

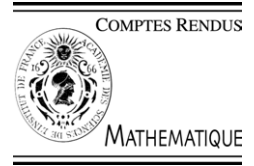


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## Mathematical Analysis/Harmonic Analysis

# Pointwise regularity criteria

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### Abstract

A wavelet characterization of the pointwise regularity condition  $T_u^p(x_0)$  of Calderón and Zygmund is obtained. The extremal case (a pointwise BMO condition) yields the sharpest wavelet condition which is implied by pointwise Hölder regularity; in particular, this criterium is sharper than the usual two-microlocal condition. *To cite this article: S. Jaffard, C. R. Acad. Sci. Paris, Ser. I 339 (2004).*

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### Résumé

**Critères de régularité ponctuelle.** On obtient une caractérisation par ondelettes de la condition de régularité ponctuelle  $T_u^p(x_0)$  de Calderón et Zygmund. Le cas extrême (une condition de type BMO local) fournit la condition la plus précise sur les modules des coefficients d'ondelette impliquée par la régularité Hölderienne ponctuelle ; en particulier elle est plus fine que le critère deux-microlocal usuel. *Pour citer cet article : S. Jaffard, C. R. Acad. Sci. Paris, Ser. I 339 (2004).*

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### Version française abrégée

Il existe plusieurs définitions possibles de la régularité ponctuelle d'une fonction  $f : \mathbb{R}^d \rightarrow \mathbb{R}$  ; la plus couramment utilisée est la *régularité Hölderienne*.

**Définition 0.1.** Soient  $f \in L_{\text{loc}}^\infty$ ,  $x_0 \in \mathbb{R}^d$  et  $\alpha \geq 0$  ; alors  $f \in C^\alpha(x_0)$  s'il existe  $R > 0$ ,  $C > 0$ , et un polynôme  $P$  de degré inférieur à  $\alpha$  tels que

$$\text{si } |x - x_0| \leq R \text{ alors } |f(x) - P(x - x_0)| \leq C|x - x_0|^\alpha. \quad (1)$$

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La condition  $L_{\text{loc}}^\infty$  est nécessaire ; en effet (1) implique que  $f$  est bornée au voisinage de  $x_0$ . Donc cette définition n'est pas utilisable si le cadre naturel pour  $f$  est  $L_{\text{loc}}^p$  pour un  $p < \infty$ . Un autre inconvénient a été découvert par Calderón et Zygmund en 1961, cf. [2] : la condition (1) n'est pas conservée sous l'action des opérateurs pseudodifférentiels classiques d'ordre 0. De plus, les ondelettes ne sont pas des bases inconditionnelles de l'espace  $C^\alpha(x_0)$ , cf. [3]. Aussi Calderón et Zygmund ont introduit les conditions de régularité ponctuelle  $T_u^p(x_0)$  qui ne souffrent pas de ces inconvénients. Soit  $p \in (1, \infty)$ ,  $f \in L_{\text{loc}}^p$  et  $u \geq -d/p$  ; on dit que  $f \in T_u^p(x_0)$  s'il existe  $R, C > 0$  et un polynôme  $P$  de degré inférieur à  $u$  tels que

$$\forall r \leq R, \quad \left( \frac{1}{r^d} \int_{B(x_0, r)} |f(x) - P(x - x_0)|^p dx \right)^{1/p} \leq Cr^u.$$

En tenant compte du fait que les extensions « naturelles » (du point de vue de l'analyse harmonique) des espaces  $L^p$  pour  $p = \infty$  et  $p \leq 1$  sont, respectivement, l'espace BMO et les espaces de Hardy réels  $H^p$ , on peut étendre les conditions  $T_u^p(x_0)$  à ces valeurs de  $p$ .

**Définition 0.2.** Soit  $p \in (0, 1]$ ,  $f \in H_{\text{loc}}^p$  et  $u \geq -d/p$  ;  $f \in T_u^p(x_0)$  s'il existe  $R, C > 0$  et un polynôme  $P$  de degré inférieur à  $u$  tels que  $\|(f - P)1_{B(x_0, r)}\|_p \leq Cr^{u+d/p}$ .

