$ with

$$
\left(\frac{1}{2} \Delta_{\mathrm{S}^{\mathrm{N}}}-\lambda \mathrm{I}\right) h=0 \quad \text { on } \quad \mathrm{S}^{\mathrm{N}} \backslash\{\mathrm{P}\}
$$

and a singularity at $P$. If we make the change of variables $q=\frac{1}{2}(1-\cos \theta)$ this becomes

$$
q(1-q) \frac{d^{2} h}{d q^{2}}+\frac{1}{2} \mathrm{~N}(1-2 q) \frac{d h}{d q}-2 \lambda h=0
$$

for $0 \leqslant q<1$. This is in the standard hypergeometric form and may be solved by a power series

$$
h=\sum_{n=0}^{\infty} a_{n} q^{n}
$$

This series has radius of convergence 1 and each $a_{n}$ is positive, so $h$ is certainly positive on $0 \leqslant q<1$. For $\lambda \neq \mathrm{N}(\mathbf{N}-2) / 8$ this formula does not define an elementary function. Although the conditioned process may be studied as in the previous section, it does not correspond to a simple process on $\mathbb{R}^{\mathbf{N}}$.

The key property of stereographic projection is that it is conformal so it alters the metric at any point only by a scale factor. We can develop the arguments above for any such conformal transformation.

Proposition 3. - Let M be an N -manifold $(\mathrm{N} \geqslant 3)$ with a Riemannian metric $g_{a b}$ and a conformally equivalent metric

$$
\tilde{g}_{a b}=\Omega^{2} g_{a b} \quad \text { with } \quad \Omega>0 .
$$

Let R and $\tilde{\mathrm{R}}$ be the scalar curvature for $g$ and $\tilde{g}$ respectively. Then the Brownian motion relative to $\tilde{g}$ can be obtained, up to a time change, by conditioning the Brownian motion relative to $g$ according to its behaviour at a negative exponential time if, and only if, $\mathrm{R}-\Omega^{2} \tilde{\mathrm{R}}$ is constant on M .

Proof. - In terms of the infinitesimal generators $\frac{1}{2} \Delta$ and $\frac{1}{2} \tilde{\Delta}$ for the Brownian motions, the Proposition states that there exists $\lambda>0$ and strictly positive functions $h$ and $c$ on M with

$$
\begin{equation*}
\frac{1}{2} \tilde{\Delta} u=c^{2} h^{-1}\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right)(h \cdot u) \tag{1}
\end{equation*}
$$

if, and only if, $R-\Omega^{2} \tilde{R}$ is constant. (If we consider the second degree terms of (1) we see that the condition can only be satisfied if $g$ and $\tilde{g}$ are conformal. So there was no loss of generality in restricting ourselves to this case.)

The proof is simply a standard calculation of the scalar curvature for conformal metrics. We shall use the usual index notation for vectors and tensors on M. Let $\nabla_{a}, \tilde{\nabla}_{a}$ be the covariant derivatives relative to $g$ and $\tilde{g}$.

Then a straightforward but tedious calculation yields the formulae:

$$
\begin{aligned}
\tilde{\nabla}_{a} v_{b} & =\nabla_{a} v_{b}-\Omega^{-1}\left(v_{a} \nabla_{b} \Omega+v_{b} \nabla_{a} \Omega-g_{a b} g^{c d} v_{c} \nabla_{d} \Omega\right) \\
\tilde{\Delta} u & =\tilde{g}^{a b} \tilde{\nabla}_{a} \tilde{\nabla}_{b} u=\tilde{g}^{a b} \tilde{\nabla}_{a}\left(\nabla_{b} u\right) \\
& =\Omega^{2}\left(\Delta u+(\mathrm{N}-2) \Omega^{-1} g^{a b} \nabla_{a} \Omega \nabla_{b} u\right) \\
\Omega^{2} \tilde{\mathrm{R}} & =\mathrm{R}-2(\mathrm{~N}-1) \Omega^{-1} \Delta \Omega-(\mathrm{N}-1)(\mathrm{N}-4) \Omega^{-2} g^{a b} \nabla_{a} \Omega \nabla_{b} \Omega .
\end{aligned}
$$

Thus, for (1) to hold, we must have $c=\Omega$ and

$$
h^{-1}\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right)(h . u)=\frac{1}{2} \Delta u+\frac{1}{2}(\mathrm{~N}-2) \Omega^{-1} g^{a b} \nabla_{a} \Omega \nabla_{b} u .
$$

