# On a Liouville phenomenon for entire weak supersolutions of elliptic partial differential equations 

# Autour d'un phénomène de Liouville pour les sursolutions entières faibles d'équations aux derivées partielles elliptiques 

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#### Abstract

We study a new Liouville-type phenomenon for entire weak supersolutions of elliptic partial differential equations of the form $A(u)=0$ on $\mathbb{R}^{n}, n \geqslant 2$. Typical examples of the operator $A(u)$ are the $p$-Laplacian for $p>1$, the mean curvature operator, and their well-known modifications.


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## Résumé

Ce travail est consacré à l'étude d'un nouveau phénomène de type de Liouville pour les sursolutions entières faibles d'équations aux derivées partielles elliptiques de la forme $A(u)=0$ sur $\mathbb{R}^{n}, n \geqslant 2$. Des exemples typiques de l'opérateur $A(u)$ sont le $p$-laplacien pour $p>1$, l'opérateur de courbure moyenne, et leurs modifications bien connues. © 2006 Elsevier Masson SAS. All rights reserved.

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## 1. Introduction

Liouville's well-known theorem says that any superharmonic function on $\mathbb{R}^{2}$ bounded below by a constant is itself a constant. On the other hand it is also well known that for $n \geqslant 3$ there exist non-constant superharmonic functions on $\mathbb{R}^{n}$ bounded below by a constant. The purpose of this work is to determine for $n \geqslant 3$ the 'sharp distance at infinity' between the non-constant superharmonic functions on $\mathbb{R}^{n}$ bounded below by a constant and this constant itself in

[^0]the form of a theorem of Liouville type and to characterize basic properties of quasilinear elliptic partial differential operators which make it possible to obtain such a theorem for supersolutions of quasilinear elliptic partial differential equations of the form
\[

$$
\begin{equation*}
A(u)=0 \tag{1}
\end{equation*}
$$

\]

on $\mathbb{R}^{n}, n \geqslant 2$. Typical examples of the operator $A(u)$ are the $p$-Laplacian

$$
\begin{equation*}
\Delta_{p}(u):=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right), \quad p>1, \tag{2}
\end{equation*}
$$

its well-known modification (see, e.g., [8, p. 155])

$$
\begin{equation*}
\tilde{\Delta}_{p}(u):=\sum_{i=1}^{n} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p-2} \frac{\partial u}{\partial x_{i}}\right), \quad p>1, \tag{3}
\end{equation*}
$$

the mean curvature operator

$$
\begin{equation*}
\Xi(u):=\operatorname{div}\left(\frac{\nabla u}{\sqrt{1+|\nabla u|^{2}}}\right) \tag{4}
\end{equation*}
$$

and its well-known modifications.
Note that a Liouville theorem for solutions of linear uniformly elliptic second-order partial differential equations on $\mathbb{R}^{n}, n>2$, was first obtained, as a direct consequence of a Harnack inequality, in [1] under some continuity assumptions on the coefficients of the equations and in [12] without continuity assumptions on the coefficients of the equations. In the case of quasilinear uniformly elliptic second-order partial differential equations on $\mathbb{R}^{n}, n \geqslant 2$, a Liouville theorem is a direct consequence of a Harnack inequality first obtained in [14]. Note also that a Liouville theorem for mappings of $\mathbb{R}^{n}, n>2$, with bounded distortion was first obtained in [13] by using the same Harnack inequality from [14]. Finally, in the case of linear uniformly elliptic second-order partial differential equations on $\mathbb{R}^{2}$, a Liouville theorem is a direct consequence of a Harnack inequality first obtained in [7].

## 2. Definitions

Let $A(u)$ be a differential operator defined formally by

$$
\begin{equation*}
A(u)=\sum_{i=1}^{n} \frac{\mathrm{~d}}{\mathrm{~d} x_{i}} A_{i}(x, u, \nabla u) . \tag{5}
\end{equation*}
$$

Here and in what follows, $n \geqslant 2$. We assume that the functions $A_{i}(x, \eta, \xi), i=1, \ldots, n$, satisfy the usual Carathéodory conditions on $\mathbb{R}^{n} \times \mathbb{R}^{1} \times \mathbb{R}^{n}$; namely, they are continuous in $\eta$ and $\xi$ for almost all $x \in \mathbb{R}^{n}$ and measurable in $x$ for any $\eta \in \mathbb{R}^{1}$ and $\xi \in \mathbb{R}^{n}$.