Soit  $f \in BMO_{\text{loc}}$  ;  $f \in T_u^\infty(x_0)$  s'il existe  $R, C > 0$  et un polynôme  $P$  de degré inférieur à  $u$  tels que  $\|(f - P)1_{B(x_0, r)}\|_{BMO} \leq Cr^u$ .

Notre but est d'obtenir une caractérisation de ces conditions pour tout  $p \in (0, +\infty]$  par un critère portant sur les modules des coefficients d'ondelette de  $f$ .

Soient  $\psi^{(i)}$ ,  $i = 1, \dots, 2^d - 1$ , des fonctions  $C^A$  à support compact (où  $A$  est choisi suffisamment grand) et engendrant une base d'ondelettes, c'est-à-dire que les  $2^{dj/2}\psi^{(i)}(2^j x - k)$  ( $i = 1, \dots, 2^d - 1$ ,  $j \in \mathbb{Z}$ ,  $k \in \mathbb{Z}^d$ ) forment une base orthonormée de  $L^2(\mathbb{R}^d)$ . On utilise l'indexation des ondelettes par les cubes dyadiques (rappelée dans le texte anglais). On notera  $c_\lambda = 2^{dj} \int \psi^{(i)}(2^j x - k) f(x) dx$ . Si  $x_0 \in \mathbb{R}^d$ , on notera  $\lambda_j(x_0)$  le cube dyadique de largeur  $2^{-j}$  contenant  $x_0$ , et  $S_f(j, x_0)(x) = (\sum_{\lambda \subset 3\lambda_j(x_0)} |c_\lambda|^2 1_\lambda(x))^{1/2}$ .

**Théorème 0.3.** Soit  $p \in (0, \infty)$  and  $u > -d/p$  ; si  $f \in T_u^p(x_0)$ , alors  $\exists C \geq 0$  tel que  $\forall j \geq 0$ ,

$$\|S_f(j, x_0)\|_p \leq C2^{-j(u+d/p)}. \quad (2)$$

Si  $p = +\infty$ , cette condition devient

$$\forall \lambda \subset 3\lambda_j(x_0), \quad \left( \sum_{\lambda' \subset \lambda} 2^{-dj'} |c_{\lambda'}|^2 \right)^{1/2} \leq C2^{-dl/2} 2^{-uj}, \quad (3)$$

où la largeur de  $\lambda$  est notée  $2^{-l}$  et la largeur de  $\lambda'$  est notée  $2^{-j'}$ .

Réciproquement, si (2) est vérifiée (ou si (3) est vérifiée dans le cas  $p = +\infty$ ), et si  $u \notin \mathbb{N}$ , alors  $f \in T_u^p(x_0)$ .

On remarquera que, si  $p = 2$ , cette caractérisation se simplifie en

$$\sum_{\lambda' \subset 3\lambda_j(x_0)} 2^{-d(j'-j)} |c_{\lambda'}|^2 \leq C2^{-2uj},$$

qui avait été obtenue antérieurement par Yves Meyer (communication personnelle).

## 1. Introduction

Several definitions for the pointwise regularity of a function  $f: \mathbb{R}^d \rightarrow \mathbb{R}$  can be introduced depending on the global assumptions that are made on  $f$ . The most widely used is the *Hölder criterium*.

