

SPATIAL HETEROGENEITY IN 3D-2D DIMENSIONAL REDUCTION

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Abstract. A justification of heterogeneous membrane models as zero-thickness limits of a cylindrical three-dimensional heterogeneous nonlinear hyperelastic body is proposed in the spirit of Le Dret (1995). Specific characterizations of the 2D elastic energy are produced. As a generalization of Bouchitté *et al.* (2002), the case where external loads induce a density of bending moment that produces a Cosserat vector field is also investigated. Throughout, the 3D-2D dimensional reduction is viewed as a problem of Γ -convergence of the elastic energy, as the thickness tends to zero.

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1. INTRODUCTION

The purpose of this article is to study the behavior of a thin elastic plate, as the thickness tends to zero. This approach renders more realistic the idealized view of a film as a thin plate. The originality of the work comes from the heterogeneity of the material under consideration. Previous results have been established in the homogeneous case; our aim here is to generalize those. As we will see, accounting for inhomogeneity leads to technical difficulties linked to the equi-integrable character of the scaled gradient. We will use a “classical” approach of the theory of dimension reduction. In recent years, the investigation of dimensional reduction has focused on variational methods and used De Giorgi’s Γ -convergence (see [5, 8]) as its main tool.

As far as 3D-2D asymptotic analysis is concerned, the seminal paper is [15], in which a membrane model is derived from three-dimensional hyperelasticity. In its footsteps several studies have derived or re-derived various membrane-like models in various settings; see in particular [6] and references therein; note that in Section 3 of that paper, a transversally inhomogeneous thin domain is studied, but that in-plane-homogeneity is imposed. Because of frame indifference, it may occur that the membrane effect is not excited by the loads: this is the case for example when the lateral boundary conditions on the thin domain are compressive (see *e.g.* Th. 6.2 in [13]). Then the membrane energy, which results from a 3D-energy of the order of the thickness ε , is actually zero and lower energy modes are activated. In [11], a justification of classical nonlinear plate models for a homogeneous isotropic material is given by a formal asymptotic expansion. Recently, those results have been rigorously justified by means of variational methods for general homogeneous hyperelastic bodies. A Kirchhoff bending model in [12, 13], and a Föppl-von Kármán model in [14] have been obtained when the 3D-energy scales respectively like ε^3 and ε^5 .

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The present study falls squarely within the membrane framework in the sense that, thanks to frame indifference, the stored energy function depends only on the first fundamental form of the deformed plate mid-surface. Our goal is to rigorously derive models for heterogeneous membranes from their heterogeneous thin 3D-counterparts. The paper is devoted to a generalization of the results established in [3, 6, 15] to the case of a general inhomogeneity.

The key ingredient of this study is the equi-integrability theorem of [2] (Th. 1.1 of that reference). An alternative proof of that theorem was also communicated to the authors [4]. This theorem shows that a sequence of scaled gradients $\{(D_\alpha u_\varepsilon|_{\frac{1}{\varepsilon}D_3u_\varepsilon})\}$, which is bounded in $L^p(\Omega; \mathbb{M}^{3 \times 3})$, with $p > 1$, can be decomposed into the sum of two sequences $\{w_\varepsilon\}$ and $\{z_\varepsilon\}$ where $\{|(D_\alpha w_\varepsilon|_{\frac{1}{\varepsilon}D_3w_\varepsilon})|^p\}$ is equi-integrable and $z_\varepsilon \rightarrow 0$ in measure.

Let ω be a bounded open subset of \mathbb{R}^2 . Consider $\Omega_\varepsilon := \omega \times (-\varepsilon, \varepsilon)$, the reference configuration of a hyperelastic heterogeneous thin film, with elastic energy density given by the ε -dependent Carathéodory function $W_\varepsilon : \Omega_\varepsilon \times \mathbb{M}^{3 \times 3} \rightarrow \mathbb{R}$. We will assume *e.g.* that the body is clamped on the lateral boundary $\Gamma_\varepsilon := \partial\omega \times (-\varepsilon, \varepsilon)$ and that it is submitted to the action of surface traction densities on $\Sigma_\varepsilon := \omega \times \{-\varepsilon, \varepsilon\}$. The total energy of the system under a deformation $u : \Omega_\varepsilon \rightarrow \mathbb{R}^3$ is given by

$$\mathcal{E}(\varepsilon)(u) = \int_{\Omega_\varepsilon} W_\varepsilon(x; Du) dx - \int_{\Omega_\varepsilon} f_\varepsilon \cdot u dx - \int_{\Sigma_\varepsilon} g_\varepsilon \cdot u d\mathcal{H}^2,$$

where \mathcal{H}^2 stands for the 2-dimensional surface measure, $f_\varepsilon \in L^{p'}(\Omega_\varepsilon; \mathbb{R}^3)$ denotes an appropriate dead load and $g_\varepsilon \in L^{p'}(\Sigma_\varepsilon; \mathbb{R}^3)$ some surface traction densities ($1/p + 1/p' = 1$). We denote by $W_{\Gamma_\varepsilon}^{1,p}(\Omega_\varepsilon; \mathbb{R}^3)$ the space of kinematically admissible fields, that is the functions in $W^{1,p}(\Omega_\varepsilon; \mathbb{R}^3)$ with zero trace on Γ_ε . As is classical in hyperelasticity, the equilibrium problem is viewed as the minimization problem

$$\inf_{u-x \in W_{\Gamma_\varepsilon}^{1,p}(\Omega_\varepsilon; \mathbb{R}^3)} \mathcal{E}(\varepsilon)(u).$$

Since the integration domain depends on ε , we reformulate the problem on a fixed domain through a $1/\varepsilon$ -dilatation in the transverse direction x_3 . Let x_α the vector $(x_1, x_2) \in \omega$, we set $v(x_\alpha, \frac{x_3}{\varepsilon}) := u(x_\alpha, x_3)$ and $\mathcal{E}_\varepsilon(v) := \frac{1}{\varepsilon} \mathcal{E}(\varepsilon)(u)$, then

$$\mathcal{E}_\varepsilon(v) = \int_{\Omega} W_\varepsilon \left(x_\varepsilon; D_\alpha v(x) \Big|_{\frac{1}{\varepsilon} D_3 v(x)} \right) dx - \int_{\Omega} f_\varepsilon(x_\varepsilon) \cdot v(x) dx - \frac{1}{\varepsilon} \int_{\Sigma} g_\varepsilon(x_\varepsilon) \cdot v(x) d\mathcal{H}^2,$$

where $x_\varepsilon := (x_\alpha, \varepsilon x_3)$. We set $\Omega := \omega \times (-1, 1)$, $\Sigma := \omega \times \{-1, 1\}$, denote by $D_\alpha v$ the 3×2 matrix of partial derivatives $\frac{\partial v_i}{\partial x_\alpha}$ ($i \in \{1, 2, 3\}$, $\alpha \in \{1, 2\}$) and by $(\overline{F}|z)$, the two first columns of which are those of the matrix $\overline{F} \in \mathbb{M}^{3 \times 2}$, while the last one is the vector $z \in \mathbb{R}^3$. A formal asymptotic expansion in [11] shows that the membrane theory arises if the body forces are of order 1 and the surface loadings are of order ε . We next assume that

$$\begin{cases} W_\varepsilon(x_\alpha, \varepsilon x_3; F) = W(x_\alpha, x_3; F), \\ f_\varepsilon(x_\alpha, \varepsilon x_3) = f(x_\alpha, x_3), \\ g_\varepsilon(x_\alpha, \varepsilon x_3) = g_0(x_\alpha, x_3) + \varepsilon g(x_\alpha, x_3) \end{cases}$$

where $f \in L^{p'}(\Omega; \mathbb{R}^3)$, $g_0, g \in L^{p'}(\Sigma; \mathbb{R}^3)$ and $W : \Omega \times \mathbb{M}^{3 \times 3} \rightarrow \mathbb{R}$ is a Carathéodory function satisfying conditions of p -coercivity and p -growth: for some $0 < \beta' \leq \beta < +\infty$ and some $1 < p < \infty$,

$$\beta'|F|^p \leq W(x; F) \leq \beta(|F|^p + 1), \quad F \in \mathbb{M}^{3 \times 3}, \quad \text{for a.e. } x \in \Omega. \quad (1.1)$$

The usual Euclidian norm on the space $\mathbb{M}^{N \times m}$ of real $N \times m$ matrices is denoted by $|F|$. The minimization problem becomes

$$\inf_{v-x_\varepsilon \in W_{\Gamma}^{1,p}(\Omega; \mathbb{R}^3)} \mathcal{E}_\varepsilon(v), \quad (1.2)$$

where $W_{\Gamma}^{1,p}(\Omega; \mathbb{R}^3)$ stands for the functions in $W^{1,p}(\Omega; \mathbb{R}^3)$ with zero trace on the lateral boundary $\Gamma := \partial\omega \times (-1, 1)$.

If we denote by g_0^{\pm} (resp. g^{\pm}) the trace of g_0 (resp. g) on $\omega \times \{\pm 1\}$, in view of Remark 3.2.3 of [11], the loading vectors g_0^+ and g_0^- must satisfy $g_0^+ + g_0^- = 0$. In the second section, we assume the stronger condition that $g_0^+ = g_0^- = 0$. The physical implication of this assumption is that the plate of thickness 2ε cannot support a non vanishing resultant surface load as the thickness ε goes to zero. We generalize here the result of [6, 15] to a general inhomogeneity. In the third section, we address the general case of admissible surface loadings. It deals with a similar problem, in which the class of surface forces generates a bending moment density as in [3]; the limit behavior is not solely characterized by the limit deformations (a \mathbb{R}^3 -valued field defined on the mid-plane), but it also involves the average of the Cosserat vector also defined on the mid-plane. Once again, we generalize the result of [3] to the inhomogeneous case. The fourth and last section demonstrates that the classical membrane model can be seen as a particular case of the Cosserat model when the bending moment density is zero.

As for notation, $\mathcal{A}(\omega)$ is the family of open subsets of ω ; \mathcal{L}^N stands for the N -dimensional Lebesgue measure in \mathbb{R}^N ($N = 2$ or 3); \rightarrow always denotes strong convergence whereas \rightharpoonup (resp. $\overset{*}{\rightharpoonup}$) denotes weak (resp. weak-*) convergence. Finally, we loosely identify $L^p(\omega; \mathbb{R}^3)$ (resp. $W^{1,p}(\omega; \mathbb{R}^3)$) with those functions in $L^p(\Omega; \mathbb{R}^3)$ (resp. $W^{1,p}(\Omega; \mathbb{R}^3)$) that do not depend upon x_3 .