Definition 1. Let $\alpha>1$ be a given number. The operator $A(u)$ given by (5) belongs to the class $\mathcal{A}(\alpha)$ if for all $\eta \in \mathbb{R}^{1}$, all $\xi, \psi \in \mathbb{R}^{n}$, and almost all $x \in \mathbb{R}^{n}$ the following two inequalities hold:

$$
\begin{equation*}
0 \leqslant \sum_{i=1}^{n} \xi_{i} A_{i}(x, \eta, \xi) \tag{6}
\end{equation*}
$$

with equality only if $\xi=0$, and

$$
\begin{equation*}
\left|\sum_{i=1}^{n} \psi_{i} A_{i}(x, \eta, \xi)\right|^{\alpha} \leqslant \mathcal{K}|\psi|^{\alpha}\left(\sum_{i=1}^{n} \xi_{i} A_{i}(x, \eta, \xi)\right)^{\alpha-1} \tag{7}
\end{equation*}
$$

with $\mathcal{K}$ a certain positive constant.
It is easy to see that condition (7) is fulfilled whenever the inequality

$$
\begin{equation*}
\left(\sum_{i=1}^{n} A_{i}^{2}(x, \eta, \xi)\right)^{\alpha / 2} \leqslant \mathcal{K}\left(\sum_{i=1}^{n} \xi_{i} A_{i}(x, \eta, \xi)\right)^{\alpha-1} \tag{8}
\end{equation*}
$$

holds for all $\eta \in \mathbb{R}^{1}$, all $\xi, \psi \in \mathbb{R}^{n}$, and almost all $x \in \mathbb{R}^{n}$. Hence, the operator $A(u)$ given by (5) and satisfying conditions (6) and (8) belongs to the class $\mathcal{A}(\alpha)$.

Remark 1. Conditions (7) and (8) on the behavior of the coefficients of partial differential operators were introduced in [10].

It is not difficult to verify that for any given $p>1$ the differential operators (2) and (3) as well as the differential operator $A(u)$ given by (5) and satisfying the well-known growth conditions

$$
\begin{equation*}
\left(\sum_{i=1}^{n} A_{i}^{2}(x, \eta, \xi)\right)^{1 / 2} \leqslant \mathcal{K}_{1}|\xi|^{p-1} \tag{9}
\end{equation*}
$$

and

$$
\begin{equation*}
|\xi|^{p} \leqslant \mathcal{K}_{2} \sum_{i=1}^{n} \xi_{i} A_{i}(x, \eta, \xi) \tag{10}
\end{equation*}
$$

with $\mathcal{K}_{1}, \mathcal{K}_{2}$ positive constants, belong to the class $\mathcal{A}(\alpha)$ with $\alpha=p$.
It is also easy to see that linear divergent elliptic partial differential operators of the form

$$
\begin{equation*}
L:=\sum_{i, j=1}^{n} \frac{\partial}{\partial x_{j}}\left(a_{i j}(x) \frac{\partial}{\partial x_{i}}\right) \tag{11}
\end{equation*}
$$

with $a_{i j}(x)$ measurable bounded coefficients and with the (possibly non-uniformly) positive-definite quadratic form

$$
\begin{equation*}
\sum_{i, j=1}^{n} a_{i j}(x) \xi_{i} \xi_{j} \tag{12}
\end{equation*}
$$

belong to the class $\mathcal{A}(\alpha)$ with $\alpha=2$ but do not satisfy condition (10) for any fixed $p>1$.
In connection with this we give another example of an operator that belongs to the class $\mathcal{A}(\alpha)$ with a certain $\alpha>1$ but does not satisfy condition (10). Let $a(x, \eta, \xi)$ be a positive bounded function that satisfies the Carathéodory conditions on $\mathbb{R}^{n} \times \mathbb{R}^{1} \times \mathbb{R}^{n}$. It is easy to see that for a given $p>1$ the weighted $p$-Laplacian

$$
\begin{equation*}
\bar{\Delta}_{p}(u):=\operatorname{div}\left(a(x, u, \nabla u)|\nabla u|^{p-2} \nabla u\right) \tag{13}
\end{equation*}
$$

belongs to the class $\mathcal{A}(\alpha)$ with $\alpha=p$ but does not satisfy condition (10) for any fixed $p>1$ if the function $a(x, \eta, \xi)$ is only assumed to be positive.