**Definition 1.1.** Let  $f \in L_{\text{loc}}^\infty$ ,  $x_0 \in \mathbb{R}^d$  and  $\alpha \geq 0$ ; then  $f \in C^\alpha(x_0)$  if  $\exists R > 0$ ,  $C > 0$ , and a polynomial  $P$  of degree less than  $\alpha$  such that

$$\text{if } |x - x_0| \leq R \text{ then } |f(x) - P(x - x_0)| \leq C|x - x_0|^\alpha. \quad (4)$$

The global setting supplied by  $L_{\text{loc}}^\infty$  is implicitly required by (4); indeed, (4) implies that  $f$  is bounded in a neighbourhood of  $x_0$ . Thus Definition 1.1 cannot supply a sensible notion of pointwise regularity if the natural setting for  $f$  is  $L_{\text{loc}}^p$  for  $p < \infty$ . Another drawback is less obvious and was already pointed out by Calderón and Zygmund in 1961, see [2]: The pointwise Hölder condition is not preserved under classical pseudodifferential operators of order 0. This instability property has a counterpart in wavelet analysis: Wavelet bases are not unconditional bases of the space  $C^\alpha(x_0)$ ; even more is true: There exist two functions  $f$  and  $g$  which share the same moduli of wavelet coefficients, and nonetheless satisfy  $f \in C^\alpha(x_0)$  whereas  $\forall \beta > 0$ ,  $g \notin C^\beta(x_0)$ . Thus Definition 1.1 is unsuitable in several settings; Calderón and Zygmund introduced the following extension which makes sense in the  $L^p$  setting and is preserved under singular integral operators.

**Definition 1.2.** Let  $p \in (1, \infty)$ ,  $f \in L_{\text{loc}}^p$  and  $u \geq -d/p$ ; then  $f \in T_u^p(x_0)$  if  $\exists R, C > 0$  and a polynomial  $P$  of degree less than  $u$  such that

$$\forall r \leq R, \quad \left( \frac{1}{r^d} \int_{B(x_0, r)} |f(x) - P(x - x_0)|^p dx \right)^{1/p} \leq Cr^u. \quad (5)$$

Note that this condition can be rewritten  $\|(f - P)1_{B(x_0, r)}\|_p \leq Cr^{u+d/p}$ . If one keeps in mind the requirement of using a criterium which is invariant under pseudodifferential operators of order 0, the following definition is the natural extension of  $T_u^p(x_0)$  outside the range  $p \in (1, \infty)$ . (Recall that, if  $p \leq 1$ , then  $H^p$  denotes the real Hardy space, see [6].)

**Definition 1.3.** Let  $p \in (0, 1]$ ,  $f \in H_{\text{loc}}^p$  and  $u \geq -d/p$ ; then  $f \in T_u^p(x_0)$  if  $\exists R, C > 0$  and a polynomial  $P$  of degree less than  $u$  such that  $\|(f - P)1_{B(x_0, r)}\|_p \leq Cr^{u+d/p}$ .

Let  $f \in BMO_{\text{loc}}$ ; then  $f \in T_u^\infty(x_0)$  if  $\exists R, C > 0$  and a polynomial  $P$  of degree less than  $u$  such that  $\|(f - P)1_{B(x_0, r)}\|_{BMO} \leq Cr^u$ .

Let  $p \in (0, +\infty]$ ; then the  $p$ -exponent of  $f$  at  $x_0$  is  $h_f^p(x_0) = \sup\{u: f \in T_u^p(x_0)\}$ .

The motivations for considering this new types of pointwise conditions are of a different nature for  $p = \infty$  and for  $p \leq 1$ . If  $p = +\infty$ , then the  $T_u^\infty(x_0)$  condition is the sharpest condition which is implied by  $C^u(x_0)$  and can be characterized by a condition bearing on the moduli of the wavelet coefficients of  $f$ ; it is therefore stronger than the two-microlocal conditions  $f \in C^{u, -u}(x_0)$  of [3]. In particular, while the two-microlocal condition can be satisfied by distributions which do not coincide with a function in a neighbourhood of  $x_0$ , the  $T_u^\infty(x_0)$  wavelet characterization implies that (5) holds for any  $p < \infty$ . Another motivation is supplied by the analysis of domains with fractal boundaries; one way to understand the geometry of a domain is to use analytic tools on its characteristic function and, in particular, perform its multifractal analysis. This cannot be done using the Hölder exponent as a measure for pointwise regularity since the Hölder exponent of a characteristic function only takes the two values 0 and  $+\infty$ ; thus no characteristic function is multifractal in this sense. By contrast, the  $p$ -exponent can take any non-negative value, thus opening the way to a multifractal analysis of domains, see [5].