2. CLASSICAL NONLINEAR MEMBRANE MODEL

In this section, we assume that $g_{\varepsilon} = \varepsilon g$ with $g \in L^p(\Sigma; \mathbb{R}^3)$. Thus, the minimization problem (1.2) becomes

$$\inf_{v-x_{\varepsilon} \in W_{\Gamma}^{1,p}(\Omega; \mathbb{R}^3)} \left\{ \int_{\Omega} W \left(x; D_{\alpha} v \middle| \frac{1}{\varepsilon} D_3 v \right) dx - \int_{\Omega} f \cdot v dx - \int_{\Sigma} g \cdot v d\mathcal{H}^2 \right\}.$$

Define for any $(u; A) \in L^p(\Omega; \mathbb{R}^3) \times \mathcal{A}(\omega)$,

$$J_{\varepsilon}(u; A) := \begin{cases} \int_{A \times (-1, 1)} W \left(x; D_{\alpha} u \middle| \frac{1}{\varepsilon} D_3 u \right) dx & \text{if } u \in W^{1,p}(A \times (-1, 1); \mathbb{R}^3), \\ +\infty & \text{otherwise,} \end{cases}$$

and

$$J_{\{\varepsilon\}}(u; A) := \inf_{\{u_{\varepsilon}\}} \left\{ \liminf_{\varepsilon \rightarrow 0} J_{\varepsilon}(u_{\varepsilon}; A) : u_{\varepsilon} \rightarrow u \text{ in } L^p(A \times (-1, 1); \mathbb{R}^3) \right\}. \quad (2.1)$$

Remark 2.1. For any $A \in \mathcal{A}(\omega)$, $J_{\{\varepsilon\}}(u; A) = +\infty$ whenever $u \in L^p(\Omega; \mathbb{R}^3) \setminus W^{1,p}(A; \mathbb{R}^3)$, as is easily seen in view of the definition of J_{ε} , together with the coercivity condition (1.1).

By virtue of Remark 2.1, together with Theorem 2.5 in [6], for all sequences $\{\varepsilon\}$, there exists a subsequence $\{\varepsilon_n\}$ such that $J_{\{\varepsilon_n\}}(\cdot; A)$ defined in (2.1) is the $\Gamma(L^p)$ -limit of $J_{\varepsilon_n}(\cdot; A)$. Further, there exists a Carathéodory function $W_{\{\varepsilon_n\}} : \omega \times \mathbb{M}^{3 \times 2} \rightarrow \mathbb{R}$ such that, for all $A \in \mathcal{A}(\omega)$ and all $u \in W^{1,p}(A; \mathbb{R}^3)$

$$J_{\{\varepsilon_n\}}(u; A) = 2 \int_A W_{\{\varepsilon_n\}}(x_{\alpha}; D_{\alpha} u) dx_{\alpha}.$$

Remark 2.2. Lemma 2.6 of [6] implies that $J_{\{\varepsilon_n\}}(u; A)$ is unchanged if the approximating sequences $\{u_{\varepsilon_n}\}$ are constrained to match the lateral boundary condition of their target, i.e. $u_{\varepsilon_n} \equiv u$ on $\partial A \times (-1, 1)$.

From now onward, we will assume that $\{\varepsilon_n\}$ denotes a subsequence of $\{\varepsilon\}$ such that the $\Gamma(L^p)$ -limit of $J_{\varepsilon_n}(u; A)$ exists, in which case it coincides with $J_{\{\varepsilon_n\}}(u; A)$. Under the hypothesis that W is a homogeneous elastic energy density, it is proved in [15], Theorem 2, that $J_{\{\varepsilon_n\}}(u; A)$ does not depend upon the choice

of the sequence $\{\varepsilon_n\}$. It is given by

$$J_{\{\varepsilon_n\}}(u; A) = 2 \int_A \mathcal{Q}\overline{W}(D_\alpha u) \, dx_\alpha,$$

with for all $\overline{F} \in \mathbb{M}^{3 \times 2}$,

$$\overline{W}(\overline{F}) := \inf_{z \in \mathbb{R}^3} W(\overline{F}|z),$$

and,

$$\mathcal{Q}\overline{W}(\overline{F}) := \inf_{\phi \in W_0^{1,p}(Q'; \mathbb{R}^3)} \int_{Q'} \overline{W}(\overline{F} + D_\alpha \phi) \, dx_\alpha,$$

where $Q' := (0, 1)^2$, and $\mathcal{Q}\overline{W}$ is the 2D-quasiconvexification of \overline{W} . This result was extended to the case where W is also function of x_3 in [6], Theorem 3.4. It is proved there that, in such a case, $J_{\{\varepsilon_n\}}$ is given by

$$J_{\{\varepsilon_n\}}(u; A) = 2 \int_A \underline{W}(D_\alpha u) \, dx_\alpha,$$

with for all $\overline{F} \in \mathbb{M}^{3 \times 2}$,

$$\begin{aligned} \underline{W}(\overline{F}) := \inf_{L, \phi} & \left\{ \frac{1}{2} \int_{Q' \times (-1, 1)} W(x_3; \overline{F} + D_\alpha \phi | LD_3 \phi) \, dx_\alpha dx_3 : L > 0, \right. \\ & \left. \phi \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3), \phi = 0 \text{ on } \partial Q' \times (-1, 1) \right\}. \end{aligned}$$

We wish to extend those results to the case where W is a function of both x_3 and x_α . We set, for all $\overline{F} \in \mathbb{M}^{3 \times 2}$ and for a.e. $x_0 \in \omega$,

$$\begin{aligned} \underline{W}(x_0; \overline{F}) := \inf_{L, \phi} & \left\{ \frac{1}{2} \int_{Q' \times (-1, 1)} W(x_0, x_3; \overline{F} + D_\alpha \phi | LD_3 \phi) \, dx_\alpha dx_3 : L > 0, \right. \\ & \left. \phi \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3), \phi = 0 \text{ on } \partial Q' \times (-1, 1) \right\}. \end{aligned} \quad (2.2)$$

The following theorem holds:

Theorem 2.3. *For all $\overline{F} \in \mathbb{M}^{3 \times 2}$ and for a.e. $x_0 \in \omega$, $W_{\{\varepsilon_n\}}(x_0; \overline{F}) = \underline{W}(x_0; \overline{F})$ where \underline{W} is given by (2.2). Furthermore, for any $A \in \mathcal{A}(\omega)$, $J_\varepsilon(\cdot; A) \Gamma(L^p)$ -converges to $J_{\{\varepsilon_n\}}(\cdot; A)$ and for all $u \in L^p(\Omega; \mathbb{R}^3)$,*

$$J_{\{\varepsilon_n\}}(u; A) = \begin{cases} 2 \int_A \underline{W}(x_\alpha; D_\alpha u) \, dx_\alpha & \text{if } u \in W^{1,p}(A; \mathbb{R}^3), \\ +\infty & \text{otherwise.} \end{cases}$$

The proof of this theorem is a direct consequence of Lemmata 2.4, 2.5 below.

Lemma 2.4. *For all $\overline{F} \in \mathbb{M}^{3 \times 2}$ and for a.e. $x_0 \in \omega$,*

$$W_{\{\varepsilon_n\}}(x_0, \overline{F}) \geq \underline{W}(x_0, \overline{F}).$$

Proof. Let us fix $\overline{F} \in \mathbb{M}^{3 \times 2}$, we set $u(x_\alpha) := \overline{F}x_\alpha$ and let x_0 be a Lebesgue point of both $W_{\{\varepsilon_n\}}(\cdot; \overline{F})$ and $\underline{W}(\cdot; \overline{F})$. We denote by $Q'(x_0, r)$, the cube of \mathbb{R}^2 of center x_0 and side length r , where $r > 0$ is fixed and small enough so

that $Q'(x_0, r) \in \mathcal{A}(\omega)$. According to the equi-integrability theorem (Th. 1.1 in [2]), there exists a subsequence of $\{\varepsilon_n\}$ (not relabelled) and a sequence $\{u_n\} \subset W^{1,p}(Q'(x_0, r) \times (-1, 1); \mathbb{R}^3)$ such that

$$\left\{ \begin{array}{l} u_n \rightarrow 0 \text{ in } L^p(Q'(x_0, r) \times (-1, 1); \mathbb{R}^3), \\ \left\{ \left| \left(D_\alpha u_n \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n \right) \right|^p \right\} \text{ is equi-integrable,} \\ J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) = \lim_{n \rightarrow +\infty} \int_{Q'(x_0, r) \times (-1, 1)} W \left(x_\alpha, x_3; \bar{F} + D_\alpha u_n \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n \right) dx_\alpha dx_3. \end{array} \right.$$

Set

$$F_n(x) := \left(\bar{F} + D_\alpha u_n(x) \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n(x) \right).$$

For any $h \in \mathbb{N}$, we cover $Q'(x_0, r)$ with h^2 disjoint cubes $Q'_{i,h}$ of side length r/h . Thus $Q'(x_0, r) = \bigcup_{i=1}^{h^2} Q'_{i,h}$ and

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) = \left(\limsup_{h \rightarrow +\infty} \right) \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{Q'_{i,h} \times (-1, 1)} W(x; F_n(x)) dx. \quad (2.3)$$

Since W is a Carathéodory integrand, Scorza-Dragoni's theorem (see [9], Chap. VIII) implies the existence, for any $\eta > 0$, of a compact set $K_\eta \subset \Omega$ such that

$$\mathcal{L}^3(\Omega \setminus K_\eta) < \eta, \quad (2.4)$$

and the restriction of W to $K_\eta \times \mathbb{M}^{3 \times 3}$ is continuous. For any $\lambda > 0$, define

$$R_n^\lambda := \{x \in Q'(x_0, r) \times (-1, 1) : |F_n(x)| \leq \lambda\}.$$

By virtue of Chebyshev's inequality, there exists a constant $C > 0$ – which does not depend on n or λ – such that

$$\mathcal{L}^3([Q'(x_0, r) \times (-1, 1)] \setminus R_n^\lambda) < \frac{C}{\lambda^p}. \quad (2.5)$$

Denoting by $W^{\eta, \lambda}$ the continuous extension of W outside $K_\eta \times \bar{B}(0, \lambda)$ (defined *e.g.* in Th. 1, Sect. 1.2 in [10]), $W^{\eta, \lambda}$ is continuous on $\mathbb{R}^3 \times \mathbb{M}^{3 \times 3}$ and satisfies the following bound

$$0 \leq W^{\eta, \lambda}(x; F) \leq \max_{K_\eta \times \bar{B}(0, \lambda)} W \leq \beta(1 + \lambda^p) \quad \text{for all } (x; F) \in \mathbb{R}^3 \times \mathbb{M}^{3 \times 3}. \quad (2.6)$$

In view of (2.3), we have

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{[Q'_{i,h} \times (-1, 1)] \cap R_n^\lambda \cap K_\eta} W^{\eta, \lambda}(x; F_n(x)) dx.$$

By virtue of (2.6) and (2.4),

$$\sum_{i=1}^{h^2} \int_{[Q'_{i,h} \times (-1, 1)] \cap R_n^\lambda \setminus K_\eta} W^{\eta, \lambda}(x; F_n(x)) dx \leq \beta(1 + \lambda^p) \eta \xrightarrow{\eta \rightarrow 0} 0,$$

uniformly in (n, h) . Therefore

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{[Q'_{i,h} \times (-1, 1)] \cap R_n^\lambda} W^{\eta, \lambda}(x; F_n(x)) dx.$$