It can happen that an operator $A(u)$ given by (5) belongs simultaneously to several different classes $\mathcal{A}(\alpha)$. For example, the mean curvature operator $\Xi(u)$ given by (4) belongs to the classes $\mathcal{A}(\alpha)$ for all $1<\alpha \leqslant 2$; similarly its modification for $p \geqslant 2$,

$$
\begin{equation*}
\Xi_{p}(u):=\operatorname{div}\left(\frac{|\nabla u|^{p-2} \nabla u}{\sqrt{1+|\nabla u|^{2}}}\right), \tag{14}
\end{equation*}
$$

belongs to the classes $\mathcal{A}(\alpha)$ for all $\alpha \in(p-1, p]$ and $p \geqslant 2$. Obviously, operators given by (4) and (14) do not satisfy conditions (9)-(10) for any fixed $p \geqslant 1$.

Definition 2. Let $\alpha>1$ be a given number, and let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$. A measurable function $u: \mathbb{R}^{n} \rightarrow \mathbb{R}^{1}$ is called an entire weak supersolution of Eq. (1) on $\mathbb{R}^{n}$ if $u \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{n}\right),|\nabla u| \in L_{\mathrm{loc}}^{\alpha}\left(\mathbb{R}^{n}\right)$, and the integral inequality

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \sum_{i=1}^{n} \varphi_{x_{i}} A_{i}(x, u, \nabla u) \mathrm{d} x \geqslant 0 \tag{15}
\end{equation*}
$$

holds for every non-negative function $\varphi \in W^{1, \alpha}\left(\mathbb{R}^{n}\right)$ with compact support.

## 3. Results

Theorem 1. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $\alpha \geqslant n$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant. Then $u(x)$ is a constant on $\mathbb{R}^{n}$.

Theorem 2. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant $c$ and such that $u \in L_{\text {loc }}^{\infty}\left(\mathbb{R}^{n}\right)$. Then either $u(x)=c$ on $\mathbb{R}^{n}$ or the relation

$$
\begin{equation*}
\liminf _{r \rightarrow+\infty}\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)\right]^{\frac{n-\alpha}{\alpha-1-\nu}}=+\infty \tag{16}
\end{equation*}
$$

holds with any fixed $v \in(0, \alpha-1)$.
Theorem 3. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant $c$. Then either $u(x)=c$ on $\mathbb{R}^{n}$ or the relation

$$
\begin{equation*}
\liminf _{r \rightarrow+\infty} r^{-\alpha} \int_{r \leqslant|x| \leqslant 2 r}(u(x)-c)^{\alpha-1-v} \mathrm{~d} x=+\infty \tag{17}
\end{equation*}
$$

holds with any fixed $v \in(0, \alpha-1)$.
Due to the arbitrariness of the constant $c$ in Theorems 2 and 3, the statements of these theorems can be reformulated in a slightly different form.

Theorem $2^{\prime}$. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant and such that $u \in L_{\mathrm{loc}}^{\infty}\left(\mathbb{R}^{n}\right)$. Then either $u(x)$ is a constant on $\mathbb{R}^{n}$ or relation (16) holds with any fixed real number $c$ such that $u(x) \geqslant c$ on $\mathbb{R}^{n}$ and any fixed $v \in(0, \alpha-1)$.

Theorem 3'. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant. Then either $u(x)$ is a constant on $\mathbb{R}^{n}$ or relation (17) holds with any fixed real number $c$ such that $u(x) \geqslant c$ on $\mathbb{R}^{n}$ and any fixed $\nu \in(0, \alpha-1)$.

Remark 2. It is important to note that for any given $n \geqslant 2$ and $\alpha>1$ such that $n>\alpha$ the function

$$
\begin{equation*}
u(x)=\left(1+|x|^{\frac{\alpha}{\alpha-1}}\right)^{\frac{\alpha-n}{\alpha}} \tag{18}
\end{equation*}
$$

is an entire weak supersolution of the equation

$$
\begin{equation*}
\Delta_{p}(u)=0 \tag{19}
\end{equation*}
$$

with $p=\alpha$ that is bounded below and is such that relations (16) and (17) hold with any fixed $v \in(0, \alpha-1)$ and, at the same time, the relations

$$
\begin{equation*}
\lim _{r \rightarrow+\infty}\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-0)\right]^{r^{\frac{n-\alpha}{\alpha-1}}}=C_{1} \tag{20}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{r \rightarrow+\infty} r^{-\alpha} \int_{r \leqslant|x| \leqslant 2 r}(u(x)-0)^{\alpha-1} \mathrm{~d} x=C_{2}, \tag{21}
\end{equation*}
$$

with $C_{1}, C_{2}$ certain positive constants, also hold.