Now

$$
\Delta(h . u)=g^{a b} \nabla_{a} \nabla_{b}(h . u)=h \Delta u+2 g^{a b} \nabla_{a} h \nabla_{b} u+u \Delta h
$$

so we obtain the two conditions:

$$
h^{-1} \nabla_{a} h=\frac{1}{2}(\mathrm{~N}-2) \Omega^{-1} \nabla_{a} \Omega
$$

and

$$
\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right) h=0 .
$$

The first of these is satisfied if, and only if, $h=\mathrm{K} . \Omega^{\frac{1}{2}(\mathbb{N}-2)}$ for some constant K. In this case, the second condition becomes

$$
\begin{aligned}
0 & =\left(\frac{1}{2} \Delta-\lambda \mathrm{I}\right)\left(\Omega^{\frac{1}{2}(\mathrm{~N}-2)}\right) \\
& =\frac{1}{4}(\mathrm{~N}-2) \Omega^{\frac{1}{2} \mathrm{~N}-2} \Delta \Omega+\frac{1}{8}(\mathrm{~N}-2)(\mathrm{N}-4) \Omega^{\frac{1}{2} \mathrm{~N}-3} g^{a b} \nabla_{a} \Omega \nabla_{b} \Omega-\lambda \Omega^{\frac{1}{2} \mathrm{~N}-1} . \\
\Leftrightarrow \lambda & =\frac{1}{4}(\mathrm{~N}-2) \Omega^{-2} \Delta \Omega+\frac{1}{8}(\mathrm{~N}-2)(\mathrm{N}-4) \Omega^{-2} g^{a b} \nabla_{a} \Omega \nabla_{b} \Omega \\
& =\frac{\mathrm{N}-2}{8(\mathrm{~N}-1)} \cdot\left(\mathrm{R}-\Omega^{2} \tilde{\mathrm{R}}\right) . \quad
\end{aligned}
$$

If we take $g$ to be the Euclidean metric on $S^{N}$ and $\tilde{g}$ the metric on $S^{N}$ which corresponds under stereographic projection to the Euclidean metric on $\mathbf{R}^{\mathbf{N}}$, then

$$
\Omega=\frac{1}{1+\cos \theta}, \quad \mathrm{R}=\mathrm{N}(\mathbf{N}-1), \quad \tilde{\mathrm{R}}=0
$$

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and we recover Proposition 2. The above formula may also be usefully applied to conformal mappings from $S^{\mathrm{N}}$ to itself.

Proposition 4. - Let $\left(x_{t}: 0 \leqslant t<\mathrm{S}\right)$ be the process on $\mathrm{S}^{\mathbf{N}} \backslash\{\mathrm{P}\}$ with infinitesimal generator $\mathscr{G}_{\mathrm{P}}$ and let $\mathrm{T}: \mathrm{S}^{\mathrm{N}} \rightarrow \mathrm{S}^{\mathrm{N}}$ be a conformal automorphism of $\mathrm{S}^{\mathrm{N}}$. Then ( $\mathrm{T} x_{t}: 0 \leqslant t<\mathrm{S}$ ) is a time-changed version of the process on $\mathrm{S}^{\mathbf{N}} \backslash\{\mathrm{TP}\}$ with infinitesimal generator $\mathscr{G}_{\mathrm{TP}}$.

Proof. - Recall that the group of conformal automorphisms of $S^{N}$ is generated by the inversions in spheres orthogonal to $S^{\mathrm{N}}$. We could prove the result by direct calculation, as in $\S 1$, of the effect of such an inversion. However, it is simpler to argue indirectly.

Let $U: R^{N} \rightarrow S^{\mathbf{N}}$ be stereographic projection with centre $P$ and let $\mathrm{V}: \mathrm{R}^{\mathrm{N}} \rightarrow \mathrm{S}^{\mathrm{N}}$ be the stereographic projection with centre TP from the N -dimensional subspace of $\mathrm{R}^{\mathrm{N}+1}$ orthogonal to TP. Both of these maps are conformal, so the composite

$$
\mathrm{Q}=\mathrm{V}^{-1} \mathrm{TU}: \mathrm{R}^{\mathrm{N}} \rightarrow \mathrm{R}^{\mathrm{N}}
$$

is conformal. Since $N \geqslant 3$, the only such conformal maps are the Euclidean similarities of $\mathrm{R}^{\mathrm{N}}$. These similarities obviously preserve Brownian motion on $\mathrm{R}^{\mathrm{N}}$, to within alteration of the time scale by a constant factor. Now Proposition 1 shows that, to within a time change, $U$ maps $B M\left(R^{N}\right)$ to the process with generator $\mathscr{G}_{\mathrm{P}}$ and V maps $\mathrm{BM}\left(\mathrm{R}^{\mathrm{N}}\right)$ to the process with generator $\mathscr{G}_{\mathrm{TP}}$. Therefore, $\mathrm{T}=\mathrm{VQU}^{-1}$ does indeed transform the process with generator $\mathscr{G}_{\mathrm{P}}$ to a time-changed version of the process with generator $\mathscr{G}_{\mathrm{TP}}$.

If we combine Proposition 4 with the earlier Corollary, we see that timereversal of the process starting at Y with generator $\mathscr{G}_{\mathrm{P}}$ corresponds to the image of the process under any inversion which maps $S^{N}$ onto itself and interchanges Y and P . This should be compared with the results of M . Yor [3].

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