Since $W^{\eta,\lambda}$ is continuous, it is uniformly continuous on $\overline{\Omega} \times \overline{B}(0, \lambda)$. Thus there exists a continuous and increasing function $\omega_{\eta,\lambda} : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ satisfying $\omega_{\eta,\lambda}(0) = 0$ and such that

$$|W^{\eta,\lambda}(x; F) - W^{\eta,\lambda}(y; G)| \leq \omega_{\eta,\lambda}(|x - y| + |F - G|), \quad \forall (x; F), (y; G) \in \overline{\Omega} \times \overline{B}(0, \lambda). \quad (2.7)$$

Consequently, for all $(x_\alpha, x_3) \in [Q'_{i,h} \times (-1, 1)] \cap R_n^\lambda$ and all $y_\alpha \in Q'_{i,h}$,

$$|W^{\eta,\lambda}(x_\alpha, x_3; F_n(x_\alpha, x_3)) - W^{\eta,\lambda}(y_\alpha, x_3; F_n(x_\alpha, x_3))| \leq \omega_{\eta,\lambda}(|x_\alpha - y_\alpha|) \leq \omega_{\eta,\lambda}(c/h).$$

We get, after integration in (x, y_α) and summation,

$$\begin{aligned} \sup_{n \in \mathbb{N}} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{R_n^\lambda \cap [Q'_{i,h} \times (-1, 1)]} |W^{\eta,\lambda}(y_\alpha, x_3; F_n(x)) - W^{\eta,\lambda}(x_\alpha, x_3; F_n(x))| dx \right\} dy_\alpha \\ \leq 2r^2 \omega_{\eta,\lambda}(c/h) \xrightarrow{h \rightarrow +\infty} 0. \end{aligned}$$

Hence,

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{[Q'_{i,h} \times (-1, 1)] \cap R_n^\lambda} W^{\eta,\lambda}(y_\alpha, x_3; F_n(x)) dx \right\} dy_\alpha.$$

Define the following sets which depend on all parameters (η, λ, i, h, n) :

$$\begin{aligned} E &:= \{(y_\alpha, x_\alpha, x_3) \in Q'_{i,h} \times Q'_{i,h} \times (-1, 1) : (y_\alpha, x_3) \in K_\eta \text{ and } (x_\alpha, x_3) \in R_n^\lambda\}, \\ E_1 &:= \{(y_\alpha, x_\alpha, x_3) \in Q'_{i,h} \times Q'_{i,h} \times (-1, 1) : (y_\alpha, x_3) \notin K_\eta \text{ and } (x_\alpha, x_3) \in R_n^\lambda\}, \\ E_2 &:= \{(y_\alpha, x_\alpha, x_3) \in Q'_{i,h} \times Q'_{i,h} \times (-1, 1) : (x_\alpha, x_3) \notin R_n^\lambda\}, \end{aligned}$$

and note that $Q'_{i,h} \times Q'_{i,h} \times (-1, 1) = E \cup E_1 \cup E_2$. Since W and $W^{\eta,\lambda}$ coincide on $K_\eta \times \overline{B}(0, \lambda)$,

$$\begin{aligned} J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) &\geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_E W^{\eta,\lambda}(y_\alpha, x_3; F_n(x)) dx dy_\alpha \\ &= \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_E W(y_\alpha, x_3; F_n(x)) dx dy_\alpha. \end{aligned} \quad (2.8)$$

We will prove that the corresponding terms over E_1 and E_2 are zero. Indeed, in view of (2.4) and the p -growth condition (1.1),

$$\begin{aligned} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{E_1} W(y_\alpha, x_3; F_n(x)) dx dy_\alpha &\leq \sum_{i=1}^{h^2} \frac{h^2}{r^2} \mathcal{L}^2(Q'_{i,h}) \mathcal{L}^3([Q'_{i,h} \times (-1, 1)] \setminus K_\eta) \beta(1 + \lambda^p) \\ &= \beta(1 + \lambda^p) \mathcal{L}^3([Q'(x_0, r) \times (-1, 1)] \setminus K_\eta) \\ &< \beta(1 + \lambda^p) \eta \xrightarrow{\eta \rightarrow 0} 0, \end{aligned} \quad (2.9)$$

uniformly in (n, h) . The bound from above in (1.1), the equi-integrability of $\{|F_n|^p\}$ and (2.5) imply that

$$\begin{aligned} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{E_2} W(y_\alpha, x_3; F_n(x)) \, dx dy_\alpha &\leq \sum_{i=1}^{h^2} \frac{h^2}{r^2} \mathcal{L}^2(Q'_{i,h}) \beta \int_{[Q'_{i,h} \times (-1,1)] \setminus R_n^\lambda} (1 + |F_n(x)|^p) \, dx \\ &= \beta \int_{[Q'(x_0,r) \times (-1,1)] \setminus R_n^\lambda} (1 + |F_n(x)|^p) \, dx \xrightarrow{\lambda \rightarrow +\infty} 0, \end{aligned} \tag{2.10}$$

uniformly in (η, n, h) . Thus, in view of (2.8), (2.9), (2.10), Fatou's lemma yields

$$\begin{aligned} J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) &\geq \limsup_{h \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_n(x)) \, dx \right\} dy_\alpha \\ &\geq \limsup_{h \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_n(x)) \, dx \right\} dy_\alpha \\ &\geq \limsup_{h \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \liminf_{n \rightarrow +\infty} \int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_n(x)) \, dx \right\} dy_\alpha. \end{aligned}$$

We apply, for a.e. $y_\alpha \in Q'_{i,h}$, Theorem 3.1 in [6] to the Carathéodory function $(x_3; F) \mapsto W(y_\alpha, x_3; F)$; in particular

$$\liminf_{n \rightarrow +\infty} \int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_n(x)) \, dx \geq \frac{2r^2}{h^2} \underline{W}(y_\alpha; \overline{F}).$$

Thus

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \geq \limsup_{h \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \frac{2r^2}{h^2} \underline{W}(y_\alpha; \overline{F}) \, dy_\alpha = 2 \int_{Q'(x_0,r)} \underline{W}(y_\alpha; \overline{F}) \, dy_\alpha.$$

Dividing both sides of the previous inequality by r^2 and passing to the limit when $r \searrow 0^+$, we obtain

$$W_{\{\varepsilon_n\}}(x_0; \overline{F}) \geq \underline{W}(x_0; \overline{F}). \quad \square$$

Lemma 2.5. For all $\overline{F} \in \mathbb{M}^{3 \times 2}$ and for a.e. $x_0 \in \omega$,

$$W_{\{\varepsilon_n\}}(x_0; \overline{F}) \leq \underline{W}(x_0; \overline{F}).$$

Proof. For all $k \geq 1$, let $L_k > 0$ and $\varphi_k \in W^{1,\infty}(Q' \times (-1,1); \mathbb{R}^3)$ with $\varphi_k = 0$ on $\partial Q' \times (-1,1)$ be such that

$$Z_k(x_0; \overline{F}) := \frac{1}{2} \int_{Q' \times (-1,1)} W(x_0, x_3; \overline{F} + D_\alpha \varphi_k | L_k D_3 \varphi_k) \, dx_\alpha dx_3 \leq \underline{W}(x_0; \overline{F}) + \frac{1}{k}. \tag{2.11}$$

This is legitimate because of the density of $W^{1,\infty}(Q' \times (-1,1); \mathbb{R}^3)$ into $W^{1,p}(Q' \times (-1,1); \mathbb{R}^3)$ and the p -growth condition (1.1). We extend φ_k to $\mathbb{R}^2 \times (-1,1)$ by Q' -periodicity and set $F_k(x) := (\overline{F} + D_\alpha \varphi_k(x) | L_k D_3 \varphi_k(x))$. Then, there exists $M_k > 0$ such that

$$\|F_k\|_{L^\infty(\mathbb{R}^2 \times (-1,1); \mathbb{R}^3)} \leq M_k. \tag{2.12}$$

Let $\overline{F} \in \mathbb{M}^{3 \times 2}$ and x_0 be a Lebesgue point of $\underline{W}(\cdot; \overline{F})$ and $Z_k(\cdot; \overline{F})$ for all $k \geq 1$. We choose $r > 0$ small enough such that $Q'(x_0, r) \in \mathcal{A}(\omega)$. Fix $k \geq 1$ and set

$$\begin{cases} u(x_\alpha) := \overline{F} x_\alpha, \\ u_n^k(x_\alpha, x_3) := \overline{F} x_\alpha + L_k \varepsilon_n \varphi_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right). \end{cases}$$

Since $u_n^k \xrightarrow{n \rightarrow +\infty} u$ in $L^p(Q'(x_0, r) \times (-1, 1); \mathbb{R}^3)$,

$$\begin{aligned} J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) &\leq \liminf_{n \rightarrow +\infty} \int_{Q'(x_0, r) \times (-1, 1)} W \left(x_\alpha, x_3; D_\alpha u_n^k \Big| \frac{1}{\varepsilon_n} D_3 u_n^k \right) dx_\alpha x_3 \\ &= \liminf_{n \rightarrow +\infty} \int_{Q'(x_0, r) \times (-1, 1)} W \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) dx_\alpha x_3. \end{aligned}$$

As before, we split $Q'(x_0, r)$ into h^2 disjoint cubes $Q'_{i,h}$ of length r/h . Then,

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \leq \left(\liminf_{h \rightarrow +\infty} \right) \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{Q'_{i,h} \times (-1, 1)} W \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) dx_\alpha x_3.$$

Let K_η be like in Lemma 2.4 and $W^{\eta,k}$ be a continuous extension of W outside $K_\eta \times \overline{B}(0, M_k)$ which satisfies the analogue of (2.6) with M_k instead of λ . In view of the p -growth condition (1.1), (2.12) and (2.4), we get

$$\sup_{n \in \mathbb{N}} \sum_{i=1}^{h^2} \int_{[Q'_{i,h} \times (-1, 1)] \setminus K_\eta} W \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) dx_\alpha x_3 \leq \beta(1 + M_k^p) \eta \xrightarrow{\eta \rightarrow 0} 0.$$

Thus,

$$\begin{aligned} J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) &\leq \liminf_{\eta \rightarrow 0} \liminf_{h \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{[Q'_{i,h} \times (-1, 1)] \cap K_\eta} W^{\eta,k} \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) dx_\alpha x_3 \\ &\leq \liminf_{\eta \rightarrow 0} \liminf_{h \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \int_{Q'_{i,h} \times (-1, 1)} W^{\eta,k} \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) dx_\alpha x_3. \end{aligned}$$

Since $W^{\eta,k}$ is continuous, it is uniformly continuous on $\overline{\Omega} \times \overline{B}(0, M_k)$. Thus, there exists a continuous and increasing function $\omega_{\eta,k} : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ satisfying $\omega_{\eta,k}(0) = 0$ and the analogue of (2.7), replacing λ by M_k . Then, for every $(x_\alpha, x_3) \in Q'_{i,h} \times (-1, 1)$ and every $y_\alpha \in Q'_{i,h}$,

$$\begin{aligned} \left| W^{\eta,k} \left(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) - W^{\eta,k} \left(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \right) \right| &\leq \omega_{\eta,k}(|x_\alpha - y_\alpha|) \\ &\leq \omega_{\eta,k}(c/h). \end{aligned}$$