Remark 3. The results of this work were announced in [5]. To prove these results we further develop an approach that was proposed for solving similar problems in [6].

Remark 4. The results of Theorem 1 are new only for $\alpha=n$. Similar results to those of Theorem 1 for entire weak continuous supersolutions of (1) on $\mathbb{R}^{n}$ for $\alpha=n$ were first obtained in [11]. For $\alpha>n$, the results of Theorem 1 for entire weak supersolutions of (1) on $\mathbb{R}^{n}$, which in this case are continuous on $\mathbb{R}^{n}$ by the well-known Sobolev imbedding theory, were also first obtained in [11]. Here, we give a new proof of these results from [11] by developing an approach from [6] which does not explicitly use the continuity of entire weak supersolutions of (1) on $\mathbb{R}^{n}$.

Remark 5. In the case when $\alpha=p$ and $A(u)=\Delta_{p}(u)$, Theorem 1 coincides with well-known Liouville-type theorems for entire superharmonic and $p$-superharmonic functions locally bounded on $\mathbb{R}^{n}$ (see, e.g., [2, p. 68] and [3, p. 179]). Also, in this case, the results of Theorems 2 and 3 correlate well with certain results in the theory of entire superharmonic and $p$-superharmonic functions (see, e.g., [2, pp. 131, 139] and [3, pp. 133, 135]).

## 4. Proofs

Proof of Theorem 2. The statement of Theorem 2 follows immediately from Theorem 3. In fact, let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant $c$, i.e., $u(x) \geqslant c$ on $\mathbb{R}^{n}$, and such that $u \in L_{\text {loc }}^{\infty}\left(\mathbb{R}^{n}\right)$. Hence, by Theorem 3, either $u(x)=c$ on $\mathbb{R}^{n}$ or relation (17) holds with any fixed $v \in(0, \alpha-1)$. Further, via the trivial inequality

$$
\begin{equation*}
r^{-\alpha} \int_{r \leqslant|x| \leqslant 2 r}(u(x)-c)^{\alpha-1-v} \mathrm{~d} x \leqslant r^{-\alpha}\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)^{\alpha-1-\nu}\right] \int_{r \leqslant|x| \leqslant 2 r} \mathrm{~d} x, \tag{22}
\end{equation*}
$$

which obviously holds for any $r>0$, it follows from (17) that

$$
\begin{equation*}
\liminf _{r \rightarrow+\infty}\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)^{\alpha-1-\nu}\right] r^{n-\alpha}=+\infty . \tag{23}
\end{equation*}
$$

Then, since

$$
\begin{equation*}
\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)^{\alpha-1-v} \leqslant\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)\right]^{\alpha-1-v} \tag{24}
\end{equation*}
$$

and

$$
\begin{equation*}
\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)\right]^{\alpha-1-v} r^{n-\alpha}=\left(\left[\sup _{r \leqslant|x| \leqslant 2 r}(u(x)-c)\right]^{\frac{n-\alpha}{\alpha-1-v}}\right)^{\alpha-1-v}, \tag{25}
\end{equation*}
$$

the validity of (16) follows immediately from that of (23) and (25).
In what follows, a 'smooth' function is a $C^{\infty}$-function on $\mathbb{R}^{n}, B(r)$ is an open ball on $\mathbb{R}^{n}$ of radius $r>0$ centered at the origin, and $\overline{B(r)}$ is the closure of $B(r)$.

Proof of Theorem 3. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $n>\alpha$. Let the operator $A(u)$ given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant $c$, i.e., $u(x) \geqslant c$ on $\mathbb{R}^{n}$. Let $r$ and $\varepsilon$ be positive numbers, and let $\zeta: \mathbb{R}^{n} \rightarrow[0,1]$ be a smooth function which equals 1 on $\overline{B(r)}$ and 0 outside $B(2 r)$. Substituting, without loss of generality, $\varphi(x)=(u(x)-c+\varepsilon)^{-v} \zeta^{\alpha}(x)$ as a test function in inequality (15), where $\nu \in(0, \alpha-1)$ is arbitrary, and integrating by parts, we find

$$
\begin{align*}
& \alpha \int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} \zeta_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v} \zeta^{\alpha-1} \mathrm{~d} x \\
& \quad \geqslant v \int_{B(2 r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \zeta^{\alpha} \mathrm{d} x . \tag{26}
\end{align*}
$$

