



Partial Differential Equations/Differential Geometry

A Note on the Bernstein property of a fourth order complex partial differential equations

Sur la propriété de Bernstein des équations différentielles partielles complexes d'ordre quatre

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ABSTRACT

For a smooth strictly plurisubharmonic function u on an open set $\Omega \subset \mathbb{C}^n$ and F a C^1 nondecreasing function on \mathbb{R}_+^* , we investigate the complex partial differential equations

$$\Delta_g \log \det(u_{i\bar{j}}) = F(\det(u_{i\bar{j}})) \|\nabla_g \log \det(u_{i\bar{j}})\|_g^2,$$

where Δ_g , $\|\cdot\|_g$ and ∇_g are the Laplacian, tensor norm and the Levi-Civita connexion, respectively, with respect to the Kähler metric $g = \partial\bar{\partial}u$. We show that the above PDE's has a Bernstein property, i.e. $\det(u_{i\bar{j}})$ is constant on Ω , provided that g is complete, the Ricci curvature of g is bounded below and F satisfies $\inf_{t \in \mathbb{R}^+} (2tF'(t) + \frac{F(t)^2}{n}) > \frac{1}{4}$ and $F(\max_{B(R)} \det u_{i\bar{j}}) = o(R)$.

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RÉSUMÉ

Pour une fonction u strictement plurisouharmonique de classe C^∞ sur un ouvert Ω de \mathbb{C}^n et F une fonction de classe C^1 croissante sur \mathbb{R}^+ , on considère l'équation aux dérivées partielles complexes

$$\Delta_g \log \det(u_{i\bar{j}}) = F(\det(u_{i\bar{j}})) \|\nabla_g \log \det(u_{i\bar{j}})\|_g^2,$$

où Δ_g , $\|\cdot\|_g$ et ∇_g sont respectivement le Laplacien, la norme et la connexion de Levi-Civita par rapport à la métrique Kählerienne $g = \partial\bar{\partial}u$. On montre que l'EDP précédente vérifie la propriété de Bernstein, i.e. $\det(u_{i\bar{j}})$ est constante sur Ω , pourvu que g soit complète, la courbure de Ricci de g soit minorée et F satisfasse $\inf_{t \in \mathbb{R}^+} (2tF'(t) + \frac{F(t)^2}{n}) > \frac{1}{4}$ et $F(\max_{B(R)} \det u_{i\bar{j}}) = o(R)$.

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1. Introduction

In this Note, we investigate a class of fourth order complex nonlinear partial differential equations for a strictly plurisubharmonic function u on an open set $\Omega \subset \mathbb{C}^n$:

$$\Delta_g \log \det(u_{i\bar{j}}) = F(\det(u_{i\bar{j}})) \|\nabla_g \log \det(u_{i\bar{j}})\|_g^2 \quad (1)$$

where Δ_g , $\|\cdot\|_g$ and ∇_g are the Laplacian, tensor norm and the Levi-Civita connexion, respectively, with respect to the Kähler metric $g = \partial\bar{\partial}u$. For $n \geq 2$, Eq. (1) is the Euler–Lagrange equation of the variational problem of the functional

$$\mathcal{V}_\Phi(u) = \int_{\Omega} \Phi(\det(u_{i\bar{j}})) dV_e \quad (2)$$

for some $\Phi \in C^4(\mathbf{R}_+^*)$. If $\Phi(t) = t^{1-\beta}$, $\beta \neq 1$, the Euler–Lagrange equation of \mathcal{V}_Φ is

$$\Delta_g \log \det(u_{i\bar{j}}) = \beta \|\nabla_g \log \det(u_{i\bar{j}})\|_g^2 \quad (3)$$

In [1], Chen and Li showed the following Bernstein property of the solution of (3).

Theorem 1.1. Let $\Omega \subset \mathbb{C}^n$ and $u \in C^2(\Omega)$ a strictly plurisubharmonic solution of (3). Assume that

- (i) (Ω, g) is complete,
- (ii) $Ric(g) \geq -C_1$ ($K \geq 0$),
- (iii) $\det(u_{i\bar{j}}) < C_2$

where $C_1, C_2 > 0$ are real constants. If $\beta > \frac{\sqrt{9n^2+4n+3n}}{4}$ then $\det(u_{i\bar{j}})$ is constant on Ω .