Integration and summation yield in turn

$$\begin{aligned} \sup_{n \in \mathbb{N}} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{Q'_{i,h} \times (-1, 1)} \left| W^{\eta,k}(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) - W^{\eta,k}(x_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) \right| dx \right\} dy_\alpha \\ \leq 2r^2 \omega_{\eta,k}(c/h) \xrightarrow{h \rightarrow +\infty} 0. \end{aligned}$$

Hence,

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \leq \liminf_{\eta \rightarrow 0} \liminf_{h \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{Q'_{i,h} \times (-1, 1)} W^{\eta,k}(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) dx_\alpha dx_3 \right\} dy_\alpha.$$

According to (2.6) and (2.4),

$$\sup_{n \in \mathbb{N}} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{[Q'_{i,h} \times (-1,1)] \cap K_\eta} W^{\eta,k}(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) dy_\alpha dx_3 \right\} dx_\alpha \leq \beta(1 + M_k^p) \eta \xrightarrow{\eta \rightarrow 0} 0.$$

Since $W^{\eta,k}$ coincides with W on $K_\eta \times \overline{B}(0, M_k)$, we get

$$\begin{aligned} J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) &\leq \liminf_{\eta \rightarrow 0} \liminf_{h \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \left\{ \int_{[Q'_{i,h} \times (-1,1)] \cap K_\eta} W(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) dy_\alpha dx_3 \right\} dx_\alpha \\ &\leq \liminf_{h \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \limsup_{n \rightarrow +\infty} \int_{Q'_{i,h}} \left\{ \int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right)) dy_\alpha dx_3 \right\} dx_\alpha. \end{aligned}$$

Riemann-Lebesgue's lemma applied to the Q' -periodic function $\int_{Q'_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_k(\cdot, x_3)) dy_\alpha dx_3$ implies that

$$J_{\{\varepsilon_n\}}(u; Q'(x_0, r)) \leq \liminf_{h \rightarrow +\infty} \sum_{i=1}^{h^2} \frac{h^2}{r^2} \int_{Q'_{i,h}} \frac{2r^2}{h^2} Z_k(y_\alpha; \overline{F}) dy_\alpha = 2 \int_{Q'(x_0, r)} Z_k(y_\alpha; \overline{F}) dy_\alpha.$$

Dividing both sides of the inequality by r^2 and letting $r \searrow 0^+$, we get in view of the definition of x_0 and (2.11),

$$W_{\{\varepsilon_n\}}(x_0; \overline{F}) \leq Z_k(x_0; \overline{F}) \leq \underline{W}(x_0; \overline{F}) + \frac{1}{k}.$$

Passing to the limit when $k \nearrow +\infty$ yields the desired result. \square

Proof of Theorem 2.3. For a.e. $x_0 \in \omega$ and for all $\overline{F} \in \mathbb{M}^{3 \times 2}$, $W_{\{\varepsilon_n\}}(x_0; \overline{F}) = \underline{W}(x_0; \overline{F})$. Since the $\Gamma(L^p)$ -limit does not depend upon the choice of sequence $\{\varepsilon_n\}$, appealing to Proposition 7.11 in [5] we conclude that for any $A \in \mathcal{A}(\omega)$, the whole sequence $J_\varepsilon(\cdot; A)$ $\Gamma(L^p)$ -converges to $J_{\{\varepsilon\}}(\cdot; A)$ and we have,

$$J_{\{\varepsilon\}}(u; A) = 2 \int_A \underline{W}(x_\alpha; D_\alpha u) dx_\alpha,$$

for all $u \in W^{1,p}(A; \mathbb{R}^3)$. \square

Remark 2.6. Proposition 4.1 gives another expression for the energy density \underline{W} .

Remark 2.7. By construction and thanks to Remark 3.3 of [6], Theorem 2.3 generalizes both Theorem 2 of [15] and Theorem 3.4 of [6].

3. COSSERAT NONLINEAR MEMBRANE MODEL

In this section, we assume as in [3] that $g_\varepsilon := g_0 + \varepsilon g$ with $g_0, g \in L^{p'}(\Sigma; \mathbb{R}^3)$ and $g_0^+ + g_0^- = 0$. Thus, the minimization problem (1.2) reads as

$$\inf_{v - x_\varepsilon \in W_{\Gamma}^{1,p}(\Omega; \mathbb{R}^3)} \left\{ \int_{\Omega} W \left(x; D_\alpha v \middle| \frac{1}{\varepsilon} D_3 v \right) dx - L_\varepsilon(v) \right\},$$

with

$$L_\varepsilon(v) := \int_{\Omega} f \cdot v dx + \int_{\Sigma} g \cdot v d\mathcal{H}^2 + \int_{\omega} g_0^\pm \cdot \left(\frac{v^+ - v^-}{\varepsilon} \right) dx_\alpha, \quad v^\pm(x_\alpha) := v(x_\alpha, \pm 1).$$

If $v_\varepsilon \rightarrow v$ in $L^p(\Omega; \mathbb{R}^3)$ is a minimizing sequence and if $b_\varepsilon := \frac{1}{\varepsilon} D_3 v_\varepsilon$, then

$$L_\varepsilon(v_\varepsilon) = \int_\Omega f \cdot v_\varepsilon \, dx + \int_\Sigma g \cdot v_\varepsilon \, d\mathcal{H}^2 + 2 \int_\omega g_0^+ \cdot \bar{b}_\varepsilon \, dx_\alpha, \quad \text{where } \bar{b}_\varepsilon = \frac{1}{2} \int_{-1}^1 b_\varepsilon(\cdot, x_3) \, dx_3.$$

By virtue of the coercivity condition (1.1), we deduce that the sequence $\{v_\varepsilon\}$ is uniformly bounded in $W^{1,p}(\Omega; \mathbb{R}^3)$ and that, for a subsequence of $\{\varepsilon\}$ still labelled $\{\varepsilon\}$, $v_\varepsilon \rightarrow v$ in $W^{1,p}(\Omega; \mathbb{R}^3)$ and $b_\varepsilon \rightarrow b$ in $L^p(\Omega; \mathbb{R}^3)$ with $v \in W^{1,p}(\omega; \mathbb{R}^3)$. As previously, v is associated to the mid-plane deformation, whereas b is the Cosserat vector. In any case, $\lim_{\varepsilon \rightarrow 0} L_\varepsilon(v_\varepsilon) = L(v, \bar{b})$, with

$$L(v, \bar{b}) := \int_\omega (2\bar{f} + g^+ + g^-) \cdot v \, dx_\alpha + 2 \int_\omega g_0^+ \cdot \bar{b} \, dx_\alpha, \quad (3.1)$$

where $\bar{b}(x_\alpha) := \frac{1}{2} \int_{-1}^1 b(x_\alpha, x_3) \, dx_3$ and $\bar{f}(x_\alpha) := \frac{1}{2} \int_{-1}^1 f(x_\alpha, x_3) \, dx_3$. The desired membrane model should thus depend on the average, \bar{b} , of b with respect to x_3 . Once we establish our Γ -convergence result, we will be in a position to conclude that v and \bar{b} are truly independent and that the corresponding model is a Cosserat type membrane model.

To this end, we define, for all $(u, \bar{b}; A) \in L^p(\Omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3) \times \mathcal{A}(\omega)$,

$$\mathcal{J}_\varepsilon(u, \bar{b}; A) := \begin{cases} \int_{A \times (-1,1)} W\left(x; D_\alpha u \Big| \frac{1}{\varepsilon} D_3 u\right) \, dx & \text{if } \begin{cases} u \in W^{1,p}(A \times (-1,1); \mathbb{R}^3), \\ \frac{1}{2\varepsilon} \int_{-1}^1 D_3 u(\cdot, x_3) \, dx_3 = \bar{b}, \end{cases} \\ +\infty & \text{otherwise,} \end{cases} \quad (3.2)$$

and

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) := \inf_{\{u_\varepsilon, \bar{b}_\varepsilon\}} \left\{ \liminf_{\varepsilon \rightarrow 0} \mathcal{J}_\varepsilon(u_\varepsilon, \bar{b}_\varepsilon; A) : u_\varepsilon \rightarrow u \text{ in } L^p(A \times (-1,1); \mathbb{R}^3) \text{ and } \bar{b}_\varepsilon \rightarrow \bar{b} \text{ in } L^p(A; \mathbb{R}^3) \right\}. \quad (3.3)$$

Remark 3.1. Let $(u, \bar{b}; A) \in L^p(\Omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3) \times \mathcal{A}(\omega)$ and suppose that $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) < +\infty$. Arguing as in Remark 2.1, we deduce that $u \in W^{1,p}(A; \mathbb{R}^3)$. Hence, if $u \in L^p(\Omega; \mathbb{R}^3) \setminus W^{1,p}(A; \mathbb{R}^3)$, then $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = +\infty$.

Remark 3.2. Whenever $u \in W^{1,p}(A; \mathbb{R}^3)$, one has $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) < +\infty$, this is easily obtained by considering the sequence $\{u_\varepsilon(x_\alpha, x_3) := u(x_\alpha) + \varepsilon x_3 \bar{b}_\varepsilon(x_\alpha)\}$, where $\bar{b}_\varepsilon \in C_c^\infty(A; \mathbb{R}^3)$ and $\bar{b}_\varepsilon \rightarrow \bar{b}$ strongly in $L^p(A; \mathbb{R}^3)$.

Theorem 1.2 in [3] shows that, if W is a homogeneous elastic energy density, then $\mathcal{J}_{\{\varepsilon\}}$ is the $\Gamma(L^p)$ -limit of \mathcal{J}_ε , by which we mean, from now onward, the Γ -limit with respect to, respectively, the strong topology of $L^p(\Omega; \mathbb{R}^3)$, and the weak topology of $L^p(\omega; \mathbb{R}^3)$. Furthermore, for all $(u, \bar{b}; A) \in W^{1,p}(\omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3) \times \mathcal{A}(\omega)$,

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = 2 \int_A \mathcal{Q}^* W(D_\alpha u | \bar{b}) \, dx_\alpha,$$

where, for all $\bar{F} \in \mathbb{M}^{3 \times 2}$ and $z \in \mathbb{R}^3$,

$$\mathcal{Q}^* W(\bar{F} | z) := \inf_{L > 0, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1,1)} W(\bar{F} + D_\alpha \varphi | LD_3 \varphi) \, dx_\alpha \, dx_3 : \varphi \in W^{1,p}(Q' \times (-1,1); \mathbb{R}^3), \right. \\ \left. \varphi(\cdot, x_3) \text{ } Q' \text{-periodic for a.e. } x_3 \in (-1,1), \frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi \, dx = z \right\}.$$