Estimating the left-hand side of (26) by using condition (7) on the coefficients of the operator $A(u)$, we have

$$
\begin{align*}
& \alpha \mathcal{K}^{1 / \alpha} \int_{B(2 r) \backslash B(r)}\left(\sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)\right)^{(\alpha-1) / \alpha}|\nabla \zeta|(u-c+\varepsilon)^{-v} \zeta^{\alpha-1} \mathrm{~d} x \\
& \quad \geqslant\left|\alpha \int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} \zeta_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v} \zeta^{\alpha-1} \mathrm{~d} x\right| \tag{27}
\end{align*}
$$

Further, estimating the left-hand side of (27) by Hölder's inequality, we arrive at

$$
\begin{align*}
& \alpha \mathcal{K}^{1 / \alpha}\left(\int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{\alpha-1-v} \mathrm{~d} x\right)^{1 / \alpha} \\
& \quad \times\left(\int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \zeta^{\alpha} \mathrm{d} x\right)^{(\alpha-1) / \alpha} \\
& \geqslant\left|\alpha \int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} \zeta_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v} \zeta^{\alpha-1} \mathrm{~d} x\right| \tag{28}
\end{align*}
$$

In turn, (26) and (28) imply the inequality

$$
\begin{align*}
& \alpha \mathcal{K}^{1 / \alpha}\left(\int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{\alpha-1-v} \mathrm{~d} x\right)^{1 / \alpha} \\
& \quad \times\left(\int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-\nu-1} \zeta^{\alpha} \mathrm{d} x\right)^{(\alpha-1) / \alpha} \\
& \geqslant v \int_{B(2 r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-\nu-1} \zeta^{\alpha} \mathrm{d} x \tag{29}
\end{align*}
$$

and, therefore, the inequality

$$
\begin{equation*}
\alpha^{\alpha} \mathcal{K} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{\alpha-1-v} \mathrm{~d} x \geqslant v^{\alpha} \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x . \tag{30}
\end{equation*}
$$

It is easy to see that the right-hand side of (30) increases monotonically if $\varepsilon>0$ decreases strongly monotonically to zero. Therefore, it follows from (30) that the inequality

$$
\begin{equation*}
\alpha^{\alpha} \mathcal{K} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{-\nu+\alpha-1} \mathrm{~d} x \geqslant \nu^{\alpha} \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\delta)^{-\nu-1} \mathrm{~d} x \tag{31}
\end{equation*}
$$

holds with any $\delta>0$ and any $\varepsilon \in(0, \delta]$. Since for any sequence $\varepsilon_{k}>0$ monotonically decreasing to zero as $k \rightarrow+\infty$ the sequence of functions

$$
\begin{equation*}
\Phi_{k}(x):=|\nabla \zeta|^{\alpha}\left(u-c+\varepsilon_{k}\right)^{\alpha-1-v} \tag{32}
\end{equation*}
$$

measurable on $\mathbb{R}^{n}$ converges a.e. on $\mathbb{R}^{n}$ to the function

$$
\begin{equation*}
\Phi(x):=|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \tag{33}
\end{equation*}
$$

measurable on $\mathbb{R}^{n}$, since for sufficiently large $k$

$$
\begin{equation*}
\left|\Phi_{k}(x)\right| \leqslant|\nabla \zeta|^{\alpha}(u-c+1)^{\alpha-1-v} \tag{34}
\end{equation*}
$$

on $\mathbb{R}^{n}$, and since the function

$$
\begin{equation*}
|\nabla \zeta|^{\alpha}(u-c+1)^{\alpha-1-v} \tag{35}
\end{equation*}
$$

is locally integrable on $\mathbb{R}^{n}$, then, by Lebesgue's theorem (see, e.g., [4, p. 303]), for $\varepsilon=\varepsilon_{k}>0$ monotonically decreasing to zero we can pass to the limit as $k \rightarrow+\infty$ on the left-hand side of (31). As a result, we obtain the inequality

$$
\begin{equation*}
\alpha^{\alpha} \mathcal{K} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x \geqslant \nu^{\alpha} \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\delta)^{-v-1} \mathrm{~d} x, \tag{36}
\end{equation*}
$$

which holds with any $\delta>0$. Then, for any $r>0$ and any sequence $\varepsilon_{k}>0$ monotonically decreasing to zero as $k \rightarrow+\infty$, it follows from (36), by letting $\delta=\varepsilon_{k}$ and

$$
\begin{equation*}
\Psi_{k}(x):=\sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)\left(u-c+\varepsilon_{k}\right)^{-v-1} \tag{37}
\end{equation*}
$$