If  $p < 1$ , the condition  $f \in H_{loc}^p$ , allows one to deal with singularities such as  $|x - x_0|^{-a}$  near  $x_0$  for  $a < d/p$ ; therefore using arbitrarily small values of  $p$  allows one to deal with singularities of arbitrarily large exponent  $a$ , which is needed in some applications, see [1].

Clearly,  $C^u(x_0) \leftrightarrow T_u^\infty(x_0)$  and, if  $+\infty \geq p \geq q > 0$ , then  $T_u^p(x_0) \leftrightarrow T_u^q(x_0)$ . Using the classical interpolation results between  $L^p$  and/or  $H^p$  spaces, it follows that, if  $f \in T_u^p(x_0) \cap T_v^q(x_0)$ , and if  $r$  is such that  $\frac{1}{r} = \frac{\alpha}{p} + \frac{1-\alpha}{q}$  with  $0 < \alpha < 1$ , then  $f \in T_r^w(x_0)$  with  $w = \alpha u + (1-\alpha)v$ . Thus, for  $x_0$  given, the function  $q \rightarrow h_f^{1/q}(x_0)$  is defined on an interval of the form  $[q_0, +\infty)$  or  $(q_0, +\infty)$ , where it is concave and increasing.

**2. Wavelet characterization**

Let  $\psi^{(i)}$ ,  $i = 1, \dots, 2^d - 1$ , be compactly supported  $C^A$  functions (where  $A$  is large enough) generating a wavelet basis, i.e. the  $2^{dj/2}\psi^{(i)}(2^j x - k)$  ( $i = 1, \dots, 2^d - 1, j \in \mathbb{Z}, k \in \mathbb{Z}^d$ ) form an orthonormal basis of  $L^2(\mathbb{R}^d)$ . Wavelets will be indexed by dyadic cubes as follows: We can consider that  $i$  takes values among all dyadic subcubes  $\lambda_i$  of  $[0, 1)^d$  of width  $1/2$  except for  $[0, 1/2)^d$ ; thus, the set of indices  $(i, j, k)$  can be relabelled using dyadic cubes:  $\lambda$  denotes the cube  $\{x: 2^j x - k \in \lambda_i\}$ ; we note  $\psi_\lambda(x) = \psi^{(i)}(2^j x - k)$  (an  $L^\infty$  normalization is used), and  $c_\lambda = 2^{dj} \int \psi_\lambda(x) f(x) dx$ ;  $\exists C: \text{supp}(\psi_\lambda) \subset C\lambda$  where  $C\lambda$  denotes the cube of same center as  $\lambda$  and  $C$  times wider. If  $x_0 \in \mathbb{R}^d$ , then  $\lambda_j(x_0)$  denotes the unique dyadic cube of width  $2^{-j}$  which contains  $x_0$ , and the local square function is  $S_f(j, x_0)(x) = (\sum_{\lambda \subset 3\lambda_j(x_0)} |c_\lambda|^2 1_\lambda(x))^{1/2}$ .

**Theorem 2.1.** Let  $p \in (0, \infty)$  and  $u > -d/p$ ; if  $f \in T_u^p(x_0)$ , then  $\exists C \geq 0$  such that  $\forall j \geq 0$ ,

$$\|S_f(j, x_0)\|_p \leq C 2^{-j(u+d/p)}. \tag{6}$$

If  $p = +\infty$ , this condition becomes

$$\forall \lambda \subset 3\lambda_j(x_0), \left( \sum_{\lambda' \subset \lambda} 2^{-dj'} |c_{\lambda'}|^2 \right)^{1/2} \leq C 2^{-dj/2} 2^{-uj}, \tag{7}$$

where  $2^{-l}$  is the width of  $\lambda$  and  $2^{-j'}$  is the width of  $\lambda'$ .