For more details about Bernstein property, one can see the paper [1] and references therein. In this Note, we will show the Bernstein property for the solutions of Theorem 1.1. Our main result is the following:

Theorem 1.2. Let $\Omega \subset \mathbb{C}^n$ be an open set and $u \in C^2(\Omega)$ a strictly plurisubharmonic solution of (1). Assume that

- (i) (Ω, g) is complete,
- (ii) $Ric(g) \geq -K$ ($K \geq 0$),
- (iii) $F \in C^1([0, +\infty[)$ nondecreasing satisfying

$$\inf_{t>0} \left(2tF'(t) + \frac{F(t)^2}{n} \right) > \frac{1}{4} \quad \text{and} \quad F\left(\max_{B(R)} \det(u_{i\bar{j}})\right) = o(R) \quad (4)$$

where $B(R) = \{z \in \Omega, d_g(z, z_0) < R\}$. Then $\det(u_{i\bar{j}})$ is constant on Ω .

If $F(t) = \beta$, the following corollary extends Theorem 1.1.

Corollary 1.3. Let $\Omega \subset \mathbb{C}^n$ be an open set and $u \in C^2(\Omega)$ a strictly plurisubharmonic solution of (3). Assume that (Ω, g) is complete and $Ric(g) \geq -K$ ($K \geq 0$). If $|\beta| > \frac{\sqrt{n}}{2}$ then $\det(u_{i\bar{j}})$ is constant on Ω .

Examples. Examples of function F which satisfies conditions of Theorem 1.2:

1. $F(t) = \log^{2\alpha} t$ where $\alpha \in]\alpha_0(n), \frac{1}{2}[$, $\alpha_0(n) > 0$.
2. $F(t) = (\log t) \log(\alpha + (\log t)^2)$ where $\alpha > e^{\frac{1}{4}}$.
3. $F(t) = f(\log t)$ where f is a solution of the Riccati equation $2y' + \frac{y^2}{n} = \alpha$, $\alpha > \frac{1}{4}$.

2. Proof of Theorem 1.2

Proof. We denote $\Psi = |\nabla U|^2$ where $U = \log \det(u_{i\bar{j}})$ and $f(t) = F(e^t)$. By Bochner formulae

$$\begin{aligned} \Delta_g \Psi &= 2f'(U)\Psi^2 + f(U)(\langle \nabla_g \Psi, \nabla_g U \rangle + \langle \nabla_g U, \nabla_g \Psi \rangle) \\ &\quad + \sum u^{i\bar{j}} u^{k\bar{l}} U_{i\bar{l}} U_{j\bar{k}} + \sum u^{i\bar{j}} u^{k\bar{l}} U_{i\bar{l}} U_{j\bar{k}} + \sum u^{i\bar{j}} u^{m\bar{s}} U_{i\bar{l}} U_m U_{j\bar{s}} R_{i\bar{s}} \end{aligned} \quad (5)$$

where (u^{ij}) is the inverse matrix of (u_{ij}) and $R_{i\bar{s}}$ is the Ricci curvature of g

$$R_{i\bar{s}} = -\frac{\partial^2 \log \det(u_{k\bar{l}})}{\partial z_i \partial \bar{z}_s} = \frac{\partial^2 \log U}{\partial z_i \partial \bar{z}_s}$$

Let $z \in \Omega$ fixed, we can choose normal coordinates (z_1, \dots, z_n) at z such that

$$u_{i\bar{j}} = \delta_{ij}, \quad U_1 = U_{\bar{1}}, \quad U_i = U_{\bar{i}} = 0, \quad \forall i > 1$$

At z the formula (5) becomes

$$\begin{aligned} \frac{\Delta_g \Psi}{\Psi} &= 2f'(U)\Psi + \left(f(U) - \frac{1}{2}\right) \left(\frac{\Psi_1}{\Psi} U_{\bar{1}} + U_1 \frac{\Psi_{\bar{1}}}{\Psi}\right) + \frac{1}{2} \left(\frac{\Psi_1}{\Psi} U_{\bar{1}} + U_1 \frac{\Psi_{\bar{1}}}{\Psi}\right) \\ &\quad + \frac{\sum U_{ik} U_{j\bar{k}}}{\Psi} + \frac{\sum U_{i\bar{k}} U_{\bar{i}k}}{\Psi} - \frac{U_{1\bar{1}} U_1 U_{\bar{1}}}{\Psi} \end{aligned} \quad (6)$$

Since $\sum U_{ik} U_{j\bar{k}} \geq U_{11} U_{\bar{1}\bar{1}}$ and

$$\begin{aligned} \frac{1}{2} \left(\frac{\Psi_1}{\Psi} U_{\bar{1}} + U_1 \frac{\Psi_{\bar{1}}}{\Psi}\right) - \frac{U_{1\bar{1}} U_1 U_{\bar{1}}}{\Psi} &= \frac{1}{2\Psi} 2\Re(U_{11} U_{\bar{1}} U_{\bar{1}}) \\ &\geq \frac{U_{11} U_{\bar{1}\bar{1}}}{\Psi} - \frac{1}{4}\Psi \end{aligned}$$

we obtain

$$\frac{1}{2} \left(\frac{\Psi_1}{\Psi} U_{\bar{1}} + U_1 \frac{\Psi_{\bar{1}}}{\Psi}\right) - \frac{U_{1\bar{1}} U_1 U_{\bar{1}}}{\Psi} + \frac{\sum U_{ik} U_{j\bar{k}}}{\Psi} \geq -\frac{1}{4}\Psi \quad (7)$$