that the sequence of integrals

$$
\begin{equation*}
\int_{B(r)} \Psi_{k}(x) \mathrm{d} x \tag{38}
\end{equation*}
$$

is bounded above by the positive constant

$$
\begin{equation*}
c_{1}=\mathcal{K}\left(\frac{\alpha}{v}\right)^{\alpha} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x, \tag{39}
\end{equation*}
$$

which does not depend on $\varepsilon_{k}$. Hence, since

$$
\begin{equation*}
\Psi_{1}(x) \leqslant \Psi_{2}(x) \leqslant \cdots \leqslant \Psi_{k}(x) \leqslant \cdots \tag{40}
\end{equation*}
$$

on $\mathbb{R}^{n}$, then by Beppo Levi's theorem (see, e.g., [4, p. 305]), for any $r>0$ there exists a function $\Theta_{r}: B(r) \rightarrow \mathbb{R}^{1}$ integrable on $B(r)$ and such that the sequence of functions $\Psi_{k}(x)$ converges a.e. to $\Theta_{r}(x)$ on $B(r)$ and

$$
\begin{equation*}
\lim _{k \rightarrow+\infty} \int_{B(r)} \Psi_{k}(x) \mathrm{d} x=\int_{B(r)} \Theta_{r}(x) \mathrm{d} x \tag{41}
\end{equation*}
$$

Further, it is easy to see that the family of functions $\left\{\Theta_{r}\right\}_{r>0}$ uniquely determines a function $\psi: \mathbb{R}^{n} \rightarrow \mathbb{R}^{1}$ which is non-negative, measurable, locally integrable on $\mathbb{R}^{n}$ and is such that $\Psi(x)=\Theta_{r}(x)$ on $B(r)$ for all $r>0$. Therefore, the sequence of functions $\Psi_{k}(x)$ given by (37) converges a.e. to $\Psi(x)$ on $\mathbb{R}^{n}$ for any sequence $\varepsilon_{k}>0$ monotonically decreasing to zero as $k \rightarrow+\infty$. Then, by choosing $\delta=\varepsilon_{k}$ in (36), where the sequence $\varepsilon_{k}>0$ converges monotonically to zero as $k \rightarrow+\infty$, and passing to the limit on the right-hand side of (36), we find, due to (41), the inequality

$$
\begin{equation*}
\alpha^{\alpha} \mathcal{K} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x \geqslant v^{\alpha} \int_{B(r)} \Psi(x) \mathrm{d} x . \tag{42}
\end{equation*}
$$

We divide the rest of the proof into three cases according to the behavior of the right-hand side of (42), which can monotonically approach zero, $+\infty$, or some positive number $I$ as $r$ strongly monotonically approaches $+\infty$.

If the right-hand side of (42) approaches zero as $r \rightarrow+\infty$, then, due to the non-negativity of the function $\Psi(x)$, we have that $\Psi(x)=0$ on $\mathbb{R}^{n}$. Further, since by (37) and (40) the inequality

$$
\begin{equation*}
\sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)\left(u-c+\varepsilon_{k}\right)^{-\nu-1} \leqslant \Psi(x) \tag{43}
\end{equation*}
$$

holds on $\mathbb{R}^{n}$ for any sequence $\varepsilon_{k}>0$ monotonically decreasing to zero as $k \rightarrow+\infty$, then, again, due to the nonnegativity of the left-hand side of (43), we obtain that

$$
\begin{equation*}
\sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)\left(u-c+\varepsilon_{k}\right)^{-v-1}=0 \tag{44}
\end{equation*}
$$

on $\mathbb{R}^{n}$. Hence, by condition (6) on the coefficients of the operator $A(u)$, the supersolution $u(x)=$ const. on $\mathbb{R}^{n}$, and, therefore, either $u(x)=c$ on $\mathbb{R}^{n}$ or relation (17) holds with any fixed $v \in(0, \alpha-1)$.