Conversely, if (6) holds (or if (7) holds in the case  $p = +\infty$ ) and if  $u \notin \mathbb{N}$ , then  $f \in T_u^p(x_0)$ .

**Proof of Theorem 2.1.** Assume first that  $p < \infty$ ; then (see [6])  $f \in L^p(\mathbb{R}^d)$  if  $p > 1$ , or  $f \in H^p(\mathbb{R}^d)$  if  $p \leq 1$  if and only if  $(\sum_\lambda |c_\lambda|^2 1_\lambda(x))^{1/2} \in L^p$ . The direct part of the theorem follows by applying this characterization to  $g(x) = (f(x) - P(x - x_0)) 1_{B(x_0, D2^{-j})}(x)$  and noticing that, if  $D$  is large enough and  $\lambda \subset 3\lambda_j(x_0)$ , then the corresponding wavelet coefficients of  $f$  and  $g$  coincide. If  $p = +\infty$ , the argument is the same using the characterization of BMO, see [6]:  $\exists C, \forall \lambda, \sum_{\lambda' \subset \lambda} 2^{-dj'} |c_{\lambda'}|^2 \leq C \text{Meas}(\lambda)$ .

Let us now prove the converse part. We can forget the ‘low frequency component’ of  $f$  corresponding to  $j < 0$  in its wavelet decomposition, since its contribution belongs locally to  $C^A(\mathbb{R}^d)$ . Let  $\Lambda_j$  denote the set of dyadic cubes of width  $2^{-j}$ ,  $\Delta_j f = \sum_{\lambda \in \Lambda_j} c_\lambda \psi_\lambda$ , and let  $P_j(x - x_0)$  denote the Taylor polynomial of  $\Delta_j f$  of degree  $[u]$  at  $x_0$ ; (6) or (7) imply that,

$$\text{if } \text{dist}(\lambda, x_0) \leq D 2^{-j}, \text{ then } |c_\lambda| \leq C 2^{-uj}. \tag{8}$$

Let  $\rho > 0$  be fixed and let  $J$  be defined by  $2^{-J} \leq \rho < 2 \cdot 2^{-J}$  and  $L$  be a constant which will be fixed later, but depends only on the size of the support of the wavelets. If  $j \leq J + L$ , then at most  $C$  of the  $\psi_\lambda$  have a support intersecting  $B = B(x_0, \rho)$  and each of them satisfies (8). It follows from Taylor’s formula that, if  $x \in B$  and  $j \leq J + L$ , then  $|\Delta_j f(x) - P_j(x - x_0)| \leq C \rho^{|u|+1} 2^{j(|u|+1-u)}$ , and therefore

$$\sum_{j=0}^{J+L} |\Delta_j f(x) - P_j(x - x_0)| \leq C \rho^u. \tag{9}$$

It follows also from (8) that, if  $|k| \leq [u] + 1$ , then,  $\forall j \geq 0$ ,  $|\Delta_j^{(k)} f(x_0)| \leq 2^{(kl-u)j}$ ; therefore the series  $P(x - x_0) = \sum_{j=0}^{\infty} P_j(x - x_0) = \sum_{j=0}^{\infty} \sum_{|k| < u} \frac{\Delta_j f^{(k)}(x_0)}{k!} (x - x_0)^k$  converges and, if  $|x - x_0| \leq \rho$ , then

$$\sum_{j=J+L}^{\infty} |P_j(x - x_0)| \leq C \sum_{j=J+L}^{\infty} \sum_{|k| < u} 2^{(kl-u)j} \rho^k \leq C \rho^u. \tag{10}$$