Let G be the function defined on $B(R) = B_g(z_0, R)$ by

$$G(z) = (R^2 - r(z)^2)^2 \Psi$$

where $r(z) = d_g(z, z_0)$. Since G vanishes at the boundary $\partial B(R)$, the maximum of G is attained on $w \in B(R)$. By the maximum principle

$$\nabla_g G(w) = 0$$

$$\Delta_g G(w) \leq 0$$

A computation at w gives

$$\frac{\Psi_i}{\Psi} = \frac{4rr_i}{R^2 - r^2} \quad (8)$$

$$\begin{aligned} \Delta_g G &= 2\Psi \|\nabla_g(R^2 - r^2)\|^2 + 2(R^2 - r^2)\Psi \Delta_g(R^2 - r^2) \\ &\quad + 2(R^2 - r^2) \sum (R^2 - r^2)_i \Psi_i + (R^2 - r^2)^2 \Delta_g \Psi \\ &= 2\Psi r^2 \|\nabla_g r\|^2 - 2\Psi(R^2 - r^2) \Delta_g r^2 + (R^2 - r^2)^2 \Delta_g \Psi \\ &\quad - 4r(R^2 - r^2) \sum_i r_i \Psi_i \leq 0 \end{aligned} \quad (9)$$

Since $\|\nabla_g r\|_g = 1$, (8) and (9) imply

$$\frac{\Delta_g \Psi}{\Psi} \leq \frac{2\Delta_g r^2}{R^2 - r^2} + \frac{14r^2}{(R^2 - r^2)^2} \quad (10)$$

Let $\eta >$ fixed such that $\inf_{t \in \mathbb{R}} (2f'(t) + \frac{f(t)^2}{n}) > 2\eta + \frac{1}{4}$. By Schwarz inequality

$$\begin{aligned} \left(f(U) - \frac{1}{2}\right) \left(\frac{\Psi_1}{\Psi} U_{\bar{1}} + U_1 \frac{\Psi_{\bar{1}}}{\Psi}\right) &= 2\Re \left(U_1 \frac{4rr_1(f(U) - \frac{1}{2})}{R^2 - r^2}\right) \\ &\geq -\eta \Psi - \frac{16(f(U) - \frac{1}{2})^2 r^2}{\eta(R^2 - r^2)^2} \end{aligned} \quad (11)$$

Since $\Delta U = f(U)$ we have

$$\frac{\sum U_{\bar{i}\bar{k}}U_{\bar{i}\bar{k}}}{\Psi} \geq \frac{\sum (U_{\bar{i}\bar{i}})^2}{\Psi} \geq \frac{(\sum U_{\bar{i}\bar{i}})^2}{n\Psi} = \frac{(\Delta U)^2}{n\Psi} = \frac{(f(U))^2}{n\Psi} \quad (12)$$

The inequalities (6), (11) and (12) give

$$\frac{\Delta_g \Psi}{\Psi} \geq \eta\Psi - \frac{16(f(U) - \frac{1}{2})^2 r^2}{\eta(R^2 - r^2)^2} \quad (13)$$

By Laplacian comparison theorem [2]: $\Delta_g r^2 \leq 2n + 2(n-1)\sqrt{K}r$. Since $|f(U) - \frac{1}{2}|^2 \leq 2f(U)^2 + \frac{1}{2}$ and f is nondecreasing, the inequalities (10) and (11) give on $B(R)$

$$(R^2 - r^2)^2 \Psi \leq G(w) \leq C_1(1+R)R^2 + C_2 \left(f \left(\max_{B(R)} U \right) \right)^2 R^2$$

Hence if $z \in B(R/2)$, we have

$$\Psi(z) \leq C_1 \frac{1+R}{R^2} + C_2 \left(\frac{f(\max_{B(R)} U)}{R} \right)^2$$

Since (Ω, g) is complete, letting R to infinity, we obtain $\Psi(z) = \|\nabla_g U(z)\|^2 = 0$, i.e. $\det(u_{i\bar{j}}) = C$ on Ω . \square

References

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- [2] R.E. Greene, H. Wu, Function Theory on Manifolds Which Possess a Pole, Lecture Notes in Math., vol. 699, 1979.