If the right-hand side of (42) approaches $+\infty$ as $r \rightarrow+\infty$, then, due to monotonicity, (42) yields that

$$
\begin{equation*}
\liminf _{r \rightarrow+\infty} \int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x=+\infty . \tag{45}
\end{equation*}
$$

Finally, if the right-hand side of (42) monotonically approaches a certain positive number $I$ as $r$ approaches $+\infty$, i.e.,

$$
\begin{equation*}
\lim _{r \rightarrow+\infty} v^{\alpha} \int_{B(r)} \Psi(x) \mathrm{d} x=I>0, \tag{46}
\end{equation*}
$$

we again consider inequality (29), just noting here that, due to monotonicity,

$$
\begin{equation*}
\int_{B\left(2 r_{k}\right) \backslash B\left(r_{k}\right)} \Psi(x) \mathrm{d} x \rightarrow 0 \tag{47}
\end{equation*}
$$

for any sequence $r_{k}>0$ such that $r_{k} \rightarrow+\infty$. First, we have from (29) the inequality

$$
\begin{align*}
& \alpha \mathcal{K}^{1 / \alpha}\left(\int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{\alpha-1-v} \mathrm{~d} x\right)^{1 / \alpha} \\
& \quad \times\left(\int_{B(2 r) \backslash B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x\right)^{(\alpha-1) / \alpha} \\
& \geqslant v \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{48}
\end{align*}
$$

In (48), let $\varepsilon=\varepsilon_{k}>0$ converge monotonically to zero as $k \rightarrow+\infty$. Then, by Lebesgue's theorem (see, e.g., [4, p. 303]), we can pass to the limit on both sides of (48). Namely, we know from the above that for any sequence $\varepsilon_{k}>0$ monotonically decreasing to zero as $k \rightarrow+\infty$ the sequences of functions $\Phi_{k}(x)$ and $\Psi_{k}(x)$ measurable and locally integrable on $\mathbb{R}^{n}$ and given, respectively, by (32) and (37), converge a.e. on $\mathbb{R}^{n}$, respectively, to the functions $\Phi(x)$ and $\Psi(x)$ measurable and locally integrable on $\mathbb{R}^{n}$. Further, arguing as above and letting $\varepsilon=\varepsilon_{k}>0$ monotonically decrease to zero as $k \rightarrow+\infty$, by Lebesgue's theorem (see, e.g., [4, p. 303]) we can pass to the limit on both sides of (48). As a result, we arrive at the inequality

$$
\begin{equation*}
\alpha \mathcal{K}^{1 / \alpha}\left(\int_{B(2 r) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x\right)^{1 / \alpha}\left(\int_{B(2 r) \backslash B(r)} \Psi(x) \mathrm{d} x\right)^{(\alpha-1) / \alpha} \geqslant v \int_{B(r)} \Psi(x) \mathrm{d} x . \tag{49}
\end{equation*}
$$

In (49), for $r=r_{k}>0$ monotonically increasing to $+\infty$, by passing to the limit as $r_{k} \rightarrow+\infty$, we obtain from (46), (47), and (49) that

$$
\begin{equation*}
\lim _{r_{k} \rightarrow+\infty} \int_{B\left(2 r_{k}\right) \backslash B\left(r_{k}\right)}|\nabla \zeta|^{\alpha}(u-c)^{\alpha-1-v} \mathrm{~d} x=+\infty . \tag{50}
\end{equation*}
$$

Thus, due to the arbitrariness in the choice of the sequence $r_{k}$ in (50), we again arrive at relation (45).
Now, without loss of generality, we choose in (45) the function $\zeta(x)$ in the form $\zeta(x)=\psi(|x| /(2 r))$, where $\psi:[0,+\infty) \rightarrow[0,1]$ is a smooth function that equals 1 on $[0,1 / 2]$ and 0 on $[1,+\infty)$ and is such that the inequality

$$
\begin{equation*}
|\nabla \zeta| \leqslant c_{2} r^{-1} \tag{51}
\end{equation*}
$$

holds on $\mathbb{R}^{n}$ with a certain positive constant $c_{2}$ for an arbitrary $r>0$. Relation (17) then follows immediately from (45) and (51).

Proof of Theorem 1. Let $n \geqslant 2$ and $\alpha>1$ be given numbers such that $\alpha \geqslant n$. Let the operator $A$ (u) given by (5) belong to the class $\mathcal{A}(\alpha)$, and let $u(x)$ be an entire weak supersolution of (1) on $\mathbb{R}^{n}$ bounded below by a constant $c$, i.e., $u(x) \geqslant c$ on $\mathbb{R}^{n}$. Let $r, R$, and $\varepsilon$ be positive numbers such that $R>r$, and let $\zeta: \mathbb{R}^{n} \rightarrow[0,1]$ be a smooth function which equals 1 on $\overline{B(r)}$ and 0 outside $B(R)$. Substituting, without loss of generality, $\varphi(x)=(u(x)-c+\varepsilon)^{-v} \zeta^{\alpha}(x)$ as a test function in inequality (15), where $v>\alpha-1$ is an arbitrary positive number, and integrating by parts, we have the inequality

$$
\begin{align*}
& \alpha \int_{B(R) \backslash B(r)} \sum_{i=1}^{n} \zeta_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v} \zeta^{\alpha-1} \mathrm{~d} x \\
& \quad \geqslant v \int_{B(R)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \zeta^{\alpha} \mathrm{d} x \tag{52}
\end{align*}
$$