We propose to extend this result to the heterogeneous case. We set, for all $\bar{F} \in \mathbb{M}^{3 \times 2}$, $z \in \mathbb{R}^3$ and *a.e.* $x_0 \in \omega$,

$$\begin{aligned} \mathcal{Q}^*W(x_0; \bar{F}|z) := \inf_{L > 0, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1,1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 : \right. \\ \left. \varphi \in W^{1,p}(Q' \times (-1,1); \mathbb{R}^3), \varphi(\cdot, x_3) \text{ } Q'\text{-periodic for a.e. } x_3 \in (-1,1), \right. \\ \left. \text{and } \frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi dx = z \right\} \quad (3.4) \end{aligned}$$

where, for *a.e.* $x \in \Omega$ and all $F \in \mathbb{M}^{3 \times 3}$, $\mathcal{Q}W(x; \cdot)$, the 3D-quasiconvexification of $W(x; \cdot)$ is defined as

$$\mathcal{Q}W(x; F) = \inf_{\phi \in W_0^{1,p}(Q; \mathbb{R}^3)} \int_Q W(x; F + D\phi(y)) dy$$

with $Q := (0,1)^3$. Since $\mathcal{Q}W(x; \cdot)$ is quasiconvex and satisfies a p -growth condition, for all $F_1, F_2 \in \mathbb{M}^{3 \times 3}$ and for *a.e.* $x \in \Omega$,

$$|\mathcal{Q}W(x; F_1) - \mathcal{Q}W(x; F_2)| \leq \beta (1 + |F_1|^{p-1} + |F_2|^{p-1}) |F_1 - F_2| \quad (3.5)$$

(see [7], Lem. 2.2). Elementary properties of \mathcal{Q}^*W are summarized in the following proposition:

Proposition 3.3.

i) For all $(\bar{F}, z) \in \mathbb{M}^{3 \times 2} \times \mathbb{R}^3$ and *a.e.* $x_0 \in \omega$,

$$0 \leq \mathcal{Q}^*W(x_0; \bar{F}|z) \leq \beta (|\bar{F}|^p + |z|^p + 1). \quad (3.6)$$

ii) \mathcal{Q}^*W is a Carathéodory function.

Proof.

i). We take $\varphi(x) := zx_3/L$ as test function in (3.4) and use the p -growth condition (1.1).

ii). It suffices to show that $\mathcal{Q}^*W(x_0; \cdot)$ is continuous for *a.e.* $x_0 \in \omega$. Let $\bar{F}_n \rightarrow \bar{F}$ and $z_n \rightarrow z$. We first prove that $\mathcal{Q}^*W(x_0; \cdot)$ is upper semicontinuous. For any $\delta > 0$, set $L > 0$ and $\varphi \in W^{1,p}(Q' \times (-1,1); \mathbb{R}^3)$ Q' -periodic satisfying $\frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi dx = z$ such that

$$\mathcal{Q}^*W(x_0; \bar{F}|z) \leq \frac{1}{2} \int_{Q' \times (-1,1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 \leq \mathcal{Q}^*W(x_0; \bar{F}|z) + \delta.$$

The sequence $\{\varphi_n(x) := \varphi(x) + x_3(z_n - z)/L\}$ is in $W^{1,p}(Q' \times (-1,1); \mathbb{R}^3)$ and it is Q' -periodic. Furthermore, $D_\alpha \varphi = D_\alpha \varphi_n$ and $\frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi_n dx = z_n$. Since

$$\|(\bar{F} + D_\alpha \varphi | LD_3 \varphi) - (\bar{F}_n + D_\alpha \varphi_n | LD_3 \varphi_n)\|_{L^\infty(Q' \times (-1,1); \mathbb{M}^{3 \times 3})} \leq |(\bar{F}|z) - (\bar{F}_n|z_n)| \rightarrow 0$$

while $\{D\varphi_n\}$ is bounded in $L^p(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})$, (3.5) together with Hölder's inequality, yields

$$\begin{aligned}
& \limsup_{n \rightarrow +\infty} \mathcal{Q}^*W(x_0; \bar{F}_n | z_n) - \mathcal{Q}^*W(x_0; \bar{F} | z) - \delta \\
& \leq \limsup_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1, 1)} |\mathcal{Q}W(x_0, x_3; \bar{F}_n + D_\alpha \varphi_n | LD_3 \varphi_n) - \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi)| \, dx \\
& \leq \limsup_{n \rightarrow +\infty} \frac{\beta}{2} \int_{Q' \times (-1, 1)} \left(1 + |(\bar{F}_n + D_\alpha \varphi_n | LD_3 \varphi_n)|^{p-1} + |(\bar{F} + D_\alpha \varphi | LD_3 \varphi)|^{p-1} \right) \\
& \quad \times |(\bar{F} + D_\alpha \varphi | LD_3 \varphi) - (\bar{F}_n + D_\alpha \varphi_n | LD_3 \varphi_n)| \, dx \\
& \leq \limsup_{n \rightarrow +\infty} C \left(1 + \|D\varphi\|_{L^p(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})}^{\frac{p-1}{p}} + \|D\varphi_n\|_{L^p(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})}^{\frac{p-1}{p}} \right) \\
& \quad \times \|(\bar{F} + D_\alpha \varphi | LD_3 \varphi) - (\bar{F}_n + D_\alpha \varphi_n | LD_3 \varphi_n)\|_{L^p(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})} = 0.
\end{aligned}$$

Passing to the limit when $\delta \searrow 0^+$ yields the desired upper semicontinuity. Let us prove now that $\mathcal{Q}^*W(x_0; \cdot)$ is lower semicontinuous. For every $n \geq 1$, choose $L_n > 0$ and $\varphi_n \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3)$ Q' -periodic satisfying $\frac{L_n}{2} \int_{Q' \times (-1, 1)} D_3 \varphi_n \, dx = z_n$ such that

$$\frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F}_n + D_\alpha \varphi_n | L_n D_3 \varphi_n) \, dx_\alpha dx_3 \leq \mathcal{Q}^*W(x_0; \bar{F}_n | z_n) + \frac{1}{n}.$$

Set $\tilde{\varphi}_n(x) := \varphi_n(x) + x_3(z - z_n)/L_n$, then $\tilde{\varphi}_n \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3)$ is Q' -periodic and satisfies $\frac{L_n}{2} \int_{Q' \times (-1, 1)} D_3 \tilde{\varphi}_n \, dx = z$. Since

$$\|(\bar{F} + D_\alpha \tilde{\varphi}_n | L_n D_3 \tilde{\varphi}_n) - (\bar{F}_n + D_\alpha \varphi_n | L_n D_3 \varphi_n)\|_{L^\infty(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})} \leq |(\bar{F} | z) - (\bar{F}_n | z_n)| \rightarrow 0$$

while, in view of the coercivity condition (1.1), the sequences $\{(D_\alpha \varphi_n | L_n D_3 \varphi_n)\}$ and $\{(D_\alpha \tilde{\varphi}_n | L_n D_3 \tilde{\varphi}_n)\}$ are bounded in $L^p(Q' \times (-1, 1); \mathbb{M}^{3 \times 3})$ uniformly in n , (3.5) implies that

$$\begin{aligned}
\mathcal{Q}^*W(x_0; \bar{F} | z) & \leq \liminf_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \tilde{\varphi}_n | L_n D_3 \tilde{\varphi}_n) \, dx_\alpha dx_3 \\
& \leq \liminf_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F}_n + D_\alpha \varphi_n | L_n D_3 \varphi_n) \, dx_\alpha dx_3 \\
& \leq \liminf_{n \rightarrow +\infty} \mathcal{Q}^*W(x_0; \bar{F}_n | z_n).
\end{aligned}$$

Thus $\mathcal{Q}^*W(x_0; \cdot)$ is lower semicontinuous and the continuity follows. \square

We propose to establish the following Γ -convergence result.

Theorem 3.4. *For all $A \in \mathcal{A}(\omega)$, $\mathcal{J}_\varepsilon(\cdot, \cdot; A)$ $\Gamma(L^p)$ -converges to $\mathcal{J}_{\{\varepsilon\}}(\cdot, \cdot; A)$. Further, for all $(u, \bar{b}) \in L^p(\Omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3)$,*

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = \begin{cases} 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) \, dx_\alpha & \text{if } u \in W^{1,p}(A; \mathbb{R}^3), \\ +\infty & \text{otherwise,} \end{cases}$$

where \mathcal{Q}^*W is given by (3.4).

We first note, as in [6] p. 1374, that, if $\mathcal{R}(\omega)$ is the countable family of all finite unions of open squares in ω with faces parallel to the axes, centered at rational points and with rational edge lengths, then there exists a subsequence $\{\varepsilon_n\} \subset \{\varepsilon\}$ such that $\mathcal{J}_{\{\varepsilon_n\}}(\cdot, \cdot; C)$ is, for all $C \in \mathcal{R}(\omega)$, the $\Gamma(L^p)$ -limit of $\mathcal{J}_{\varepsilon_n}(\cdot, \cdot; C)$.

Then, the analogue of step 2 in the proof of Theorem 2.5 of [6] holds, namely

Lemma 3.5. *For any $A \in \mathcal{A}(\omega)$ and $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(A; \mathbb{R}^3)$, there exists a sequence $\{u_n\} \subset W^{1,p}(A \times (-1, 1); \mathbb{R}^3)$ satisfying*

$$\begin{cases} u_n \rightarrow u \text{ in } L^p(A \times (-1, 1); \mathbb{R}^3), \\ \bar{b}_n := \frac{1}{2\varepsilon_n} \int_{-1}^1 D_3 u_n(\cdot; x_3) dx_3 \rightarrow \bar{b} \text{ in } L^p(A; \mathbb{R}^3), \\ \lim_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(u_n, \bar{b}_n; A) = \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A). \end{cases} \quad (3.7)$$

Proof. The coercivity condition (1.1) implies that whenever $u \in W^{1,p}(C; \mathbb{R}^3)$, we can choose the attainment sequence $\{u_n, \bar{b}_n\}$, so that (3.7) holds true. Now let us fix $\delta > 0$ and choose a subset C^δ of A in $\mathcal{R}(\omega)$ such that $\overline{C^\delta} \subset A$ and

$$\int_{A \setminus C^\delta} (1 + |D_\alpha u|^p) dx_\alpha < \frac{\delta}{2\beta}.$$

Consider a sequence $\{v_n^\delta\} \subset W^{1,p}(C^\delta \times (-1, 1); \mathbb{R}^3)$ satisfying

$$\begin{cases} v_n^\delta \xrightarrow{n \rightarrow +\infty} u \text{ in } L^p(C^\delta \times (-1, 1); \mathbb{R}^3), \\ \bar{b}_n^\delta := \frac{1}{2\varepsilon_n} \int_{-1}^1 D_3 v_n^\delta(\cdot, x_3) dx_3 \xrightarrow{n \rightarrow +\infty} \bar{b} \text{ in } L^p(C^\delta; \mathbb{R}^3), \\ \lim_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(v_n^\delta, \bar{b}_n^\delta; C^\delta) = \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; C^\delta). \end{cases}$$