Let now  $g_J(x) = \sum_{j=J+L}^{\infty} \Delta_j f(x)$ ; then  $\|g_J 1_B\|_p \leq \|\sum_{j=J+L}^{\infty} \sum_{\lambda \subset B(x_0, 2\rho)} c_{\lambda} \psi_{\lambda}\|_p$  where  $L$  has been picked large enough so that both functions coincide on  $B$ . Using the wavelet characterization of  $L^p$ , the right hand side is bounded by

$$C \left\| \left( \sum_{j=J+L}^{\infty} \sum_{\lambda \subset B(x_0, 2\rho)} |c_{\lambda}|^2 1_{\lambda} \right)^{1/2} \right\|_p \leq S_f(j - L, x_0) \leq C 2^{-j(u+d/p)}. \tag{11}$$

The required estimate for  $\|(f - P(x - x_0)) 1_{B(x_0, \rho)}\|_p$  follows immediately from (9), (10) and (11).

The case  $p = \infty$  is completely similar.

### 3. Remarks and implications in multifractal analysis

If  $p = 2$ , this characterization boils down to a local  $l^2$  condition on the wavelet coefficients

$$\sum_{\lambda' \subset 3\lambda_j(x_0)} 2^{-d(j'-j)} |c_{\lambda'}|^2 \leq C 2^{-2uj} \tag{12}$$

which was previously obtained by Yves Meyer (personal communication) using an alternative proof.

If  $p = +\infty$ , and if  $1 \leq p < +\infty$ , then Theorem 2.1 improves previous results of, respectively, [3] and [5]; up to now, the converse part required a uniform regularity assumption  $f \in B_p^{\epsilon, p}$  for and  $\epsilon > 0$ , which turns out to be unnecessary. Note also that, if  $f$  satisfies (7), then  $f \in T_u^p(x_0)$  for any  $p < \infty$ . This is in sharp contrast with the two-microlocal condition obtained in [3] as a consequence of  $C^u(x_0)$  regularity which does not imply any  $T_u^p(x_0)$  regularity result (or even that  $f$  locally coincides with a function).

If  $p \neq 2$ , then (6) is not a local  $l^p$  condition on the wavelet coefficients; however, the embeddings between Sobolev and Besov spaces supply the following conditions which are easier to use in practice:

If  $p \geq 2$ , then  $L^p \hookrightarrow B_p^{0, p}$ ; thus if  $f \in T_u^p(x_0)$  for  $p \geq 2$ , then  $\sum_{\lambda' \subset 3\lambda_j(x_0)} 2^{-d(j'-j)} |c_{\lambda'}|^p \leq C 2^{-puj}$ . Similarly, if  $p \leq 2$ , then  $B_p^{0, p} \hookrightarrow L^p$ ; thus if  $\sum_{\lambda' \subset 3\lambda_j(x_0)} 2^{-d(j'-j)} |c_{\lambda'}|^p \leq C 2^{-puj}$ , then  $f \in T_u^p(x_0)$ .

The two-microlocal condition  $C^{\alpha, -\alpha}(x_0)$  is ‘far’ from the Hölder condition  $C^{\alpha}(x_0)$  in the sense that it can be satisfied by distributions which are not functions. However, it is ‘close’ if a uniform regularity condition holds; indeed, let  $\alpha > \epsilon > 0$ ; if  $f \in C^{\alpha, -\alpha}(x_0) \cap C^{\epsilon}(\mathbb{R}^d)$ , then  $\forall \beta < \alpha$ ,  $f \in C^{\beta}(x_0)$ , see [3]. The following result shows that  $T_u^p$  regularity is ‘farther’ from Hölder regularity under this respect.

**Proposition 3.1.** *Let  $f \in T_{\alpha}^p(x_0) \cap C^{\epsilon}(\mathbb{R}^d)$ , with  $\alpha > \epsilon + (d/p)$  and let  $\beta = \alpha\epsilon p / (\epsilon p + d)$ ; then  $f \in C^{\beta}(x_0)$  and this result is optimal.*

The proof is similar to the proof of the converse part in Theorem 2.1; the only difference lies in the estimate of  $\sup |\Delta_j f|$  on  $B = B(x_0, 2^{-J})$  for  $j \geq J$ . The uniform regularity assumption implies that  $\sup(|\Delta_j f|) \leq C 2^{-\epsilon j}$ ; the  $T_{\alpha}^p(x_0)$  assumption implies that  $\|\Delta_j f\|_{L^p(B)} \leq C 2^{-\alpha j}$ , which, using Bernstein’s inequalities, implies that  $\sup |\Delta_j f| \leq C 2^{-\alpha j} 2^{dj/p}$ . The conclusion follows by picking the best of these two estimates according to the value of  $j$ .