Further, we repeat the proof of Theorem 3 word for word from (26) to (30). As a result, we arrive at the inequality

$$
\begin{equation*}
\alpha^{\alpha} \mathcal{K} \int_{B(R) \backslash B(r)}|\nabla \zeta|^{\alpha}(u-c+\varepsilon)^{\alpha-1-v} \mathrm{~d} x \geqslant v^{\alpha} \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{53}
\end{equation*}
$$

It follows immediately from (53) that the inequality

$$
\begin{equation*}
\alpha^{\alpha} \varepsilon^{\alpha-1-v} \mathcal{K} \int_{B(R) \backslash B(r)}|\nabla \zeta|^{\alpha} \mathrm{d} x \geqslant v^{\alpha} \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{54}
\end{equation*}
$$

holds with any fixed $\varepsilon>0$ and $v>\alpha-1$.
Now, first let $\alpha>n$. In (54), choosing $R=2 r$ and the function $\zeta(x)$ in the form $\zeta(x)=\psi(|x| / R)$, where $\psi:[0,+\infty) \rightarrow[0,1]$ is a smooth function that equals 1 on $[0,1 / 2]$ and 0 on $[1,+\infty)$ and is such that the inequality (51) holds on $\mathbb{R}^{n}$ with a certain positive constant $c_{2}$ for an arbitrary $R>0$, we obtain from (51) and (54) the inequality

$$
\begin{equation*}
c_{3} r^{n-\alpha} \geqslant \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{55}
\end{equation*}
$$

which holds with a certain positive constant $c_{3}$ that does not depend on $r$. Passing to the limit as $r \rightarrow+\infty$ in (55), we find, due to the non-negativity of the integrand, that the equality

$$
\begin{equation*}
\sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1}=0 \tag{56}
\end{equation*}
$$

holds on $\mathbb{R}^{n}$, and, therefore, by condition (6) on the coefficients of the operator $A(u)$, that $u(x)=$ const. on $\mathbb{R}^{n}$.
If $\alpha=n$, we choose in (54) the function $\zeta(x)$ in the form $\zeta(x)=\psi\left(\frac{\ln (|x| / r)}{\ln (R / r)}\right)$ with arbitrary $R>r>1$, where $\psi:[-\infty,+\infty) \rightarrow[0,1]$ is a smooth function which equals 1 on $[-\infty, 0]$ and 0 on $[1,+\infty)$. It is not difficult to understand (see, e.g., [9, p. 12]) that the inequality

$$
\begin{equation*}
|\nabla \zeta(x)| \leqslant \frac{c_{4}}{|x| \ln (R / r)} \tag{57}
\end{equation*}
$$

holds on $\mathbb{R}^{n}$ with a certain positive constant $c_{4}$ for arbitrary $R>r>1$. It then follows from (54) and (57) that the inequality

$$
\begin{equation*}
c_{5} \int_{B(R) \backslash B(r)}(|x| \ln (R / r))^{-n} \mathrm{~d} x \geqslant \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{58}
\end{equation*}
$$

holds, and, therefore, so does the inequality

$$
\begin{equation*}
c_{6}(\ln (R / r))^{-n+1} \geqslant \int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-v-1} \mathrm{~d} x \tag{59}
\end{equation*}
$$

with arbitrary $R>r>1$ and certain positive constants $c_{5}$ and $c_{6}$ that do not depend on $R$. Passing to the limit as $R \rightarrow+\infty$ in (59), we find that the equality

$$
\begin{equation*}
\int_{B(r)} \sum_{i=1}^{n} u_{x_{i}} A_{i}(x, u, \nabla u)(u-c+\varepsilon)^{-\nu-1} \mathrm{~d} x=0 \tag{60}
\end{equation*}
$$

holds with an arbitrary $r>1$. Passing to the limit as $r \rightarrow+\infty$ in (60), we again obtain, due to the non-negativity of the integrand in (60) and by condition (6), that $u(x)=$ const. on $\mathbb{R}^{n}$.

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