In view of Lemma 2.2 in [3] (the proof in our context is identical to that of the homogeneous case), there exists a subsequence of $\{\varepsilon_n\}$ (not relabelled) and a sequence $\{\hat{v}_n^\delta\}$ in $W^{1,p}(C^\delta \times (-1, 1); \mathbb{R}^3)$ satisfying $\hat{v}_n^\delta = u$ on a neighborhood of $\partial C^\delta \times (-1, 1)$ such that

$$\begin{cases} \hat{v}_n^\delta \xrightarrow{n \rightarrow +\infty} u \text{ in } L^p(C^\delta \times (-1, 1); \mathbb{R}^3), \\ \bar{b}_n^\delta := \frac{1}{2\varepsilon_n} \int_{-1}^1 D_3 \hat{v}_n^\delta(\cdot; x_3) dx_3 \xrightarrow{n \rightarrow +\infty} \bar{b} \text{ in } L^p(C^\delta; \mathbb{R}^3), \\ \lim_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(\hat{v}_n^\delta, \bar{b}_n^\delta; C^\delta) = \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; C^\delta). \end{cases} \quad (3.8)$$

We extend \hat{v}_n^δ as u outside C^δ (and correspondingly extend \bar{b}_n^δ as 0). Since $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)$ is an increasing set function, we have $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; C^\delta) \leq \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A)$ and thus,

$$\begin{aligned} \limsup_{\delta \rightarrow 0^+} \limsup_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(\hat{v}_n^\delta, \bar{b}_n^\delta; A) &\leq \limsup_{\delta \rightarrow 0^+} \left\{ \lim_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(\hat{v}_n^\delta, \bar{b}_n^\delta; C^\delta) + 2\beta \int_{A \setminus C^\delta} (1 + |D_\alpha u|^p) dx_\alpha \right\} \\ &= \limsup_{\delta \rightarrow 0^+} \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; C^\delta) \\ &\leq \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \\ &\leq \liminf_{\delta \rightarrow 0^+} \liminf_{n \rightarrow +\infty} \mathcal{J}_{\varepsilon_n}(\hat{v}_n^\delta, \bar{b}_n^\delta; A). \end{aligned}$$

Remark that (3.8), together with coercivity, implies that

$$\|D_\alpha \hat{v}_n^\delta\|_{L^p(A \times (-1, 1); \mathbb{M}^{3 \times 2})} + \|\bar{b}_n^\delta\|_{L^p(A; \mathbb{R}^3)} \leq C,$$

independently of δ, n ; in particular, $\{\bar{b}_n^\delta\}$ lies in a subset of $L^p(A; \mathbb{R}^3)$, which is metrizable for the weak L^p -topology. A simple diagonalization lemma (Lem. 7.1 in [6]) permits to conclude the existence of a decreasing sequence $\{\delta(n)\} \searrow 0^+$ such that the sequence $\{u_n := \hat{v}_n^{\delta(n)}\}$ satisfies (3.7). \square

We now recall two results that will be of use in the proof of Lemma 3.10 below. Their proof can be found in [3] in the homogeneous case and the heterogeneity does not create any additional difficulty.

Proposition 3.6. *For any sequence $\{\varepsilon\} \searrow 0^+$, there exists a subsequence $\{\varepsilon_n\}$ such that, for any $(u, \bar{b}) \in W^{1,p}(\omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3)$, the set function $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)$ defined in (3.3) is the trace on $\mathcal{A}(\omega)$ of a Radon measure, which is absolutely continuous with respect to the 2-dimensional Lebesgue measure.*

By virtue of Lemma 3.5 and Proposition 3.6, we will assume henceforth that $\{\varepsilon_n\}$ denotes a subsequence of $\{\varepsilon\}$ such that the $\Gamma(L^p)$ -limit of $\mathcal{J}_{\varepsilon_n}$ exists, in which case it coincides with $\mathcal{J}_{\{\varepsilon_n\}}$, and such that, for every $(u, \bar{b}) \in W^{1,p}(\omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3)$, the set function $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)$ is the trace on $\mathcal{A}(\omega)$ of a Radon measure, which is absolutely continuous with respect to the 2-dimensional Lebesgue measure.

Proposition 3.7. *For all $(u, \bar{b}; A) \in W^{1,p}(\omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3) \times \mathcal{A}(\omega)$, the value of $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A)$ is unchanged if W is replaced by $\mathcal{Q}W$ in (3.2).*

Remark 3.8. If W does not depend on x_α , we can show as in [3] that for all $A \in \mathcal{A}(\omega)$, $\mathcal{J}_\varepsilon(\cdot, \cdot; A)$ $\Gamma(L^p)$ -converges to $\mathcal{J}_{\{\varepsilon\}}(\cdot, \cdot; A)$ and

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = 2 \int_A \mathcal{Q}^*W(D_\alpha u | \bar{b}) \, dx_\alpha,$$

for every $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3)$, where \mathcal{Q}^*W is given by (3.4).

Theorem 3.4 is a direct consequence of the following two lemmata.

Lemma 3.9. *For all $A \in \mathcal{A}(\omega)$ with A Lipschitz and for all $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(A; \mathbb{R}^3)$,*

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \geq 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) \, dx_\alpha.$$

Lemma 3.10. *For all $A \in \mathcal{A}(\omega)$ with A Lipschitz and for all $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(A; \mathbb{R}^3)$,*

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \leq 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) \, dx_\alpha.$$

Proof of Lemma 3.9. Let $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(A; \mathbb{R}^3)$. According to the equi-integrability theorem (Th. 1.1 in [2]) together with Lemma 3.5, there exists a subsequence of $\{\varepsilon_n\}$ (not relabelled) and a sequence $\{u_n\} \subset W^{1,p}(A \times (-1, 1); \mathbb{R}^3)$ such that

$$\left\{ \begin{array}{l} u_n \rightarrow u \text{ in } L^p(A \times (-1, 1); \mathbb{R}^3), \\ \frac{1}{2\varepsilon_n} \int_{-1}^1 D_3 u_n(\cdot, x_3) \, dx_3 \rightarrow \bar{b} \text{ in } L^p(A; \mathbb{R}^3), \\ \left\{ \left| \left(D_\alpha u_n \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n \right) \right|^p \right\} \text{ is equi-integrable,} \\ \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) = \lim_{n \rightarrow +\infty} \int_{A \times (-1, 1)} W \left(x_\alpha, x_3; D_\alpha u_n \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n \right) \, dx_\alpha \, dx_3. \end{array} \right.$$

We argue as in the proof of Lemma 2.4 with $F_n(x) := \left(D_\alpha u_n(x) \Big|_{\frac{1}{\varepsilon_n}} D_3 u_n(x) \right)$ and we obtain, since $W \geq \mathcal{Q}W$

$$\begin{aligned} \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) &\geq \limsup_{h \rightarrow +\infty} \sum_{i \in I_h} \frac{1}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h}} \left\{ \liminf_{n \rightarrow +\infty} \int_{A_{i,h} \times (-1,1)} W(y_\alpha, x_3; F_n(x_\alpha, x_3)) dx_\alpha dx_3 \right\} dy_\alpha \\ &\geq \limsup_{h \rightarrow +\infty} \sum_{i \in I_h} \frac{1}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h}} \left\{ \liminf_{n \rightarrow +\infty} \int_{A_{i,h} \times (-1,1)} \mathcal{Q}W(y_\alpha, x_3; F_n(x_\alpha, x_3)) dx_\alpha dx_3 \right\} dy_\alpha, \end{aligned}$$

where $\{A_{i,h}\}_{i \in I_h}$ denotes a finite family of disjoint open sets with diameter less than $1/h$, such that $\mathcal{L}^2(A \setminus \cup_{i \in I_h} A_{i,h}) = 0$ and $\text{Card} I_h = O(h^2)$. Applying, for *a.e.* $y_\alpha \in A_{i,h}$, Remark 3.8 to the Carathéodory function $(x_3; F) \mapsto W(y_\alpha, x_3; F)$, we get

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \geq \limsup_{h \rightarrow +\infty} \sum_{i \in I_h} \frac{1}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h}} \left\{ 2 \int_{A_{i,h}} \mathcal{Q}^*W(y_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) dx_\alpha \right\} dy_\alpha.$$

By Proposition 3.3 (ii), \mathcal{Q}^*W is a Carathéodory integrand, thus Scorza-Dragoni's Theorem implies the existence, for any $\eta > 0$, of a compact set $C_\eta \subset A$, such that

$$\mathcal{L}^2(A \setminus C_\eta) < \eta, \quad (3.9)$$

and \mathcal{Q}^*W is continuous on $C_\eta \times \mathbb{M}^{3 \times 2}$. Let $S_\lambda := \{x_\alpha \in A : |(D_\alpha u(x_\alpha) | \bar{b}(x_\alpha))| \leq \lambda\}$, thanks to Chebyshev's inequality

$$\mathcal{L}^2(A \setminus S_\lambda) < \frac{C}{\lambda^p}. \quad (3.10)$$

Consequently

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \sum_{i \in I_h} \frac{1}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h} \cap C_\eta} \left\{ 2 \int_{A_{i,h} \cap C_\eta \cap S_\lambda} \mathcal{Q}^*W(y_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) dx_\alpha \right\} dy_\alpha.$$

Since \mathcal{Q}^*W is continuous on $C_\eta \times \mathbb{M}^{3 \times 2}$, it is uniformly continuous on $C_\eta \times \bar{B}(0, \lambda)$ thus there exists a increasing and continuous function $\omega_{\eta, \lambda} : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ satisfying $\omega_{\eta, \lambda}(0) = 0$ and for every $y_\alpha \in A_{i,h} \cap C_\eta$ and every $x_\alpha \in A_{i,h} \cap C_\eta \cap S_\lambda$,

$$|\mathcal{Q}^*W(y_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) - \mathcal{Q}^*W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha))| \leq \omega_{\eta, \lambda}(|x_\alpha - y_\alpha|) \leq \omega_{\eta, \lambda}(1/h).$$

Using the fact that $\mathcal{L}^2(A_{i,h}) \text{Card} I_h \leq C$, we get

$$\begin{aligned} \sum_{i \in I_h} \frac{1}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h} \cap C_\eta} \int_{A_{i,h} \cap C_\eta \cap S_\lambda} |\mathcal{Q}^*W(y_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) - \mathcal{Q}^*W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha))| dx_\alpha dy_\alpha \\ \leq C \omega_{\eta, \lambda}(1/h) \xrightarrow{h \rightarrow +\infty} 0. \end{aligned}$$

Therefore,

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \geq \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \limsup_{h \rightarrow +\infty} \sum_{i \in I_h} \frac{2\mathcal{L}^2(A_{i,h} \cap C_\eta)}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h} \cap C_\eta \cap S_\lambda} \mathcal{Q}^*W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) dx_\alpha.$$

By virtue of the p -growth condition (3.6) together with (3.9), we get

$$\begin{aligned} \sum_{i \in I_h} \frac{\mathcal{L}^2(A_{i,h} \setminus C_\eta)}{\mathcal{L}^2(A_{i,h})} \int_{A_{i,h} \cap C_\eta \cap S_\lambda} \mathcal{Q}^* W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) \, dx_\alpha &\leq \beta(1 + \lambda^p) \sum_{i \in I_h} \mathcal{L}^2(A_{i,h} \setminus C_\eta) \\ &= \beta(1 + \lambda^p) \mathcal{L}^2(A \setminus C_\eta) \\ &< \beta(1 + \lambda^p) \eta \xrightarrow{\eta \rightarrow 0} 0. \end{aligned}$$

Thus, (3.9) and (3.10) yield

$$\begin{aligned} \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) &\geq 2 \limsup_{\lambda \rightarrow +\infty} \limsup_{\eta \rightarrow 0} \int_{A \cap C_\eta \cap S_\lambda} \mathcal{Q}^* W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) \, dx_\alpha \\ &= 2 \int_A \mathcal{Q}^* W(x_\alpha; D_\alpha u(x_\alpha) | \bar{b}(x_\alpha)) \, dx_\alpha. \end{aligned} \quad \square$$

Proof of Lemma 3.10. The proof is divided into three steps. First, we address the case where u is affine and \bar{b} is constant; then, that where u is piecewise affine and continuous, and \bar{b} piecewise constant. Finally, we address the general case.