The purpose of the multifractal analysis of a function  $f$  is to determine the Hausdorff dimensions of the sets of points where  $f$  has a given pointwise regularity. Up to now, this was performed mainly in the context of Hölder pointwise regularity. We now give a result for  $T_u^p$  regularity. In that case one wishes to determine the  $p$ -spectrum  $D_p^f(H) = \dim(\{x: h_f^p(x) = H\})$  (where  $\dim$  denotes the Hausdorff dimension). Upper bounds on the  $p$ -spectrum can be derived in terms of the following quantities. Let

$$S_f^\lambda(p) = \left( \int_\lambda \left( \sum_{\lambda' \subset \lambda} |c_{\lambda'}|^2 1_{\lambda'}(x) \right)^{p/2} dx \right)^{1/p},$$

$$\eta_f^p(q) = \lim_{R \rightarrow +\infty} \liminf_{j \rightarrow +\infty} \frac{\log(2^{d(q/p-1)j} \sum_{\lambda \in \Lambda_j \cap B(0,R)} (S_f^\lambda(p))^q)}{\log(2^{-j})}.$$

**Theorem 3.2.** *Let  $f \in L_{\text{loc}}^p$ ; then  $D_p^f(H) \leq \inf_{q \neq 0} (d - \eta_f^p(q) + Hq)$ .*

**Sketch of proof.** It follows from Theorem 2.1 that

$$h_f^p(x_0) = -\frac{d}{p} + \liminf_{j \rightarrow +\infty} \left( \frac{-1}{j} \log_2 \left( \sup_{\lambda' \in \text{adj}(\lambda)} S_f^{\lambda'}(p) \right) \right) \quad (13)$$

where  $\text{adj}(\lambda)$  denotes the  $3^d$  dyadic cubes of same width as  $\lambda$  and such that  $\bar{\lambda} \cap \bar{\lambda}' \neq \emptyset$ . The proof of Theorem 3.2 is exactly the same as the upper bound for the Hölder spectrum, see [4], since the only property used in the derivation of this upper bound is a formula similar to (13). Note that, in Theorem 3.2, no global regularity assumption is needed since, in [4], this assumption is needed only to insure the validity of the formula corresponding to (13), but not in the proof of the upper bound.

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## References

- [1] A. Arneodo, B. Audit, N. Decoster, J.-F. Muzy, C. Vaillant, Wavelet-based multifractal formalism: applications to DNA sequences, satellite images of the cloud structure and stock market data, in: A. Bunde, J. Kropp, H.J. Schellnhuber (Eds.), *The Science of Disasters*, Springer, 2002, pp. 27–102.
- [2] A.P. Calderon, A. Zygmund, Local properties of solutions of elliptic partial differential equations, *Studia Math.* 20 (1961) 171–227.
- [3] S. Jaffard, Pointwise smoothness, two-microlocalization and wavelet coefficients, *Publ. Mat.* 35 (1991) 155–168.
- [4] S. Jaffard, Wavelet techniques in multifractal analysis, in: M. Lapidus, M. van Frankenhuysen, *Fractal Geometry and Applications: A Jubilee of Benoit Mandelbrot*, Proc. Sympos. Pure Math., American Mathematical Society, 2004, in press.
- [5] S. Jaffard, C. Melot, Wavelet analysis of fractal boundaries, Preprint, 2004.
- [6] Y. Meyer, *Ondelettes et opérateurs*, Hermann, 1990.