Step 1. Let $A \in \mathcal{A}(\omega)$, we assume that

$$\begin{cases} u(x_\alpha) = \bar{F}x_\alpha + c, & (\bar{F}, c) \in \mathbb{M}^{3 \times 2} \times \mathbb{R}^3, \\ \bar{b}(x_\alpha) = z, & z \in \mathbb{R}^3. \end{cases}$$

Thanks to the density of $W^{1,\infty}(Q' \times (-1,1); \mathbb{R}^3)$ into $W^{1,p}(Q' \times (-1,1); \mathbb{R}^3)$ and to the p -growth hypothesis (1.1), for any $k \geq 1$, there exists $L_k > 0$ and $\varphi_k \in W^{1,\infty}(Q' \times (-1,1); \mathbb{R}^3)$ Q' -periodic satisfying $\frac{L_k}{2} \int_{Q' \times (-1,1)} D_3 \varphi_k \, dx = z$ and such that

$$Z_k(x_0; \bar{F} | z) := \frac{1}{2} \int_{Q' \times (-1,1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi_k | L_k D_3 \varphi_k) \, dx_\alpha \, dx_3 \leq \mathcal{Q}^* W(x_0; \bar{F} | z) + \frac{1}{k}.$$

We extend φ_k to $\mathbb{R}^2 \times (-1,1)$ by Q' -periodicity. Choose $r > 0$ small enough so that $Q'(x_0, r) \subset A$ where x_0 is a Lebesgue point of the Radon-Nikodym derivative of $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)$ with respect to the 2-dimensional Lebesgue measure and of $Z_k(\cdot; \bar{F} | z)$ for all $k \geq 1$. Fix k and set

$$u_n^k(x) := \bar{F}x_\alpha + c + L_k \varepsilon_n \varphi_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right).$$

Then,

$$u_n^k \xrightarrow{n \rightarrow +\infty} u \text{ in } L^p(Q'(x_0, r) \times (-1,1); \mathbb{R}^3),$$

and by virtue of Riemann-Lebesgue's lemma,

$$\frac{1}{2\varepsilon_n} \int_{-1}^1 D_3 u_n^k \, dx_3 = \frac{L_k}{2} \int_{-1}^1 D_3 \varphi_k \left(\frac{x_\alpha}{L_k \varepsilon_n}, x_3 \right) \, dx_3 \xrightarrow{n \rightarrow +\infty} \frac{L_k}{2} \int_{Q' \times (-1,1)} D_3 \varphi_k \, dx = \bar{b}.$$

So $\{u_n^k\}$ is admissible for $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; Q'(x_0, r))$ and, thanks to Proposition 3.7,

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; Q'(x_0, r)) \leq \liminf_{n \rightarrow +\infty} \int_{Q'(x_0, r) \times (-1,1)} \mathcal{Q}W \left(x_\alpha, x_3; D_\alpha u_n^k \left| \frac{1}{\varepsilon_n} D_3 u_n^k \right. \right) \, dx_\alpha \, dx_3.$$

Using an argument similar to that in the proof of Lemma 2.5, with $\mathcal{Q}W$ instead of W , we get

$$\frac{d\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)}{d\mathcal{L}^2}(x_0) \leq 2\mathcal{Q}^*W(x_0; \bar{F}|z).$$

Thus, integration over A yields

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \leq 2 \int_A \mathcal{Q}^*W(x_\alpha; \bar{F}|z) dx_\alpha.$$

Step 2. Assume that u is continuous and piecewise affine and \bar{b} is piecewise constant on A . There exists a partition A_1, \dots, A_N of A such that $u(x_\alpha) = \bar{F}_i x_\alpha + c_i$ and $\bar{b}(x_\alpha) = z_i$ on A_i , for all $i = 1, \dots, N$. Thanks to step 1, for all $i = 1, \dots, N$, we have

$$\mathcal{J}_{\{\varepsilon_n\}}(\bar{F}_i x_\alpha + c_i, z_i; A_i) \leq 2 \int_{A_i} \mathcal{Q}^*W(x_\alpha; \bar{F}_i|z_i) dx_\alpha.$$

In view of Proposition 3.6, $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; \cdot)$ is a measure and we thus get

$$\begin{aligned} \mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) &= \sum_{i=1}^N \mathcal{J}_{\{\varepsilon_n\}}(\bar{F}_i x_\alpha + c_i, z_i; A_i) \\ &\leq 2 \sum_{i=1}^N \int_{A_i} \mathcal{Q}^*W(x_\alpha; \bar{F}_i|z_i) dx_\alpha \\ &= 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) dx_\alpha. \end{aligned}$$

Step 3. Consider $A \in \mathcal{A}(\omega)$ with A Lipschitz and $u \in W^{1,p}(A; \mathbb{R}^3)$, $\bar{b} \in L^p(A; \mathbb{R}^3)$. There exists a sequence $\{u_n\}$ of continuous and piecewise affine functions in $W^{1,p}(A; \mathbb{R}^3)$ and a sequence $\{\bar{b}_n\}$ of piecewise constant functions in $L^p(A; \mathbb{R}^3)$ such that $u_n \rightarrow u$ in $W^{1,p}(A; \mathbb{R}^3)$ and $\bar{b}_n \rightarrow \bar{b}$ in $L^p(A; \mathbb{R}^3)$. Since $\mathcal{J}_{\{\varepsilon_n\}}(\cdot, \cdot; A)$ is lower semicontinuous, we get, in view of the previous step,

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \leq \liminf_{n \rightarrow +\infty} \mathcal{J}_{\{\varepsilon_n\}}(u_n, \bar{b}_n; A) \leq \liminf_{n \rightarrow +\infty} 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u_n|\bar{b}_n) dx_\alpha. \quad (3.11)$$

By Proposition 3.3 and Lebesgue's dominated convergence theorem,

$$\lim_{n \rightarrow +\infty} \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u_n|\bar{b}_n) dx_\alpha = \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) dx_\alpha. \quad (3.12)$$

Thus (3.11) and (3.12) yield

$$\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A) \leq 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) dx_\alpha. \quad \square$$

Proof of Theorem 3.4. The two previous lemmata demonstrate that, provided $A \in \mathcal{A}(\omega)$ is Lipschitz, then, for all $(u, \bar{b}) \in W^{1,p}(A; \mathbb{R}^3) \times L^p(A; \mathbb{R}^3)$, $\mathcal{J}_{\{\varepsilon_n\}}(u, \bar{b}; A)$ does not depend upon the choice of sequence $\{\varepsilon_n\}$. Thus, in light of Proposition 7.11 in [5], the whole sequence $\mathcal{J}_\varepsilon(u, \bar{b}; A) \Gamma(L^p)$ -converges to $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A)$ and

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) dx_\alpha.$$

Whenever $A \in \mathcal{A}(\omega)$ is an arbitrary open set, we define the nested sequence of Lipschitz open subsets $A_k := \{x_\alpha \in A : \text{dist}(x_\alpha, \partial A) > 1/k\}$ of A , so that $\overline{A_k} \subset A$ and $\cup_{k \geq 1} A_k = A$. Consider the sequence $\{u_k\}$ in $W^{1,p}(A; \mathbb{R}^3)$ such that $u_k = u$ on A_k . Since $\mathcal{J}_{\{\varepsilon\}}(\cdot, \cdot; A_k)$ is local and A_k is Lipschitz,

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A_k) = \mathcal{J}_{\{\varepsilon\}}(u_k, \bar{b}; A_k) = 2 \int_{A_k} \mathcal{Q}^*W(x_\alpha; D_\alpha u_k | \bar{b}) dx_\alpha = 2 \int_{A_k} \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) dx_\alpha.$$

But $\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; \cdot)$ is a measure, thus, letting $k \nearrow +\infty$,

$$\mathcal{J}_{\{\varepsilon\}}(u, \bar{b}; A) = 2 \int_A \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) dx_\alpha.$$

Then Remark 3.1 completes the proof of Theorem 3.4. \square

Remark 3.11. If W does not depend upon x , Proposition 1.1 (iii) of [3] states that

$$\begin{aligned} \mathcal{Q}^*W(\bar{F}|z) &= \inf_{L>0, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1,1)} W(\bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 : \varphi \in W^{1,p}(Q' \times (-1,1); \mathbb{R}^3), \right. \\ &\quad \left. \varphi(\cdot, x_3) \text{ } Q' \text{-periodic for a.e. } x_3 \in (-1,1), \frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi dx = z \right\} \\ &= \inf_{L>0, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1,1)} \mathcal{Q}W(\bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 : \varphi \in W^{1,p}(Q' \times (-1,1); \mathbb{R}^3), \right. \\ &\quad \left. \varphi(\cdot, x_3) \text{ } Q' \text{-periodic for a.e. } x_3 \in (-1,1), \frac{L}{2} \int_{Q' \times (-1,1)} D_3 \varphi dx = z \right\}. \end{aligned}$$

In other words, the result of [3] is recovered by Theorem 3.4.

Remark 3.12. Since \mathcal{Q}^*W is the integrand of the $\Gamma(L^p)$ -limit of \mathcal{J}_ε , which satisfies a p -coercivity condition (see (1.1)), for all $\bar{F} \in \mathbb{M}^{3 \times 2}$, for all $z \in \mathbb{R}^3$ and for a.e. $x_0 \in \omega$,

$$\beta'(|\bar{F}|^p + |z|^p) \leq \mathcal{Q}^*W(x_0; \bar{F}|z). \quad (3.13)$$

Remark 3.13. Theorem 3.4 implies that the functional

$$(u, \bar{b}) \mapsto \int_\omega \mathcal{Q}^*W(x_\alpha; D_\alpha u | \bar{b}) dx_\alpha$$

is sequentially weakly lower semicontinuous on $W^{1,p}(\omega; \mathbb{R}^3) \times L^p(\omega; \mathbb{R}^3)$. Therefore, $\mathcal{Q}^*W(x_0; \cdot | z)$ is quasiconvex and $\mathcal{Q}^*W(x_0; \bar{F} | \cdot)$ is convex. Thanks to the p -growth condition (3.6), $\mathcal{Q}^*W(x_0; \cdot | \cdot)$ is locally Lipschitz, because it is separately convex (see Th. 2.3 in [7]).

4. CLASSICAL MEMBRANE MODEL OBTAINED AS A ZERO BENDING MOMENT DENSITY

This section investigates the coherence of our results. In the absence of a bending moment density ($g_0 = 0$), we show below that Theorem 3.4 boils down to Theorem 2.3. We first give another form of the energy density \underline{W} similar to the definition of \mathcal{Q}^*W (see (3.4)). Specifically,

Proposition 4.1. For all $\bar{F} \in \mathbb{M}^{3 \times 2}$ and for a.e. $x_0 \in \omega$,

$$\underline{W}(x_0; \bar{F}) = \inf_{L, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 : \right. \\ \left. L > 0, \varphi \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3), \varphi(\cdot, x_3) \text{ } Q' \text{-periodic for a.e. } x_3 \in (-1, 1) \right\}.$$

Proof. Set

$$W^*(x_0; \bar{F}) := \inf_{L, \varphi} \left\{ \frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 : \right. \\ \left. L > 0, \varphi \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3), \varphi(\cdot, x_3) \text{ } Q' \text{-periodic for a.e. } x_3 \in (-1, 1) \right\}.$$

It is obvious that $\underline{W}(x_0; \bar{F}) \geq W^*(x_0; \bar{F})$. Conversely, for any $\delta > 0$, consider $L > 0$ and $\varphi \in W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3)$ Q' -periodic, such that

$$\frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi | LD_3 \varphi) dx_\alpha dx_3 \leq W^*(x_0; \bar{F}) + \delta.$$

We extend φ by Q' -periodicity and we set $\varphi_n(x_\alpha, x_3) := \frac{1}{n} \varphi(nx_\alpha, x_3)$. Then, Riemann-Lebesgue's Lemma applied to $\int_{-1}^1 \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi(\cdot, x_3) | LD_3 \varphi(\cdot, x_3)) dx_3$, implies that

$$\lim_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi_n | LnD_3 \varphi_n) dx_\alpha dx_3 \leq W^*(x_0; \bar{F}) + \delta. \quad (4.1)$$

For fixed n , the relaxation theorem of [1] (see Statement III.7 in [1])– applied to $Q' \times (-\frac{1}{Ln}, \frac{1}{Ln})$ and to $\psi_n(x_\alpha, x_3) := \varphi_n(x_\alpha, Ln x_3)$ – yields a sequence

$$\varphi_{n,k} \xrightarrow[k \rightarrow +\infty]{} \varphi_n \text{ in } W^{1,p}(Q' \times (-1, 1); \mathbb{R}^3)$$

such that,

$$\int_{Q' \times (-1, 1)} \mathcal{Q}W(x_0, x_3; \bar{F} + D_\alpha \varphi_n | LnD_3 \varphi_n) dx_\alpha dx_3 = \\ \lim_{k \rightarrow +\infty} \int_{Q' \times (-1, 1)} W(x_0, x_3; \bar{F} + D_\alpha \varphi_{n,k} | LnD_3 \varphi_{n,k}) dx_\alpha dx_3. \quad (4.2)$$

Thus (4.1) together with (4.2) give

$$\lim_{n \rightarrow +\infty} \lim_{k \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1, 1)} W(x_0, x_3; \bar{F} + D_\alpha \varphi_{n,k} | LnD_3 \varphi_{n,k}) dx_\alpha dx_3 \leq W^*(x_0; \bar{F}) + \delta.$$

Furthermore, we have,

$$\lim_{n \rightarrow +\infty} \lim_{k \rightarrow +\infty} \|\varphi_{n,k}\|_{L^p(Q' \times (-1, 1); \mathbb{R}^3)} = 0.$$

By a standard diagonalization process, we can find an increasing sequence $\{k(n)\}$, with $k(n) \xrightarrow{n} +\infty$ such that, if we set $\phi_n := \varphi_{n,k(n)}$,

$$\lim_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1,1)} W(x_0, x_3; \bar{F} + D_\alpha \phi_n | Ln D_3 \phi_n) dx_\alpha dx_3 \leq W^*(x_0; \bar{F}) + \delta, \quad (4.3)$$

and $\phi_n \rightarrow 0$ in $L^p(Q' \times (-1,1); \mathbb{R}^3)$. By virtue of the coercivity hypothesis (1.1),

$$\|(D_\alpha \phi_n | Ln D_3 \phi_n)\|_{L^p(Q' \times (-1,1); \mathbb{M}^{3 \times 3})} \leq C.$$

We define the following sequence of non negative bounded Radon measures

$$\lambda_n := (1 + |(D_\alpha \phi_n | Ln D_3 \phi_n)|^p) \chi_{Q' \times (-1,1)} \mathcal{L}^3.$$

The coercive character (1.1) of W permits to assert that, up to a subsequence (not relabelled), there exists a non negative bounded Radon measure λ such that

$$\lambda_n \xrightarrow{*} \lambda \text{ in } \mathcal{M}_b(\mathbb{R}^3).$$

We cut ϕ_n near the lateral boundary to obtain a sequence which vanishes on $\partial Q' \times (-1,1)$. Let $\theta_k \in \mathcal{C}_c^\infty(Q')$ a cut-off function defined by

$$\begin{cases} \theta_k(x_\alpha) := \begin{cases} 1 & \text{if } x_\alpha \in Q'(0, 1 - 1/k), \\ 0 & \text{if } x_\alpha \notin Q'(0, 1 - 1/(k+1)), \end{cases} \\ \|D_\alpha \theta_k\|_{L^\infty(Q')} \leq Ck^2. \end{cases} \quad (4.4)$$

We set $\phi_n^k := \theta_k \phi_n$, since $\phi_n^k = 0$ on $\partial Q' \times (-1,1)$, (4.3) together with (4.4) yields

$$\begin{aligned} W^*(x_0; \bar{F}) &\geq \liminf_{k \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \frac{1}{2} \int_{Q'(0, 1 - \frac{1}{k}) \times (-1,1)} W(x_0, x_3; \bar{F} + D_\alpha \phi_n^k | Ln D_3 \phi_n^k) dx_\alpha dx_3 - \delta \\ &\geq \liminf_{k \rightarrow +\infty} \liminf_{n \rightarrow +\infty} \frac{1}{2} \int_{Q' \times (-1,1)} W(x_0, x_3; \bar{F} + D_\alpha \phi_n^k | Ln D_3 \phi_n^k) dx_\alpha dx_3 \\ &\quad - \limsup_{k \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \frac{1}{2} \int_{(Q'(0, 1 - \frac{1}{k+1}) \setminus Q'(0, 1 - \frac{1}{k})) \times (-1,1)} W(x_0, x_3; \bar{F} + D_\alpha \phi_n^k | Ln D_3 \phi_n^k) dx_\alpha dx_3 \\ &\quad - \beta(1 + |\bar{F}|^p) \limsup_{k \rightarrow +\infty} \mathcal{L}^2 \left(Q' \setminus Q' \left(0, 1 - \frac{1}{k+1} \right) \right) - \delta \\ &\geq \underline{W}(x_0; \bar{F}) - \limsup_{k \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \left\{ C \lambda_n \left(\left(Q' \left(0, 1 - \frac{1}{k+1} \right) \setminus Q' \left(0, 1 - \frac{1}{k} \right) \right) \times (-1, 1) \right) \right. \\ &\quad \left. + C' k^{2p} \int_{Q' \times (-1,1)} |\phi_n|^p dx \right\} - \delta. \end{aligned} \quad (4.5)$$

Since $Q'(0, 1 - 1/k)$ is an increasing sequence of open sets, the union of which is Q' , we get

$$\begin{aligned} \limsup_{k \rightarrow +\infty} \limsup_{n \rightarrow +\infty} \lambda_n \left(\left(Q' \left(0, 1 - \frac{1}{k+1} \right) \setminus Q' \left(0, 1 - \frac{1}{k} \right) \right) \times (-1, 1) \right) \\ \leq \limsup_{k \rightarrow +\infty} \lambda \left(\overline{\left(Q' \left(0, 1 - \frac{1}{k+1} \right) \setminus Q' \left(0, 1 - \frac{1}{k} \right) \right) \times (-1, 1)} \right) \\ \leq \limsup_{k \rightarrow +\infty} \lambda \left(\left(Q' \setminus Q' \left(0, 1 - \frac{1}{k-1} \right) \right) \times [-1, 1] \right) = 0. \end{aligned}$$

Using the fact that $\phi_n \rightarrow 0$ in $L^p(Q' \times (-1, 1); \mathbb{R}^3)$ and letting δ tend to 0 in (4.5), we finally get

$$W^*(x_0; \bar{F}) \geq \underline{W}(x_0; \bar{F}). \quad \square$$

Now that \underline{W} and \mathcal{Q}^*W are expressed in near identical manner, Remarks 3.12 and 3.13 immediately imply that for all $\bar{F} \in \mathbb{M}^{3 \times 2}$ and for *a.e.* $x_0 \in \omega$, there exists $b_0 \in \mathbb{R}^3$ such that

$$\underline{W}(x_0; \bar{F}) = \min_{z \in \mathbb{R}^3} \mathcal{Q}^*W(x_0; \bar{F}|z) = \mathcal{Q}^*W(x_0; \bar{F}|b_0).$$

In the absence of bending moments, the linear form L given by (3.1) does not depend upon \bar{b} and we may perform explicitly the minimum in \bar{b} in the limit minimization problem. For $u \in W^{1,p}(\omega; \mathbb{R}^3)$, a classical measurability selection criterion (see [9], Chap. VIII, Th. 1.2), together with the coercivity condition (3.13), implies the existence of $\bar{b}_0 \in L^p(\omega; \mathbb{R}^3)$ such that for *a.e.* $x_0 \in \omega$,

$$\underline{W}(x_0; D_\alpha u(x_0)) = \min_{z \in \mathbb{R}^3} \mathcal{Q}^*W(x_0; D_\alpha u(x_0)|z) = \mathcal{Q}^*W(x_0; D_\alpha u(x_0)|\bar{b}_0(x_0)).$$

Thus,

$$\begin{aligned} \inf_{\bar{b} \in L^p(\omega; \mathbb{R}^3)} \int_\omega \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) \, dx_\alpha &\leq \int_\omega \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}_0) \, dx_\alpha \\ &= \int_\omega \underline{W}(x_\alpha; D_\alpha u) \, dx_\alpha \\ &= \int_\omega \min_{z \in \mathbb{R}^3} \mathcal{Q}^*W(x_\alpha; D_\alpha u|z) \, dx_\alpha \\ &\leq \int_\omega \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) \, dx_\alpha, \end{aligned} \quad (4.6)$$

where the last inequality holds for all $\bar{b} \in L^p(\omega; \mathbb{R}^3)$. Taking the infimum in \bar{b} in the last term of (4.6), the inequalities become equalities thus

$$\inf_{\bar{b} \in L^p(\omega; \mathbb{R}^3)} \int_\omega \mathcal{Q}^*W(x_\alpha; D_\alpha u|\bar{b}) \, dx_\alpha = \int_\omega \underline{W}(x_\alpha; D_\alpha u) \, dx_\alpha.$$

This shows that Theorem 2.3 is recovered from Theorem 3.4.

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